

# Theoretical Developments in Understanding Massive Star Formation

presented at **Massive Star Formation:  
Observations Confront Theory**



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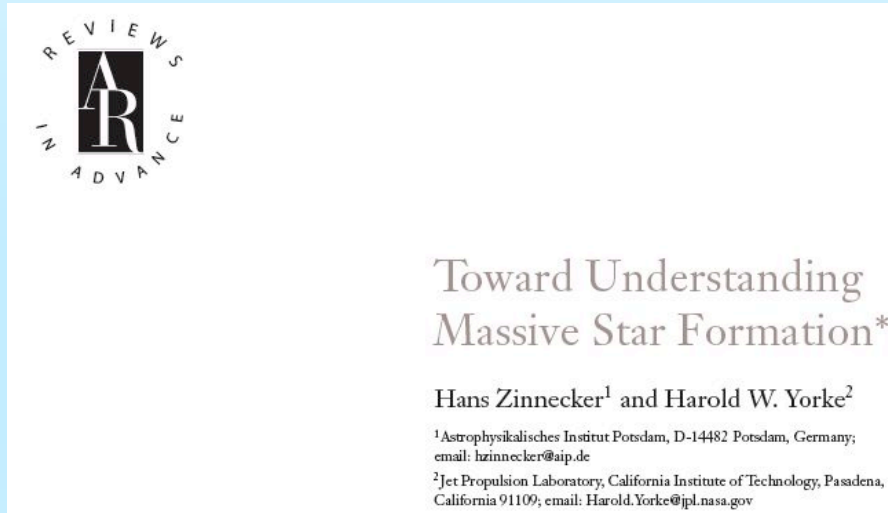
*Morse & Davidson (1996)*

Jet Propulsion Laboratory  
California Institute of Technology

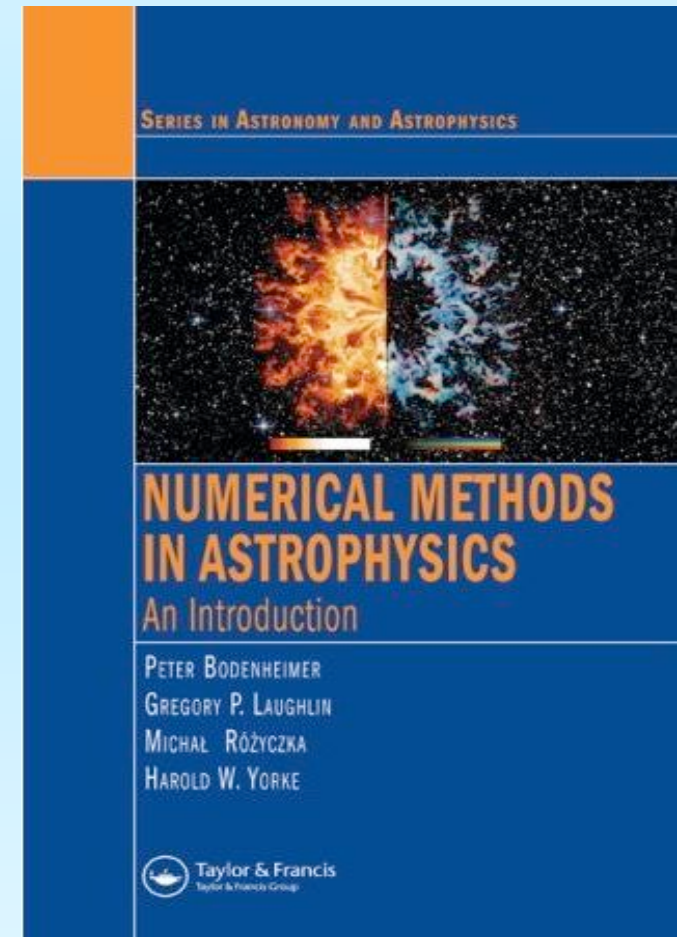
**10-14 September 2007**



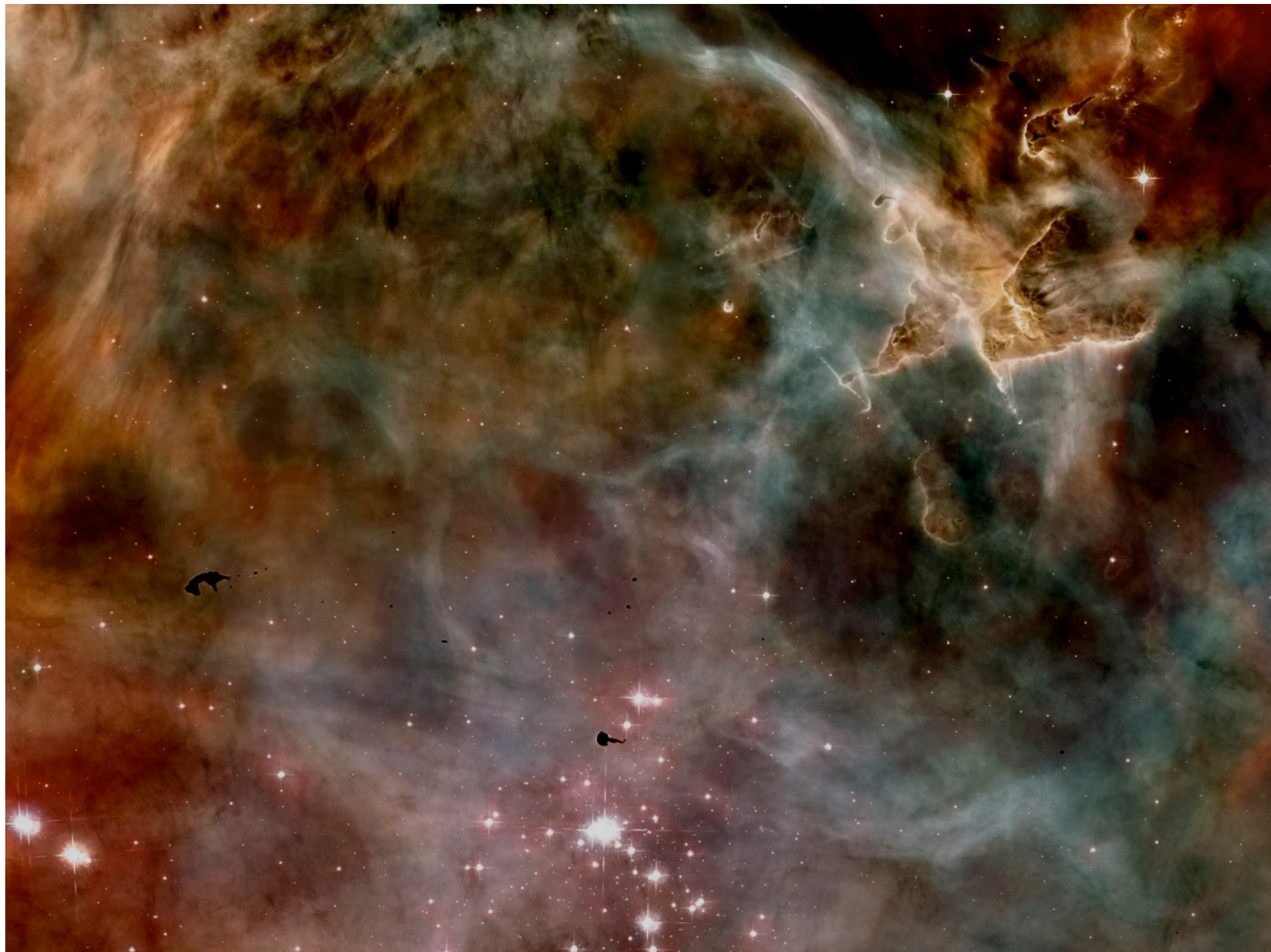
- Zinnecker & Yorke (2007) *Toward Understanding Massive Star Formation*, [Ann. Rev. Astron. Astrophys. 45: 481-563 \(arXiv:0707.1279\)](#)



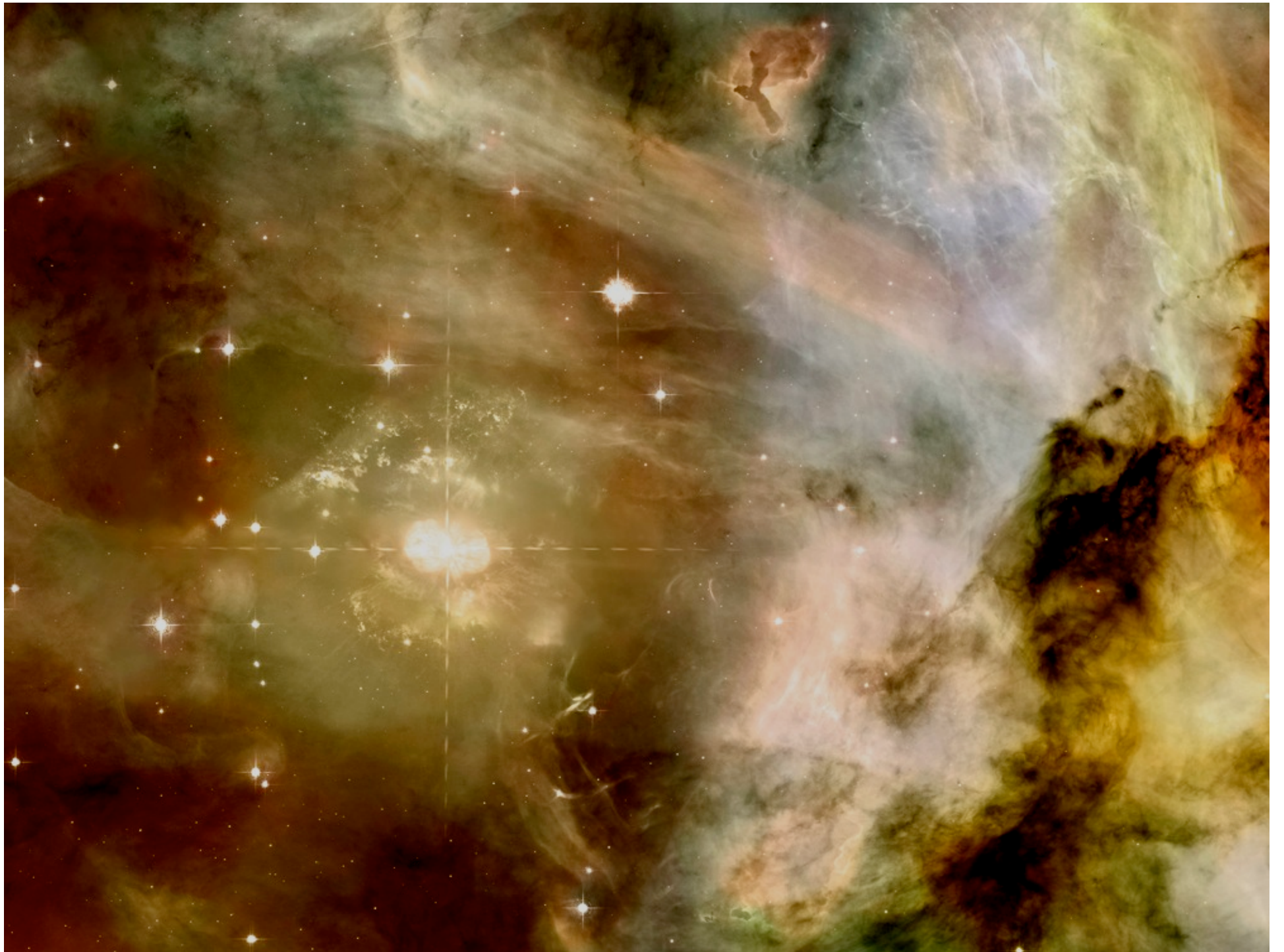
- Bodenheimer, Laughlin, Rozyczka, Yorke (2007) [Numerical Methods in Astrophysics : An Introduction](#), Taylor & Francis



