



From the Formation of Turbulent Molecular Clouds to Dense Cores

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I. Introduction

- High-mass star-forming cores are systematically denser, more massive and more “turbulent” (have higher velocity dispersions) than low-mass star-forming cores (e.g., Garay & Lizano 1999; Kurtz et al. 2000; Beuther et al. 2007).
- Understanding the formation mechanisms of low- and high-mass cores is essential for understanding SF in general.
- This talk:
 - Recall ideas of core formation and control of SF efficiency (SFE) by turbulence in molecular clouds.
 - Revisit these ideas in light of preliminary results on properties of high-density regions in simulations of MC formation and collapse out of the diffuse WNM.

II. Turbulent regulation of the SFE in isolated clouds.

- Turbulence is a *multiscale* phenomenon, with largest velocities and timescales at largest spatial scales (Kolmogorov 1941; Larson 1981; Heyer & Brunt 2004).
- Dual role of supersonic turbulence:
 - Prevent monolithic cloud collapse.
 - Promote *nonlinear* (large amplitude) small-scale density fluctuations that
 - Shorter formation and free-fall times than parent cloud's.
 - Involve only a fraction of the total cloud mass (a different kind of filter than AD-mediated cores).
 - Only a fraction of which proceeds to collapse (Elmegreen 1993; Padoan 1995; Vázquez-Semadeni et al. 1996, 2003, 2005; Klessen, Heitsch & Mac Low 2000; Heitsch, Mac Low & Klessen 2001; Padoan & Nordlund 2002; Nakamura & Li 2005).

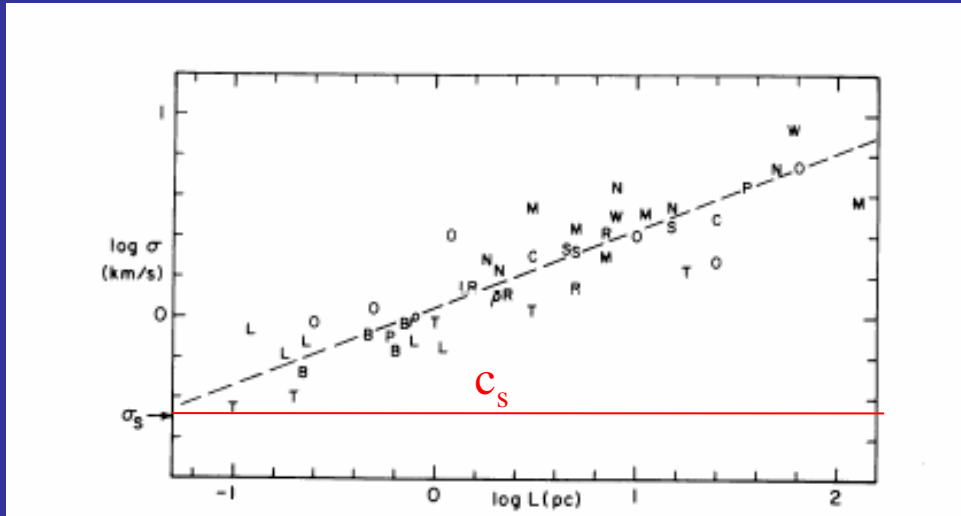
- A model for the inhibitory effect of turbulence in stationary turbulent regimes (continuously driven), is based on the *sonic scale* λ_s (Padoan 1995; Vázquez-Semadeni et al. 2003; Krumholz & McKee 2005):

λ_s : The scale across which the typical turbulent velocity difference equals the sound speed:

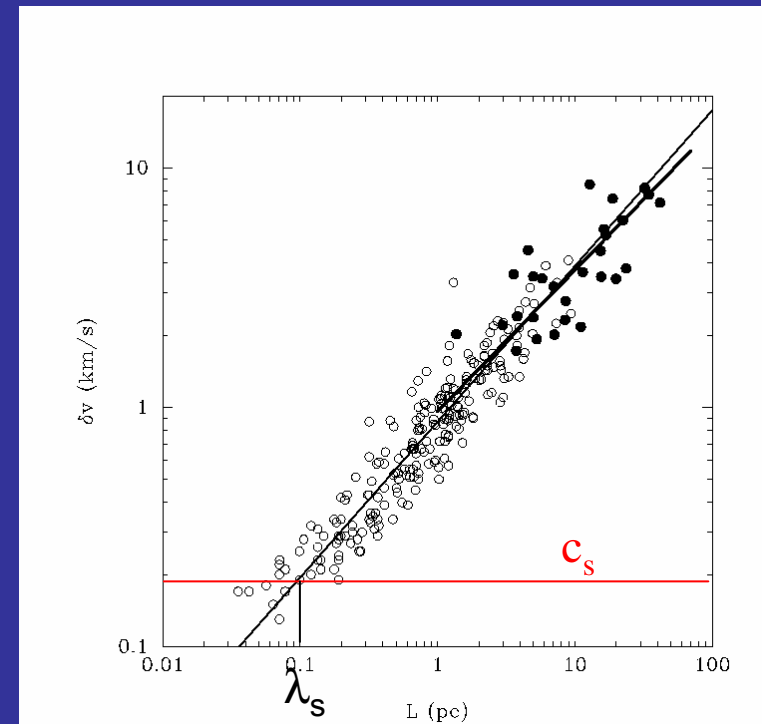
$$\Delta v \approx 0.8 \text{ km s}^{-1} \left(\frac{L}{1 \text{ pc}} \right)^\alpha$$

$$\alpha \approx 1/2$$

$$\lambda_s \sim 0.07 \text{ pc}$$



Larson 1981