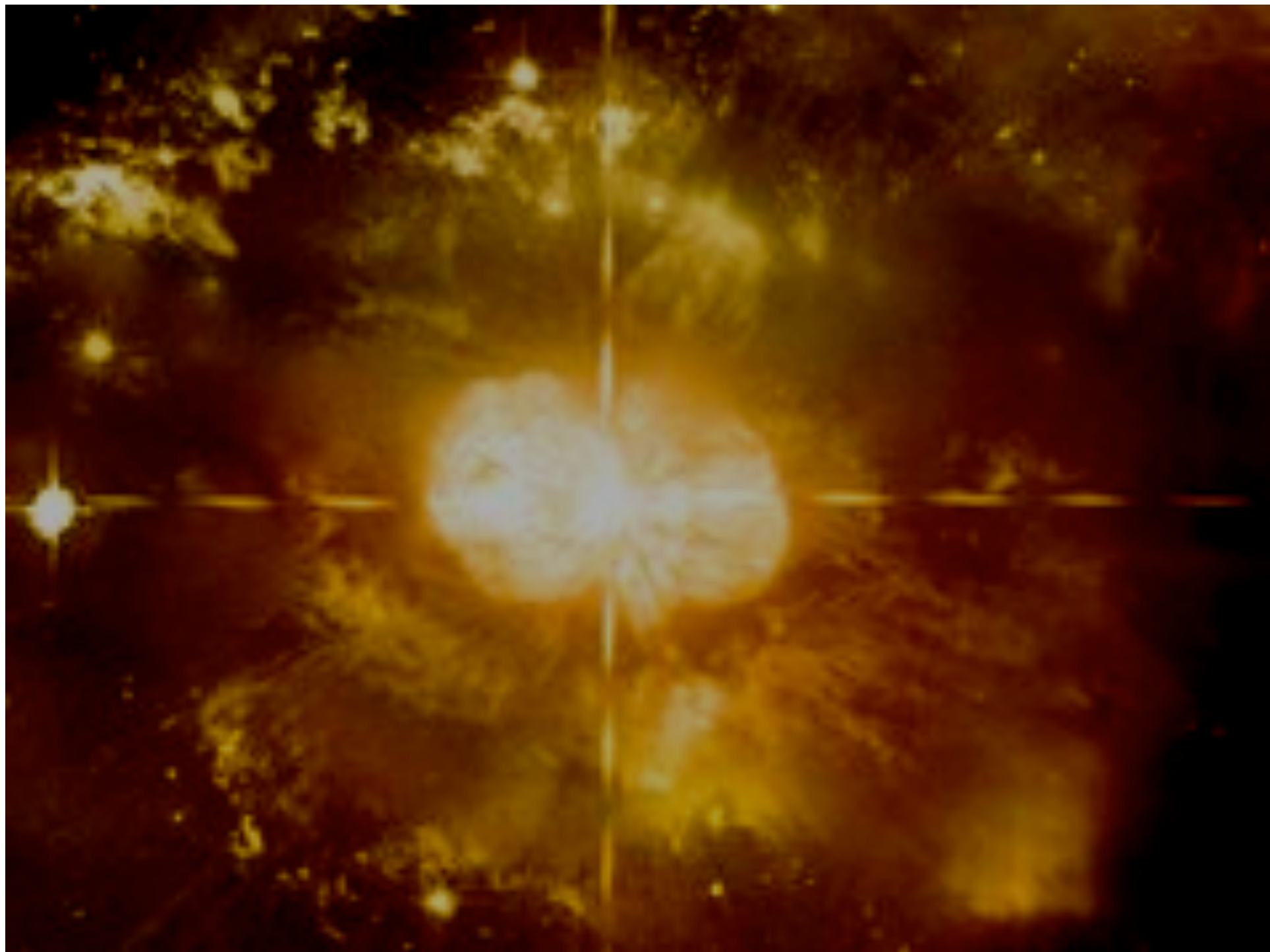














# Massive Star Formation: Observations Confound Theory (1/2)

- Universality of IMF and upper mass limit
  - No clear-cut evidence of variation of slope between high mass through intermediate mass to solar-type stars
- Multiplicity, Hierarchies, Clusters, Associations
  - High mass stars generally form in clusters & associations
    - In loose OB associations (Ori OB1a,b; Sco OB2; NGC 604)
    - In dense clusters (Ori TC; NGC 3603; 30 Dor); most O-stars located in center
    - Starburst galaxies, ULIRGs
  - Higher degree of multiplicity of high mass stars than for low mass stars
    - Average number of companions  $\sim 1.5$  for massive primary, whereas  $\sim 0.5$  for solar-type primaries
    - O-stars have preponderance of close tight binaries with  $P \sim 3-5$  days
    - Higher fraction of runaway O-stars than runaway B-stars
    - Reduced binarity among runaway O-stars compared to cluster O-stars

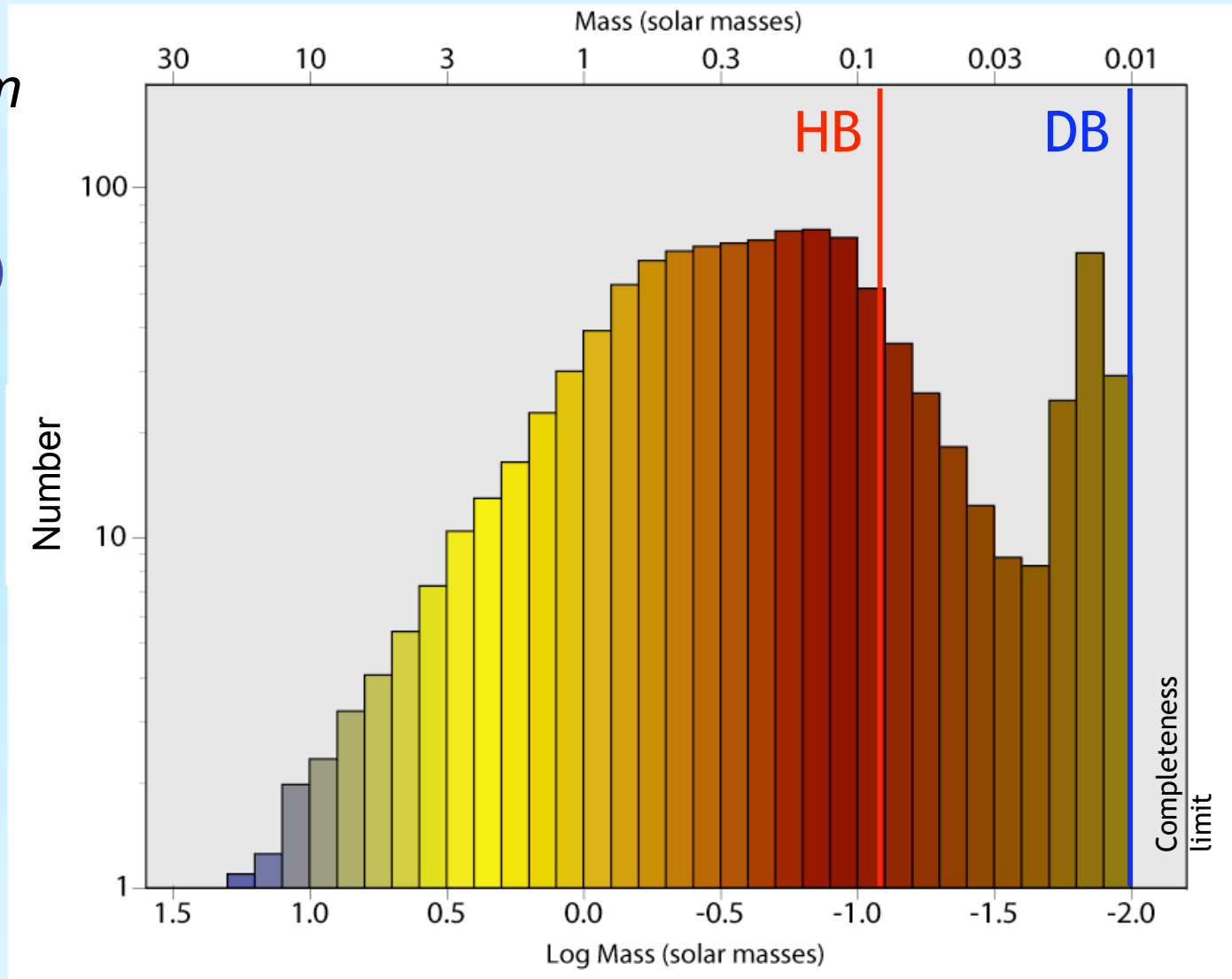


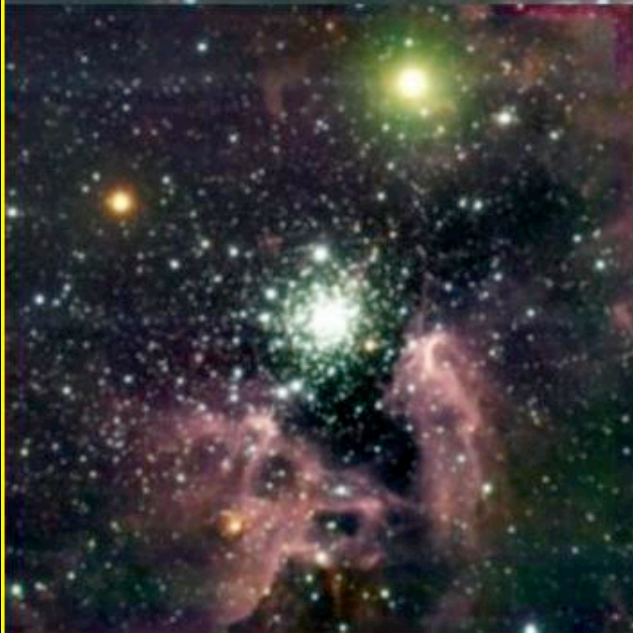
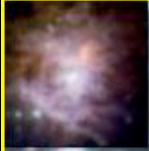


# The Initial Mass Function



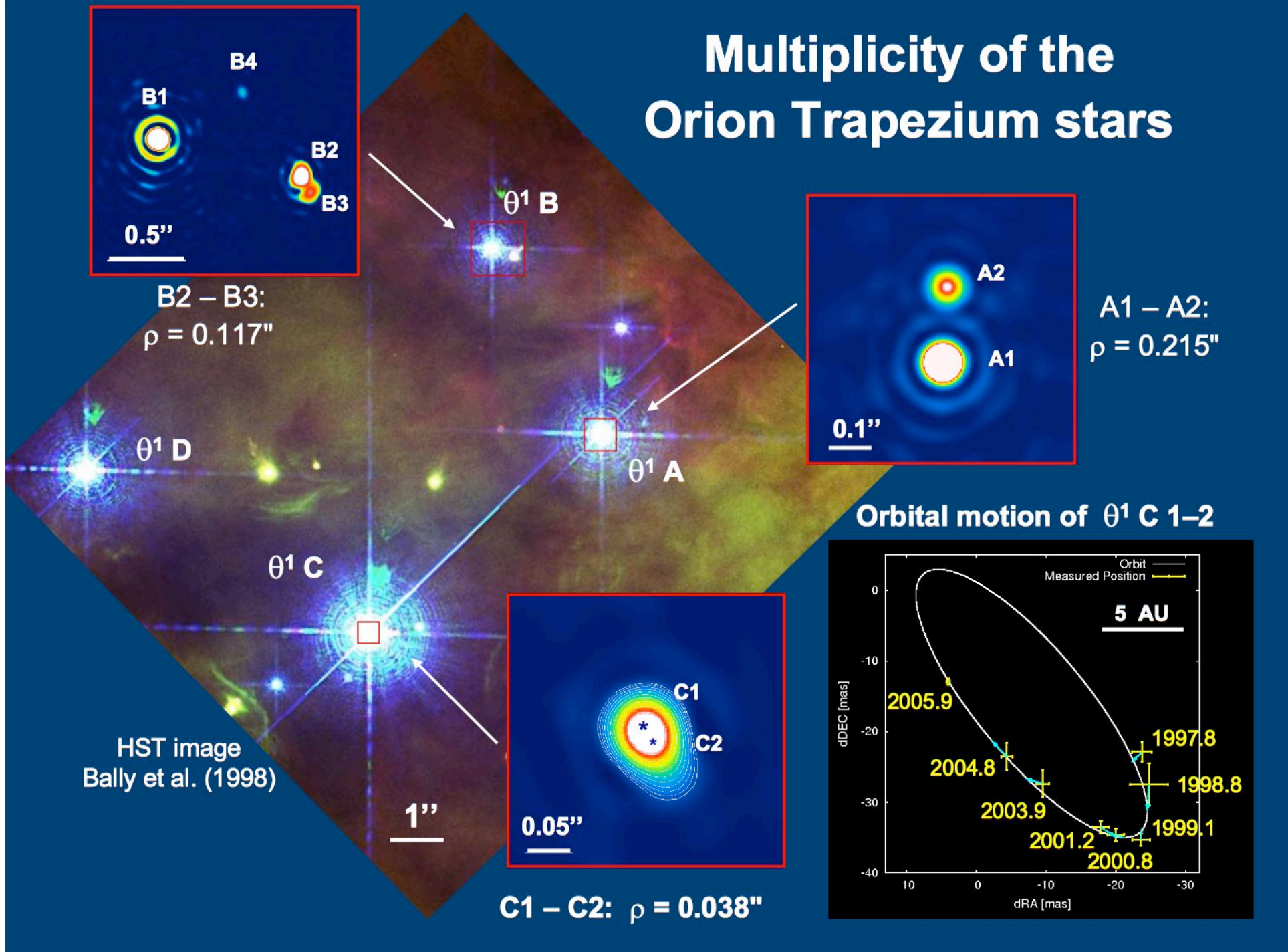
*Trapezium cluster*  
(Muench et al. 2001)







# Multiplicity of the Orion Trapezium stars





# Massive Star Formation:



## Observations Confound Theory (2/2)

- Characteristics of OB-star forming regions
  - Hot molecular cores in GMCs
  - Hypercompact and Ultracompact HII regions
  - Masers (OH, H<sub>2</sub>O, SiO, CH<sub>3</sub>OH) - in disks or in outflows?
  - Both wide-angle and collimated flows observed: jets and outflows
  - Cometary proplyds, pillars, mountains
  - Turbulence observed; enough to support clumps and cores?
- No disk around an optically visible main sequence O-star found
  - Disks around B-stars have been observed
- There are massive molecular cores or clumps without outflows
  - There are hot cores without outflows and without radio continuum
- Some magnetic field measurements; when measurable sub/super-critical within factor 2
  - $\Theta$ 1 Ori C is a magnetic star!





# The Global Picture

- Assume:  $dN = A m^{-a} dm$  with  $a=2.35$  (Salpeter) between  $m=0.1$  and  $m=100$ 
  - Average mass  $\langle m \rangle = 0.35$
- There is one  $>50 M_{O\text{-star}}$  for every 7300 stars formed
- In MWG: 2 SN/100 yr ( $m > 8$ )  $\Rightarrow$  8 stars/yr formed
  - 3  $M_{O}$ /yr converted into stars
  - Every 50 yr produce  $>8 M_{O\text{-star}}$
  - Every 200 yr produce  $>20 M_{O\text{-star}}$
  - Every 400 yr produce  $>30 M_{O\text{-star}}$
  - Every 1000 yr produce  $>50 M_{O\text{-star}}$



# Molecular Cloud Lifetime, SF Efficiency, Total Molecular Mass (1/3)



Assume:

- $10^{10} M_{\odot}$  ISM, of which 40% is molecular
- $10^7$  yr lifetime of molecular clouds

=>  $400 M_{\odot}/\text{yr}$  of molecular cloud material must be dissipated and  $400 M_{\odot}/\text{yr}$  must be newly formed

Star formation efficiency =  $3/400 = 0.7\%$



# Molecular Cloud Lifetime, SF Efficiency, Total Molecular Mass (2/3)



Assume:

- 400  $M_{\odot}$ /yr of molecular cloud material must be dissipated

Use  $>30 M_{\odot}$  O-stars, of which you have  $\sim 8000$  in MWG

$\Rightarrow$  Each O-star must dissipate  $0.05 M_{\odot}$ /yr

Champagne flows driven by O5-stars can dissipate  $0.01 M_{\odot}$ /yr (Yorke 1986)





# Molecular Cloud Lifetime, SF Efficiency, Total Molecular Mass (3/3)



Assume

- $10^{10} M_{\odot}$  ISM, of which 40% is molecular
- Star Formation efficiency of 50%

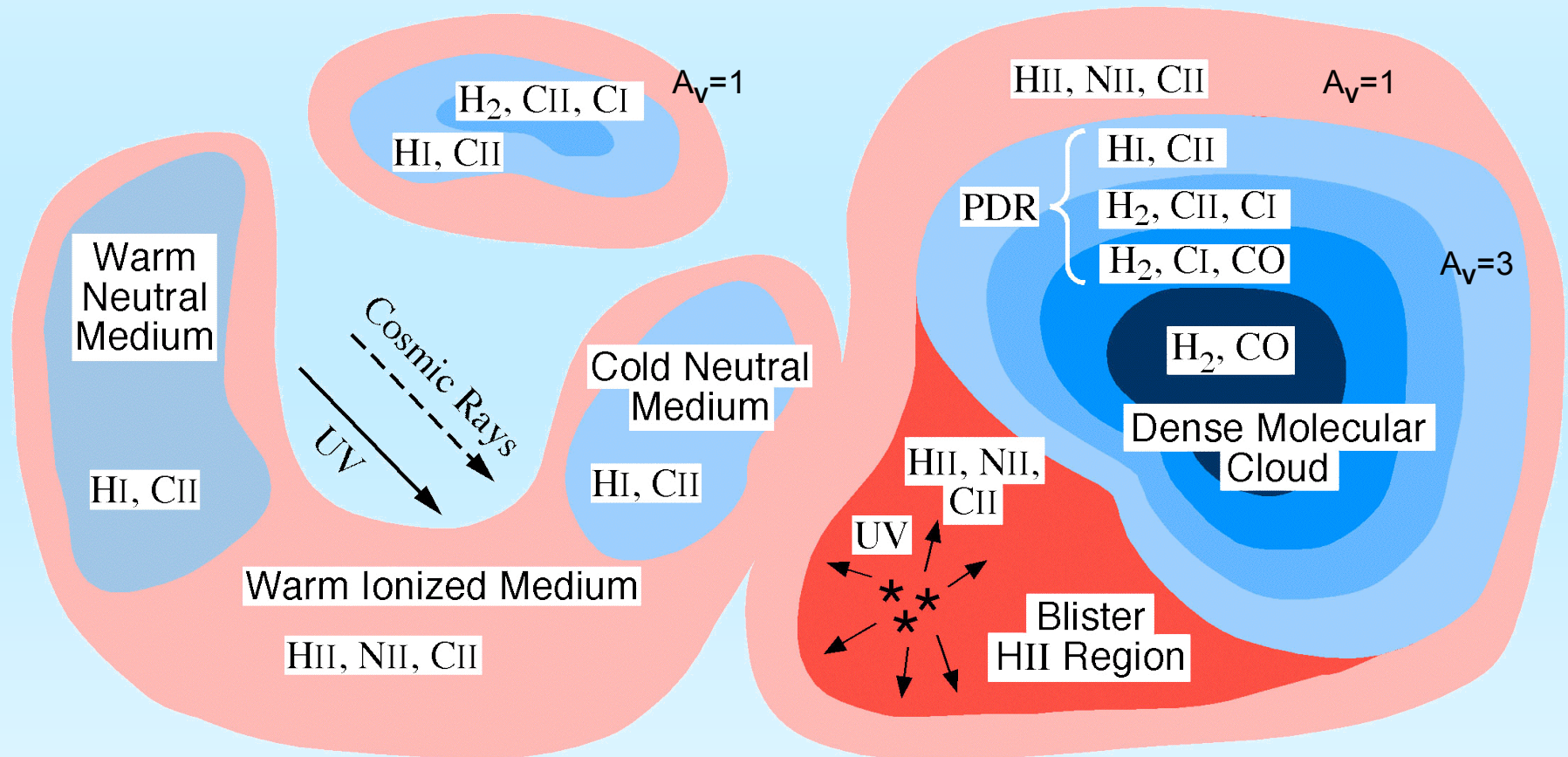
3  $M_{\odot}$ /yr of molecular cloud material will be dissipated and 6  $M_{\odot}$ /yr must be newly formed

Lifetime of molecular material:  $4 \times 10^9 / 6 = 7 \times 10^8$  yr

=> at least some molecular material is long-lived



# Is there a hidden Component of $H_2$ ?





# Steps to produce Massive Stars

- Create Giant Molecular Cloud Complexes  $\sim 10\text{-}100$  pc
- Create molecular cloud clumps  $\sim 1$  pc and cores  $\sim 0.1$  pc
- Initiate collapse of cores/filaments  $\sim 0.1$  pc
- Fragment into several sub-clumps  $\sim 0.01$  pc
- Create first hydrostatic cores  $\sim 1 R_{\odot} = 2.3 \times 10^{-8}$  pc
- Accrete onto hydrostatic cores through disks, allowing them to grow in mass
  - Accretion Disks  $\sim 10^{-3}$  pc
  - Accretion columns  $\ll 10^{-8}$  pc
- Hydrostatic cores evolve quickly to H-burning, even while accreting
- Remnant disks quickly dissipate as accretion halts







# 350 $\mu\text{m}$ Cores in the Orion Region



-1 3 7 11 15 20 Jy per beam

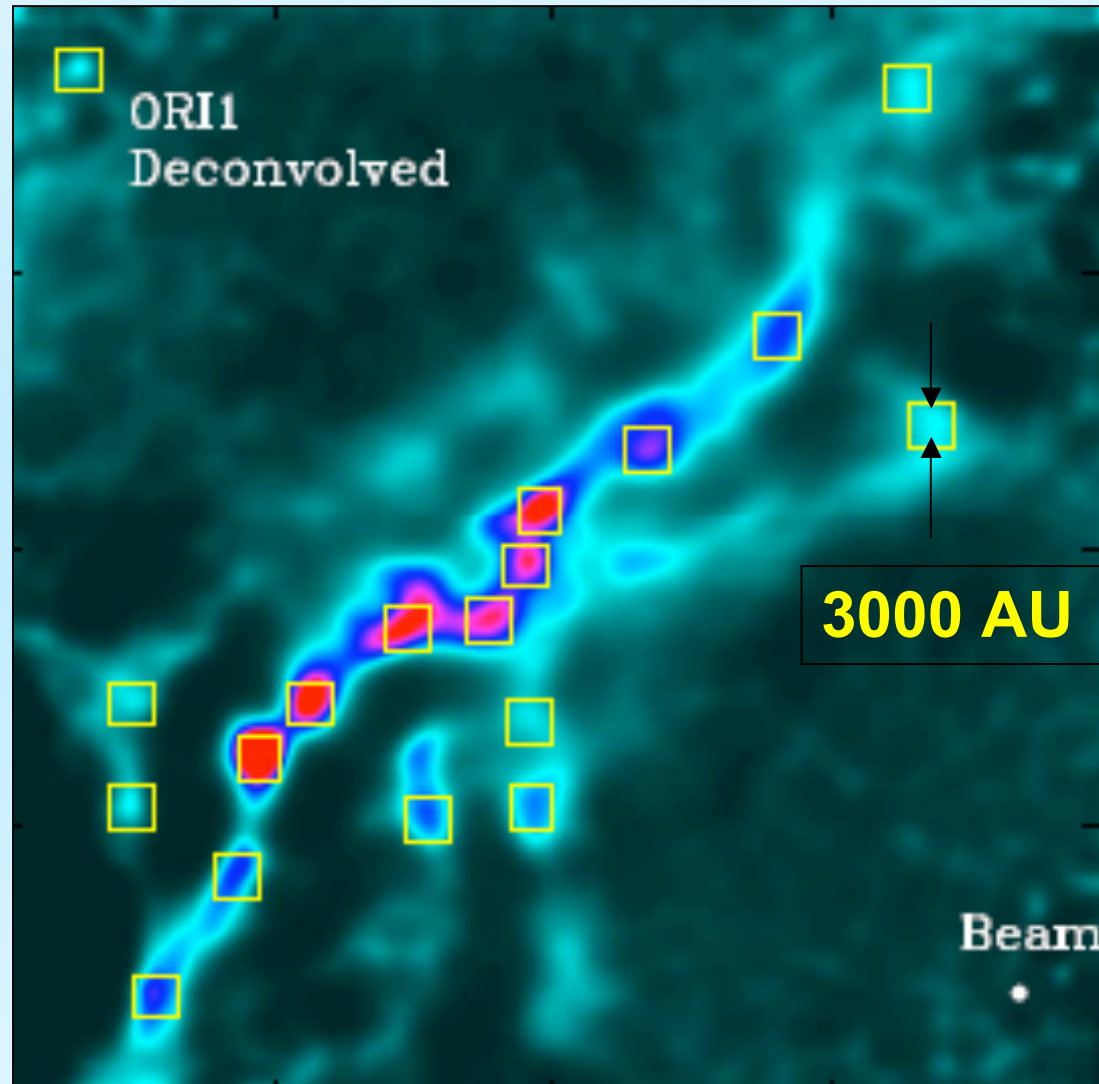


-4:57:50.4

51 cores 0.1 to 46  $M_{\text{sun}}$

Many appear to be unstable against gravitational collapse

Li, Velusamy, Goldsmith, & Langer (2007)



-5:03:09.2

5:35:28.23

H.W. Yorke 5:35:06.89



# Questions which need to be answered

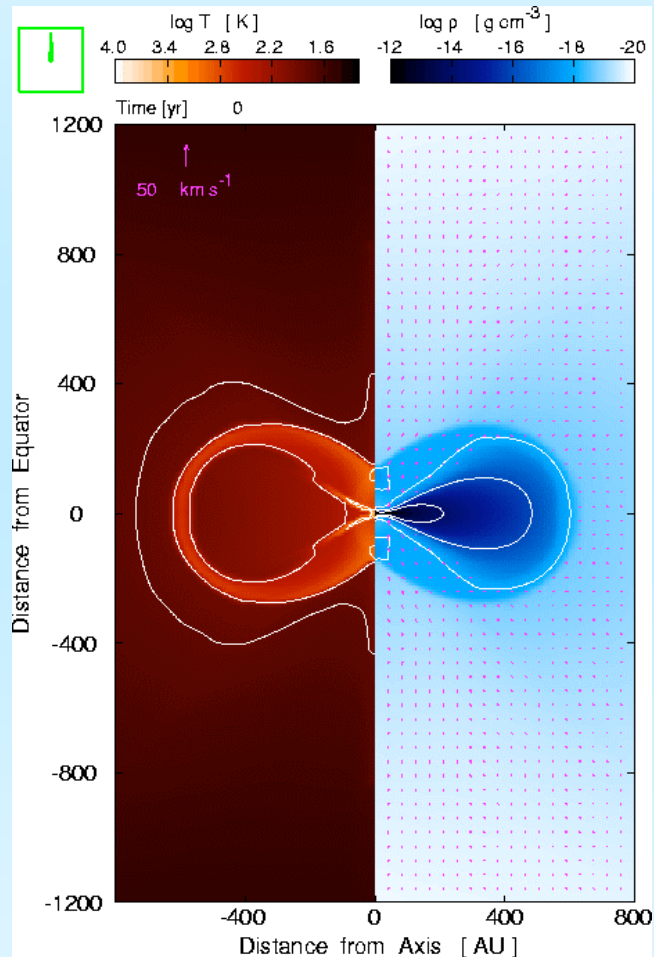


- Assuming a gravitationally unstable massive clump ( $10^{-1}$  pc), how does enough material become concentrated into a sufficiently small volume (a few  $10^{-8}$  pc) within a sufficiently short time ( $\sim 10^5$  yr)?
- How does the forming massive star influence its immediate surroundings to limit its mass?
- Today I will not try to answer these questions



# Destruction of circumstellar disks

- Destruction of circumstellar disk:



## Initial conditions

Luminosity of star:	3550 $L_{\odot}$
Mass of central star:	8.3 $M_{\odot}$
Stellar wind:	30 $\text{km s}^{-1}$ $10^{-8} M_{\odot} \text{ yr}^{-1}$
Mass of disk:	0.7 $M_{\odot}$
H-ionizing flux:	$7 \times 10^{44} \text{ s}^{-1}$
Net UV flux:	$2 \times 10^{48} \text{ s}^{-1}$

## Results

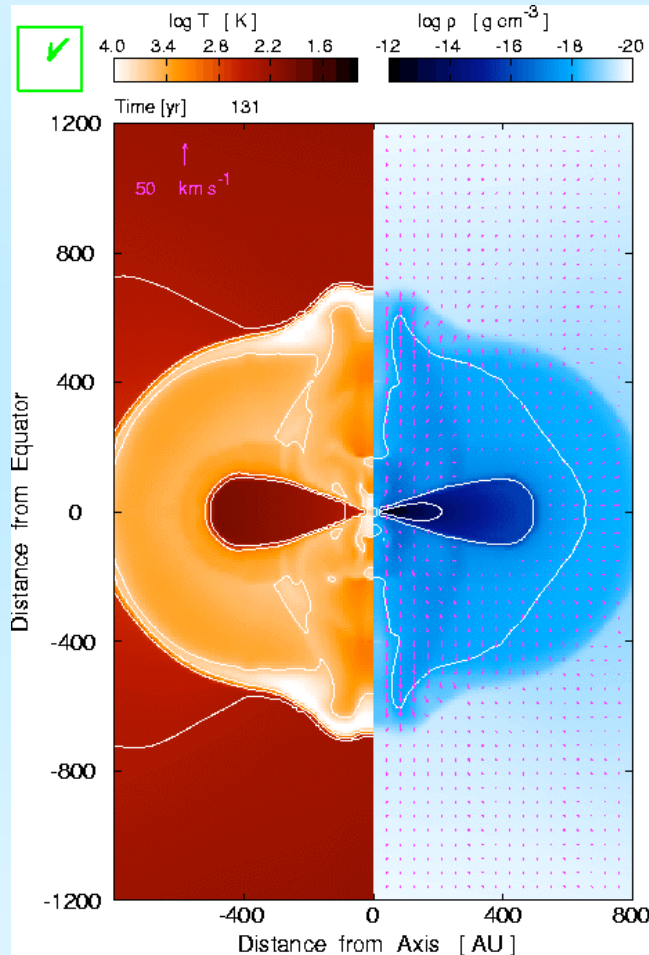
Disk wind via “external” UV heating  
Stellar wind focused through polar cavity  
Photoevaporation of disk within  $10^5$  years





# Destruction of circumstellar disks **JPL**

- Destruction of circumstellar disk:



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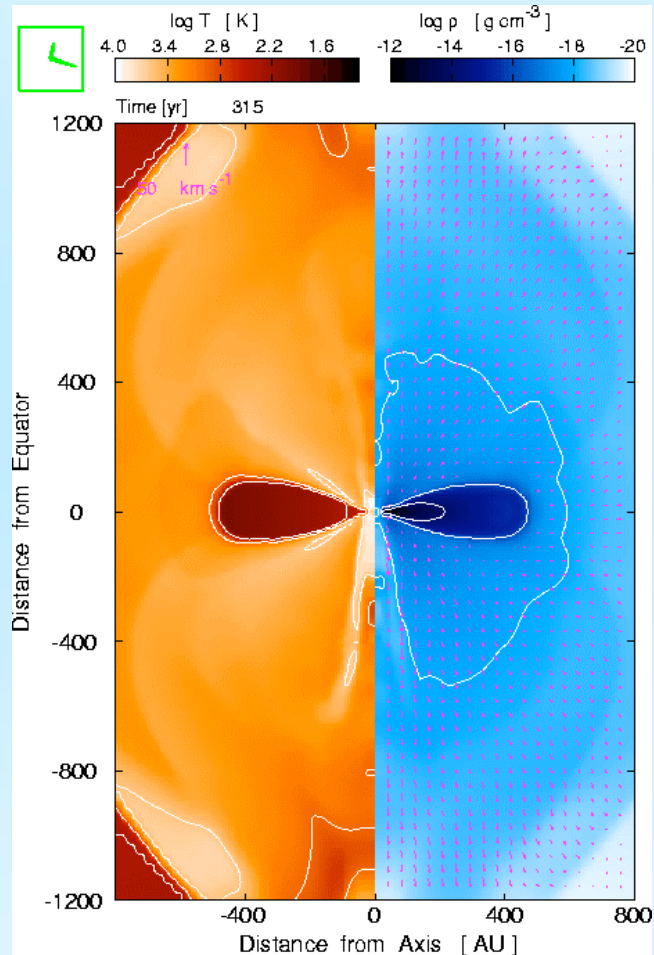
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Net UV flux:	2 x 10 <sup>48</sup> s <sup>-1</sup>

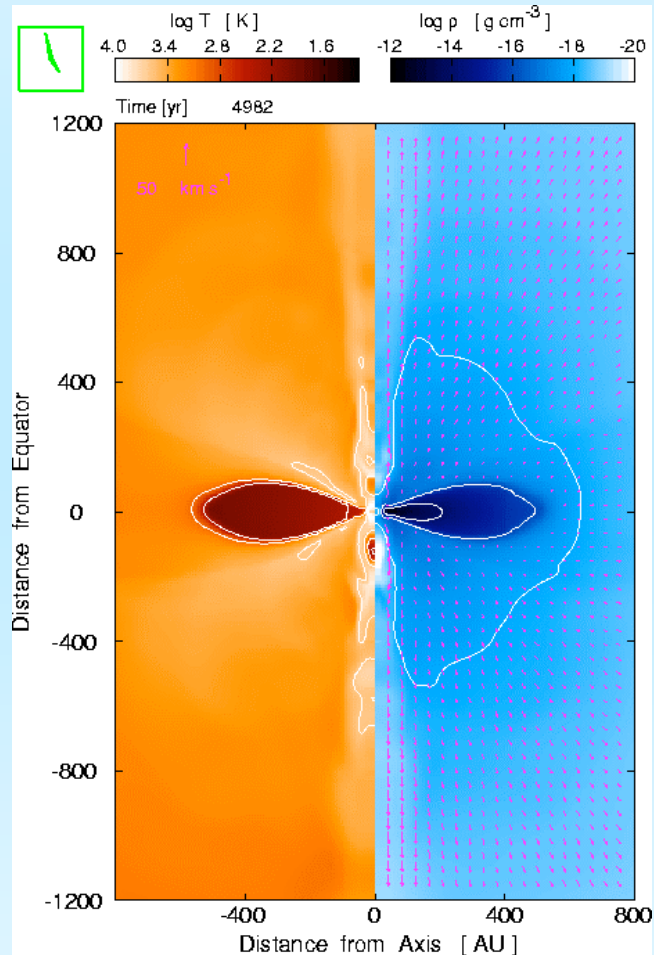
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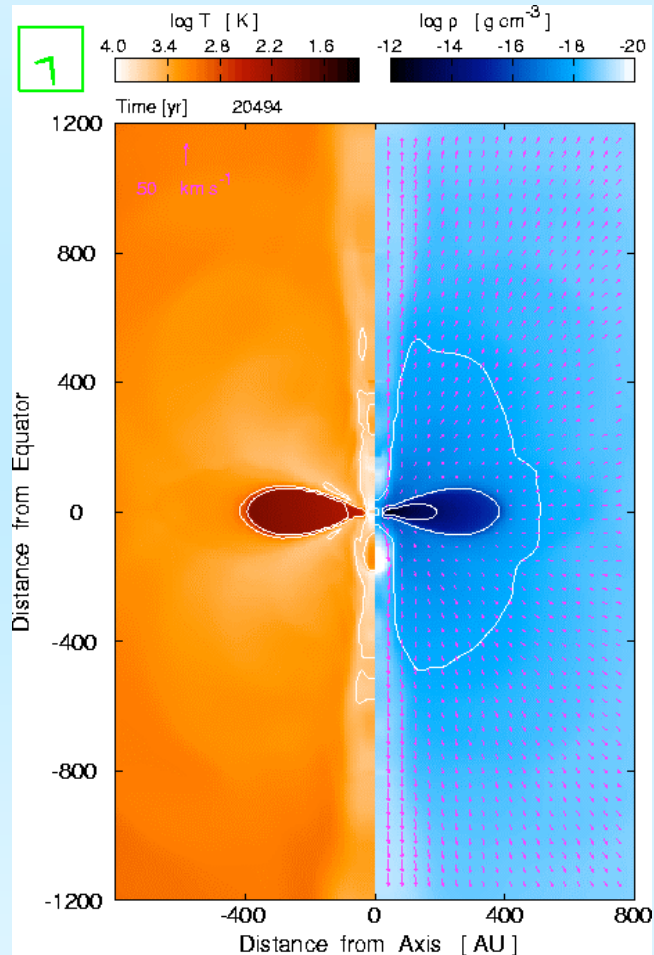
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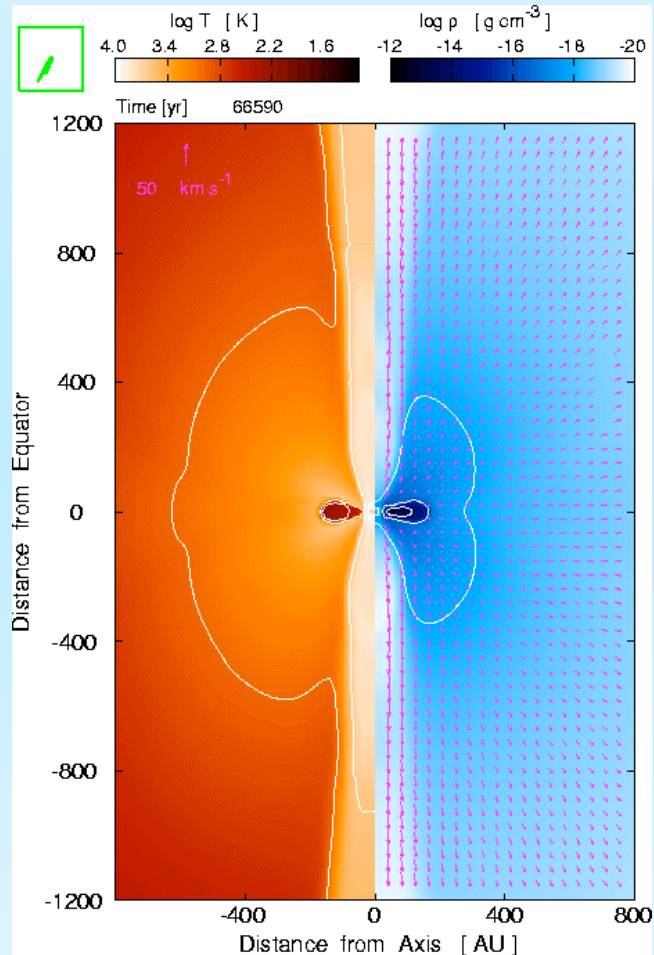
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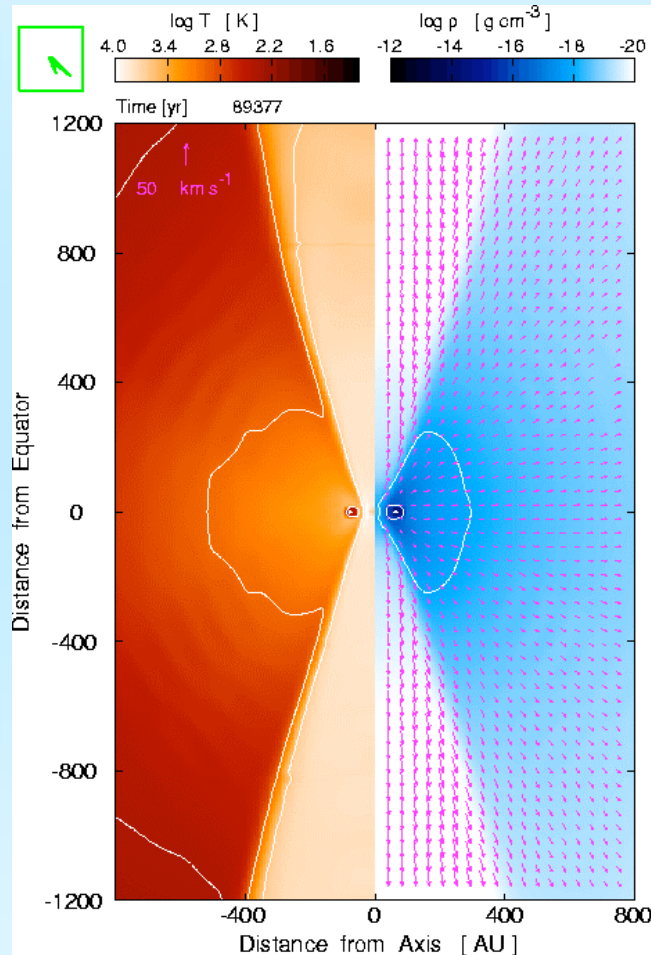
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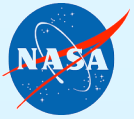


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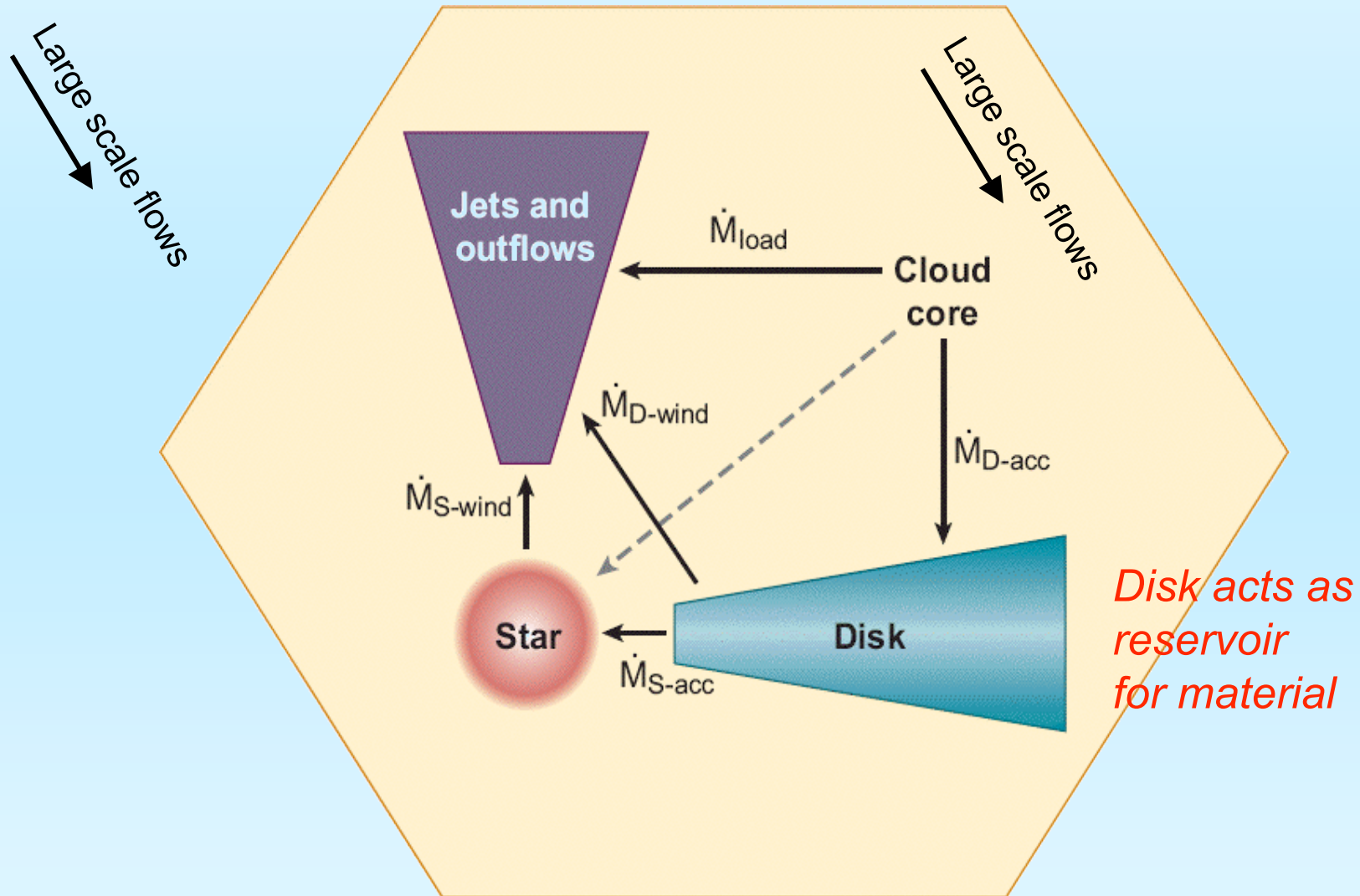
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# Accretion and mass loss as mass exchange between components







# Necessary conditions for forming stars by accretion



- The mass gained by accretion must exceed losses

$$M_* (t) = \int^t [ \dot{M}_{\text{disk}}(t') - \dot{M}_{\text{S-wind}}(t') ] dt'$$

$$M_{\text{disk}} (t) = \int^t [ \dot{M}_{\text{acc}}(t') - \dot{M}_{\text{disk}}(t') - \dot{M}_{\text{D-wind}}(t') ] dt'$$

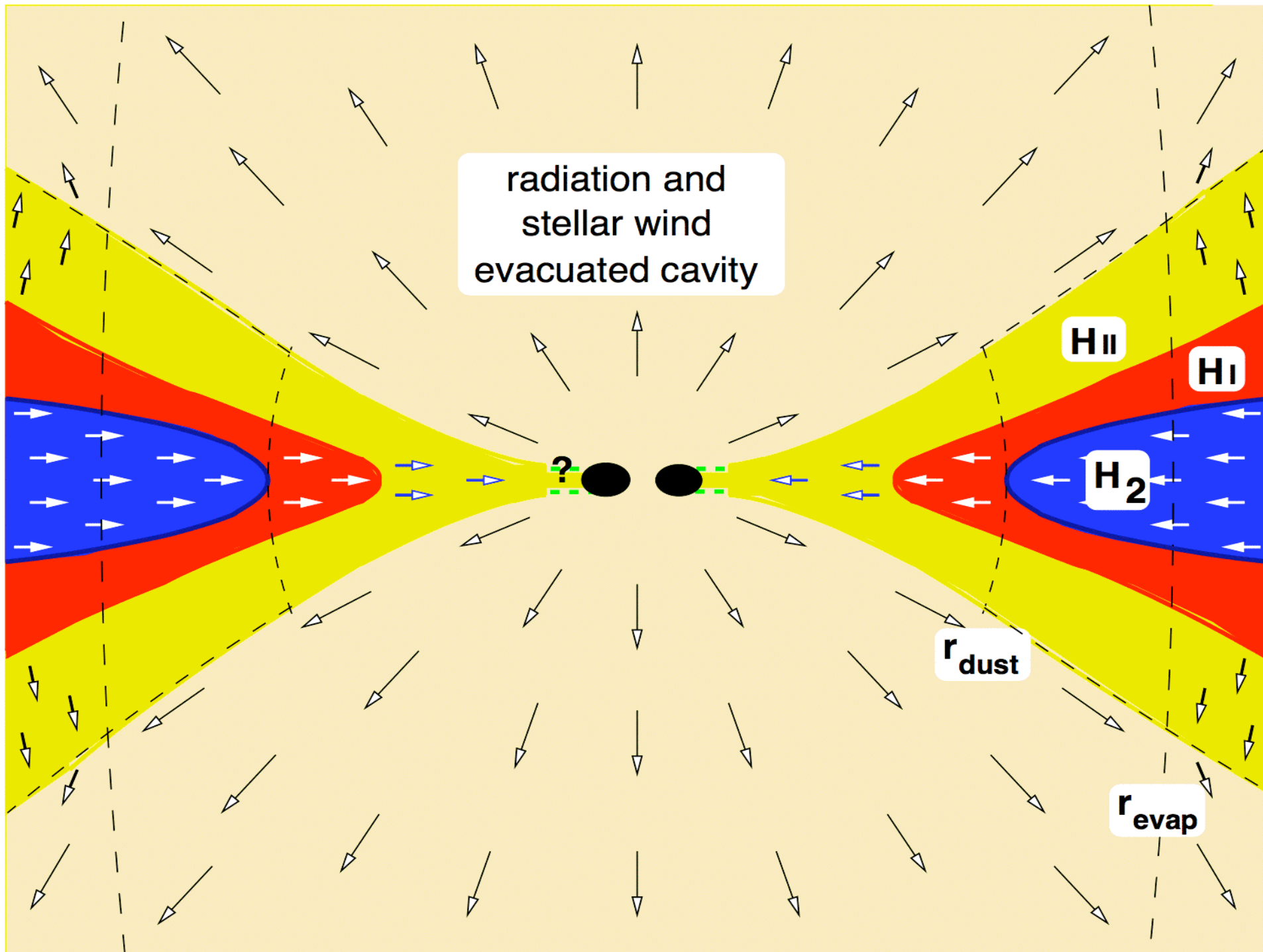
- Must accrete material within “reasonable” time

$$\dot{M}_{\text{acc}} - \dot{M}_{\text{S-wind}} - \dot{M}_{\text{D-wind}} \sim M_* / t_{\text{acc}}$$

- Gravity must be dominant force

$$\frac{GM}{r^2} > \frac{\kappa_{\text{eff}} L}{4\pi r^2 c}, \quad \text{where} \quad L = L_* + L_{\text{acc}}$$

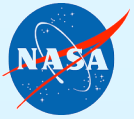
$$\kappa_{\text{eff}} < 130 \text{ cm}^2 \text{ g}^{-1} \left[ \frac{M}{10 M_{\odot}} \right] \left[ \frac{L}{1000 L_{\odot}} \right]^{-1}$$





# What determines these $dM/dt$ 's?

- $\dot{M}_{\text{disk}}$  related to angular momentum transport within disk
  - Magnetic fields
  - Tidal effects (bars, spiral arms)
  - Turbulence (photon bubbles: Turner, Quataert, Yorke 2007)
- $\dot{M}_{\text{acc}}$  determined by cloud core parameters, competitive accretion
  - $t_{\text{acc}} \sim t_{\text{ff}} \sim t_{\text{cross}}$
- $\dot{M}_{\text{D-wind}}$  has contributions from jets and photoevaporation
  - $\dot{M}_{\text{jet}} \sim f(\dot{M}_{\text{disk}})$
  - $\dot{M}_{\text{photoevap}} \sim f(F_{\text{v}})$
- $\dot{M}_{\text{S-wind}}$  related to stellar parameters



# Important relevant time scales



- Free-fall time scale

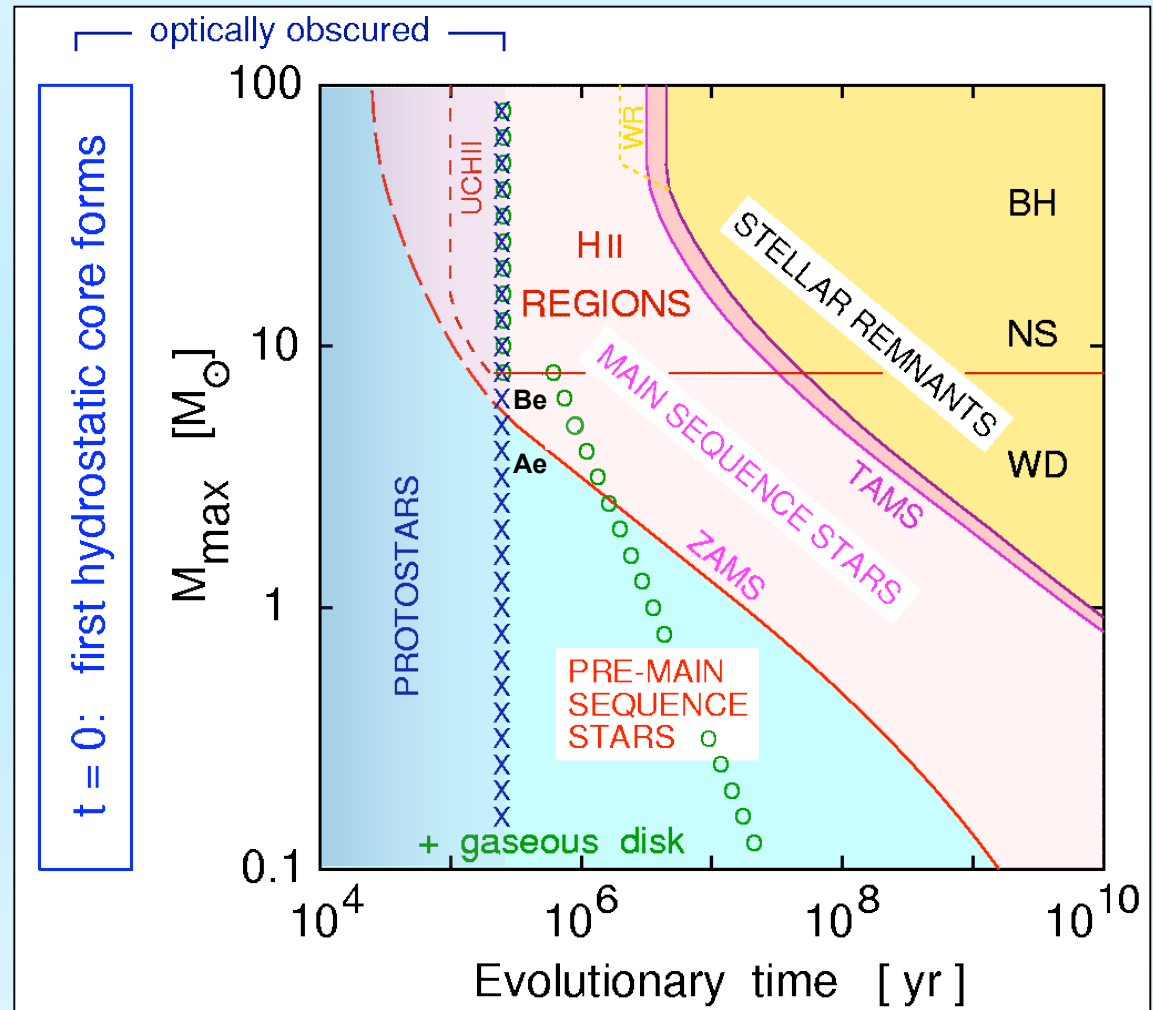
$$t_{\text{ff}} \simeq 10^5 \text{ yr} \left[ \frac{n}{10^5 \text{ cm}^{-3}} \right]^{-1/2}$$

- Accretion time scale

$$t_{\text{acc}} \simeq t_{\text{ff}}$$

- Kelvin-Helmholtz time scale (time to reach ZAMS)

$$\tau_{\text{KH}} = \frac{GM^2/R}{L} \sim M^{-2}$$





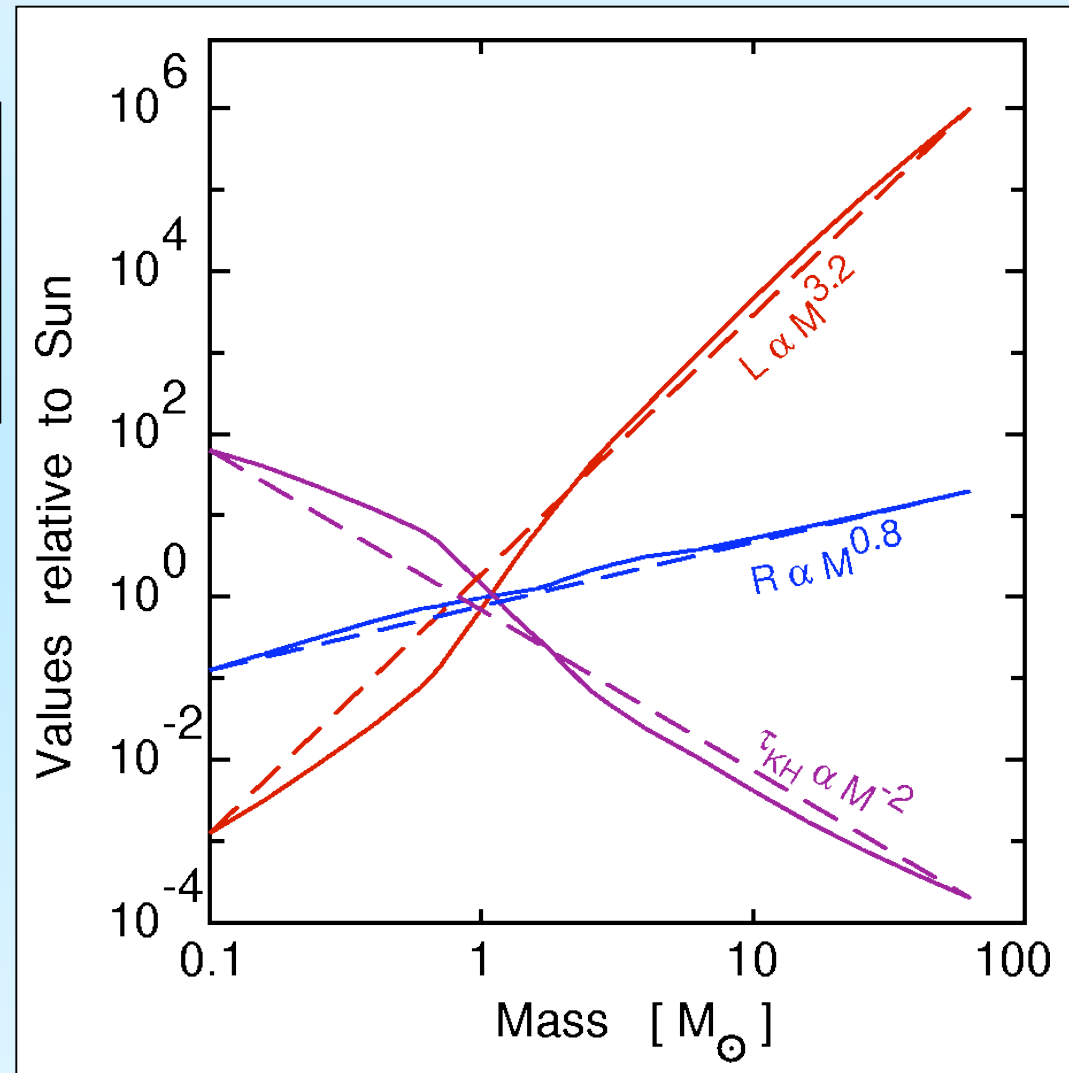


# Luminosity & Radius of ZAMS star



**Kelvin-Helmholtz timescale**  
(for thermal readjustment)

$$\tau_{\text{KH}} = \frac{GM^2/R}{L}$$





# Sources of Luminosity of accreting Stars



- Accretion luminosity:  $L_{\text{acc}} = GM_*/R_* dM/dt$

$$L_{\text{acc}} = 6000 L_{\odot} [M_*/30 M_{\odot}]^{0.2} [dM/dt / 10^{-4} M_{\odot}/\text{yr}]$$

- Deuterium burning

$$L_{\text{D}} = 400 L_{\odot} [dM/dt / 10^{-4} M_{\odot}/\text{yr}]$$

- PMS Contraction

$$L_{\text{KH}} = GM_*^2/R_*^2 dR/dt$$

- Hydrogen burning

$$L_* = 10^5 L_{\odot} [M_*/30 M_{\odot}]^{3.2}$$



# Important relevant length scales



- Size of clump to produce star of mass  $M$

$$R_{\text{clump}} \simeq 0.1 \text{ pc} \left[ \frac{M}{30 M_{\odot}} \right]^{1/3} \left[ \frac{n}{10^5 \text{ cm}^{-3}} \right]^{-1/3}$$

- Photoevaporation radius

$$r_{\text{evap}} \simeq 130 \text{ AU} \left[ \frac{M}{30 M_{\odot}} \right]$$

- Dust destruction radius

$$r_{\text{dust}} \simeq 25 \text{ AU} \left[ \frac{L}{1.5 \cdot 10^5 L_{\odot}} \right]^{1/2}$$

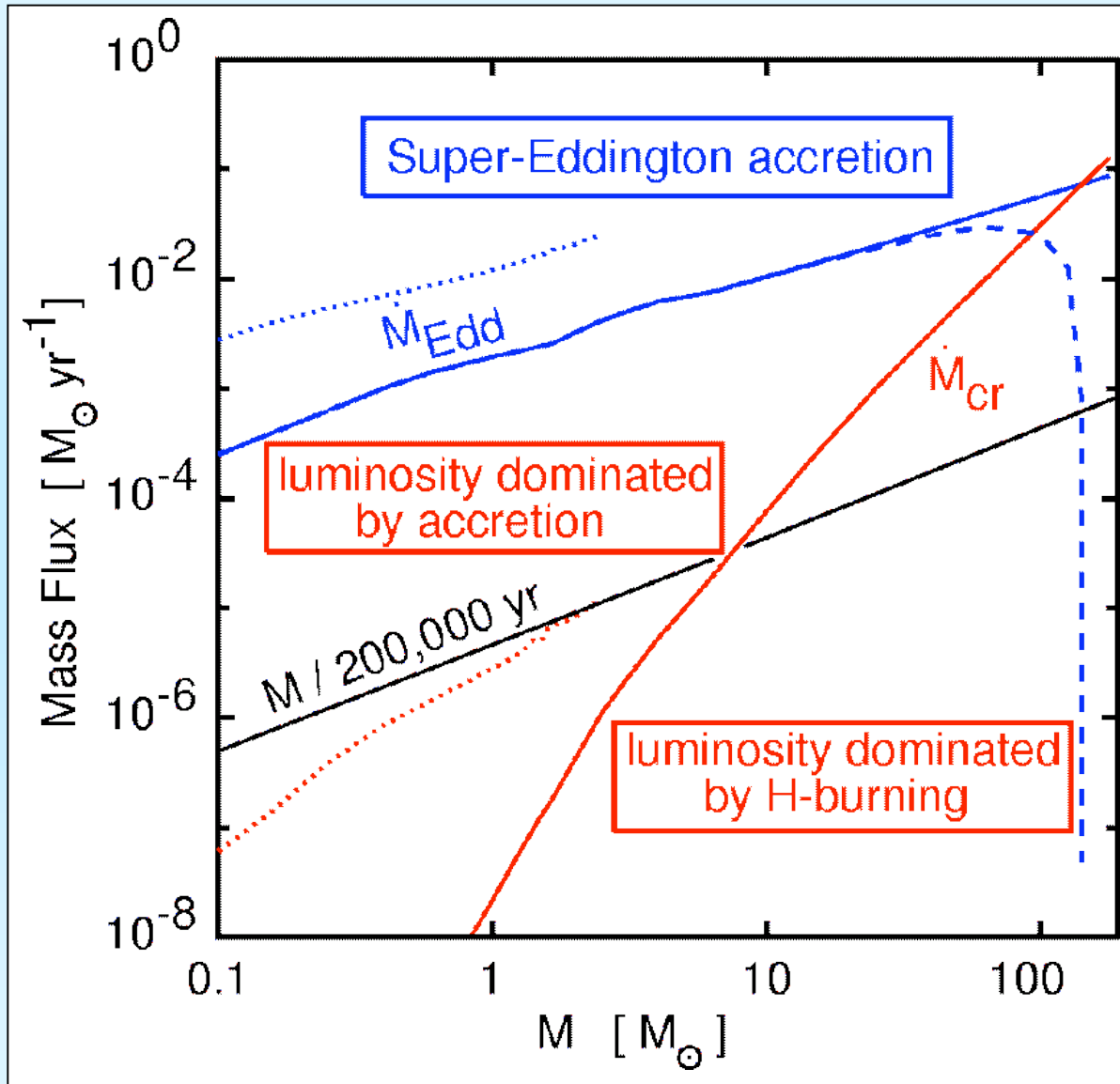
$$r_{\text{dust}} \simeq 25 \text{ AU} \left[ \frac{M}{30 M_{\odot}} \right]^{1.6}$$

- Stellar radius

$$R_{*} \simeq 12 R_{\odot} \left[ \frac{M}{30 M_{\odot}} \right]^{0.8}$$



# The Eddington accretion limit



Assume electron scattering for minimum opacity...

$$\frac{\kappa_e L_{\text{Edd}}}{4\pi r^2 c} = \frac{GM}{r^2}$$

$$L_{\text{Edd}} = L_* + \frac{GM}{R} \dot{M}_{\text{max}}$$

$$\dot{M}_{\text{max}} = \frac{4\pi R c}{\kappa_e} - \frac{RL_*}{GM}$$

$$\dot{M}_{\text{max}} = \dot{M}_{\text{Edd}} - \dot{M}_{\text{cr}}$$



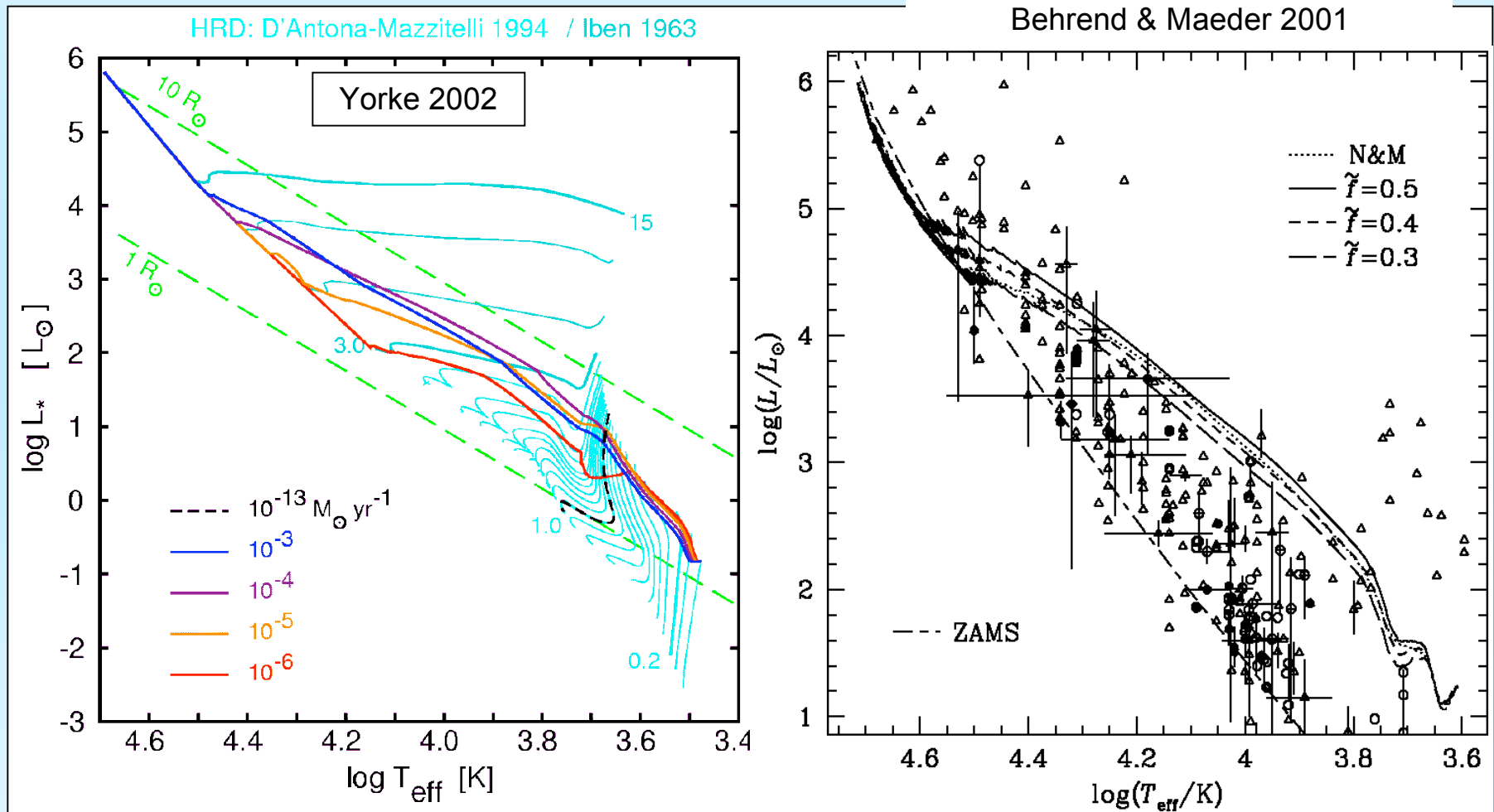


# Evolution of accreting stars in the

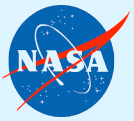


## HRD ( $dM_*/dt > 0$ )

(Yorke 2002; Behrend & Maeder 2001)



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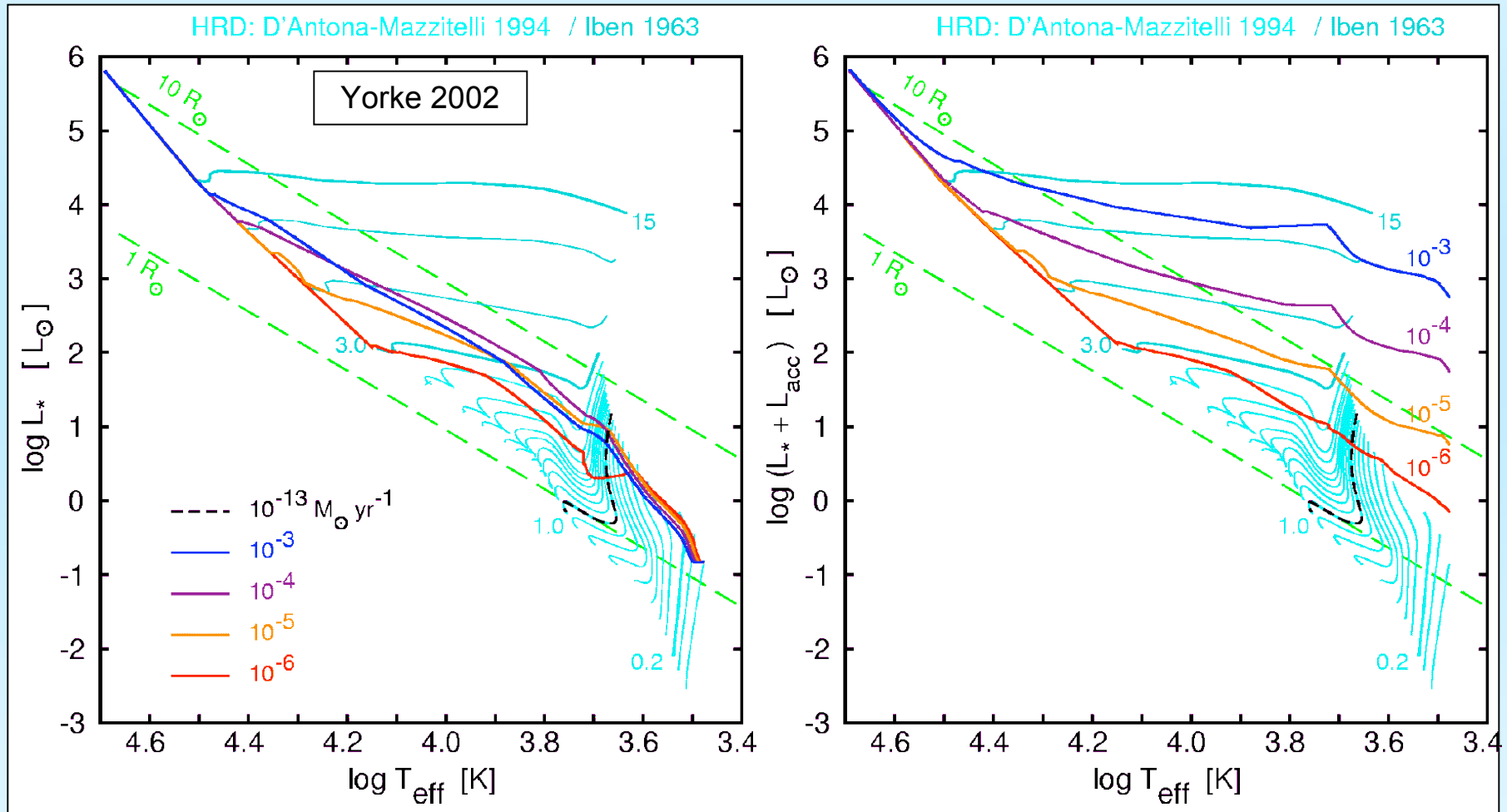


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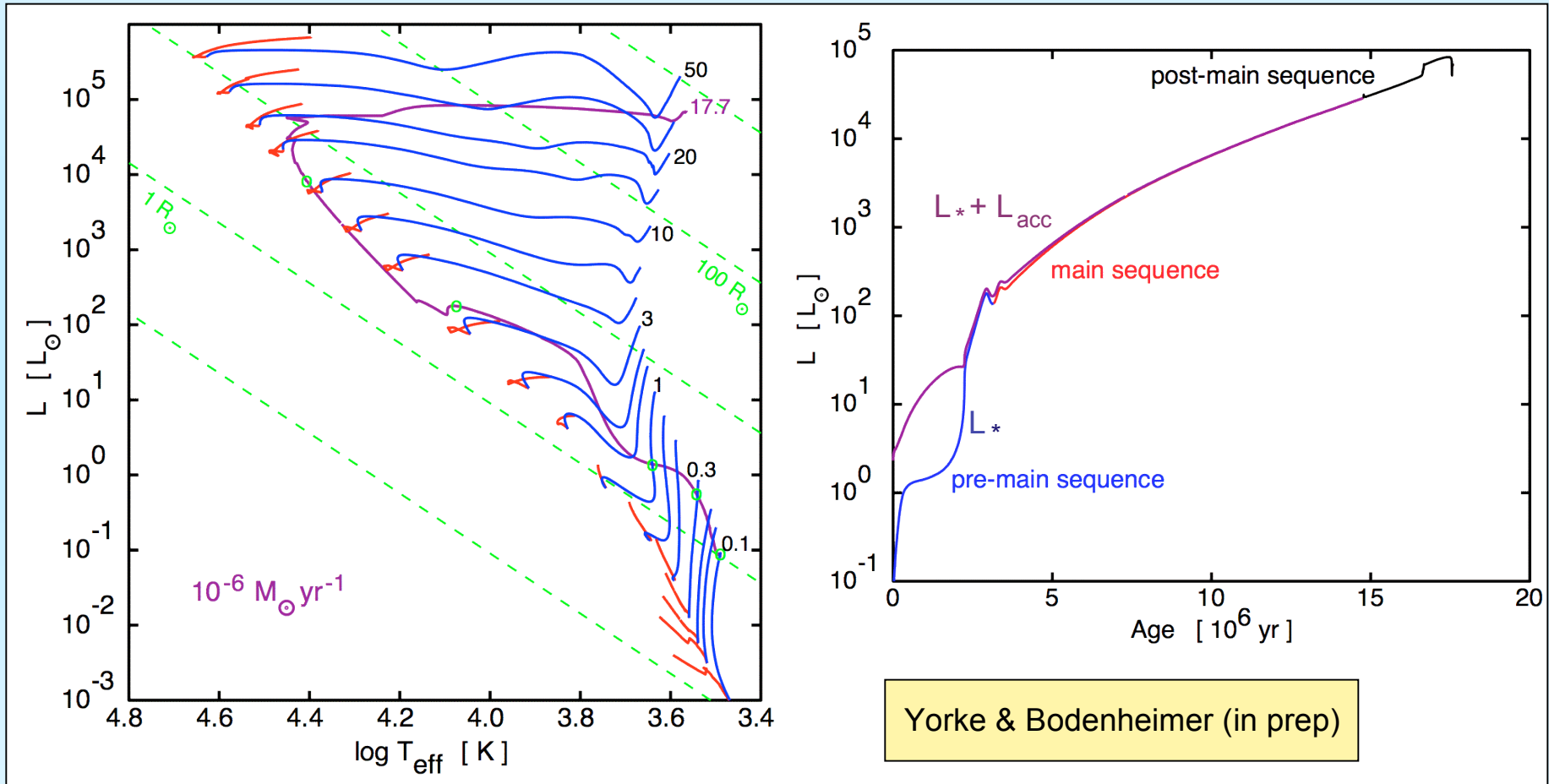
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H.W. Yorke 36

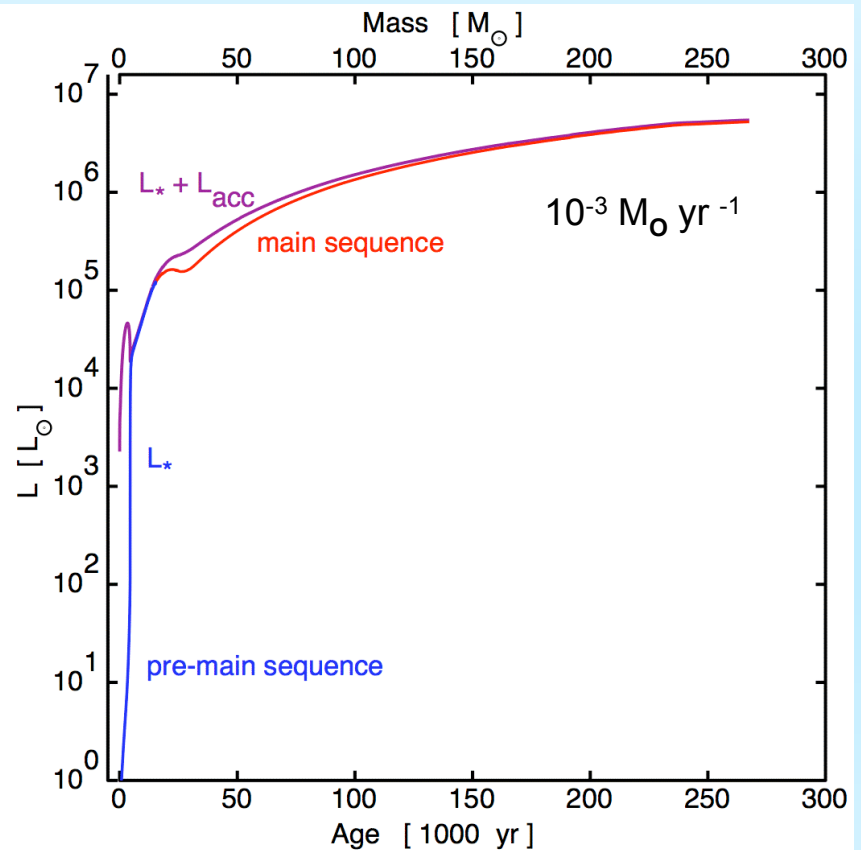
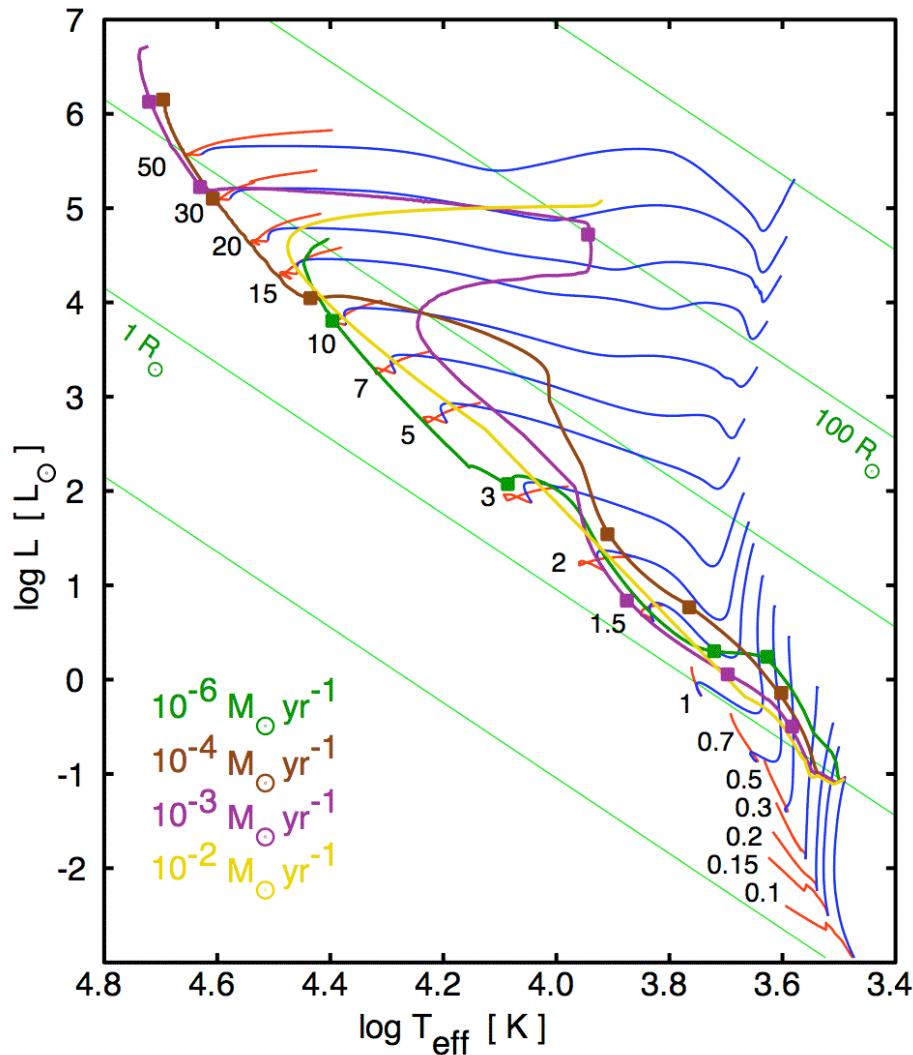


# Evolution of accreting stars in the HRD ( $dM_*/dt > 0$ )





# Evolution of accreting stars in the HRD ( $dM_*/dt > 0$ )

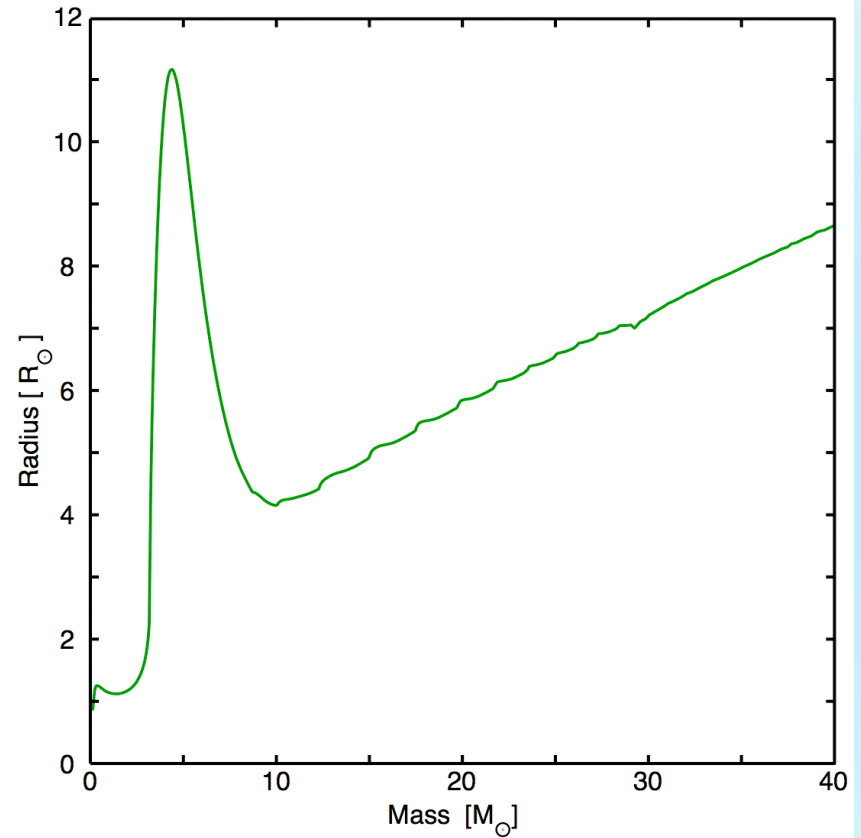
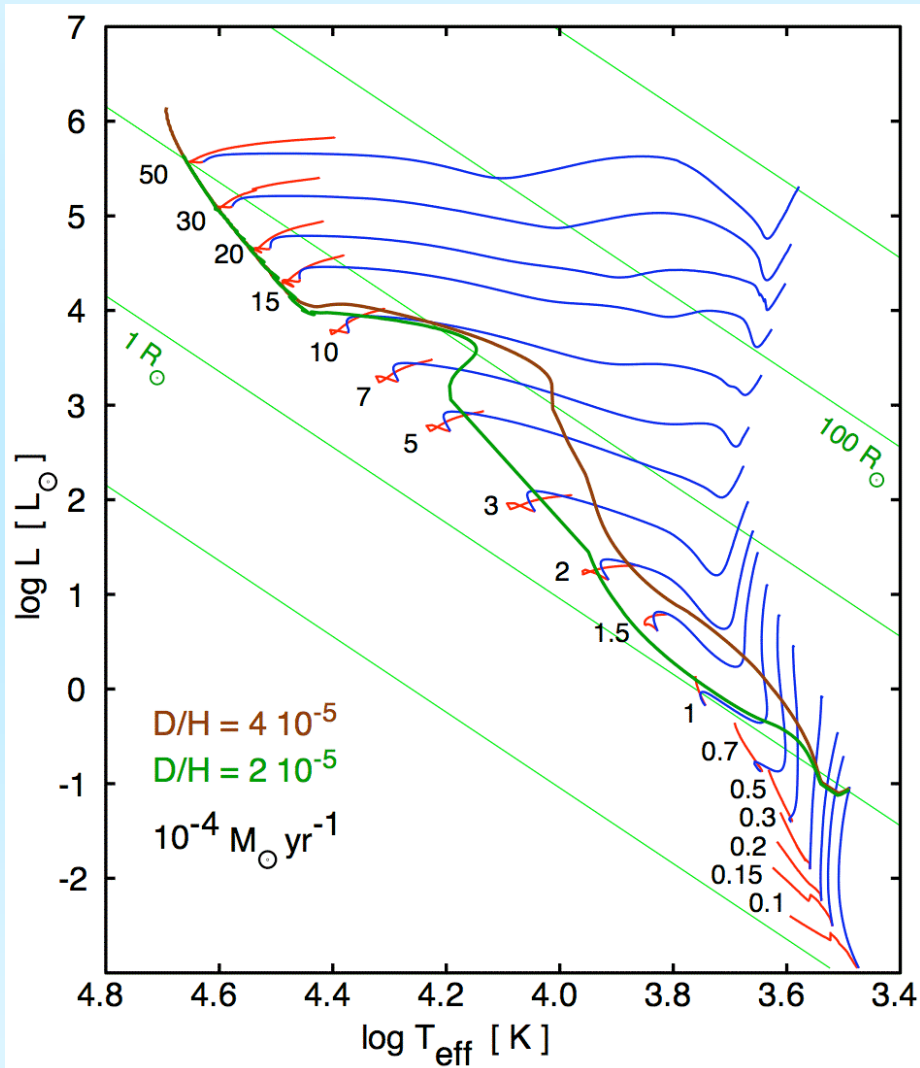


Yorke & Bodenheimer (in prep)





# Evolution of accreting stars in the HRD ( $dM_*/dt > 0$ )



Yorke & Bodenheimer (in prep)



# How likely is it to observe high mass stars during accretion?



## Assume:

- Galactic star formation rate:  $5 M_{\odot}/\text{yr}$
- $t_{\text{acc}} = 2 \times 10^5 \text{ yr}$
- Salpeter IMF
  - $N(M) dM = A M^{-\alpha} dM$
  - $0.1 M_{\odot} < M < 100 M_{\odot}$

Note that the local Galaxy ( $r < 500 \text{ pc}$ ) contains  $\sim 10^{-3}$  of Galactic star forming volume

$\Rightarrow 10^6 M_{\odot}$  in stars

$M_*$ [ $M_{\odot}$ ]	$N_*$ $\alpha = 2.35$	$N_*$ $\alpha = 2.3$
>10	5400	6300
>20	2000	2400
>30	950	1200
>50	390	480



# Concluding Remarks



- Accretion physics will be key to understanding formation of massive stars
- Stellar evolution is not dead



# Conclusions

- Massive star formation is a difficult theoretical problem
  - Magnetic, radiative forces on dust important
  - Complex microphysics (dust, degree of ionization, ...)
  - Massive stars form in groups/clusters and their winds, ionizing radiation, and supernova explosions strongly affect ongoing star formation.
- Massive stars photoionize nearby disks (including their own), eventually destroying the disks.
  - The accretion and disk destruction processes close to massive stars operate on similar timescales.
  - Photoionization can limit the final mass of the star.
- Observations continue to confound theorists



Thank you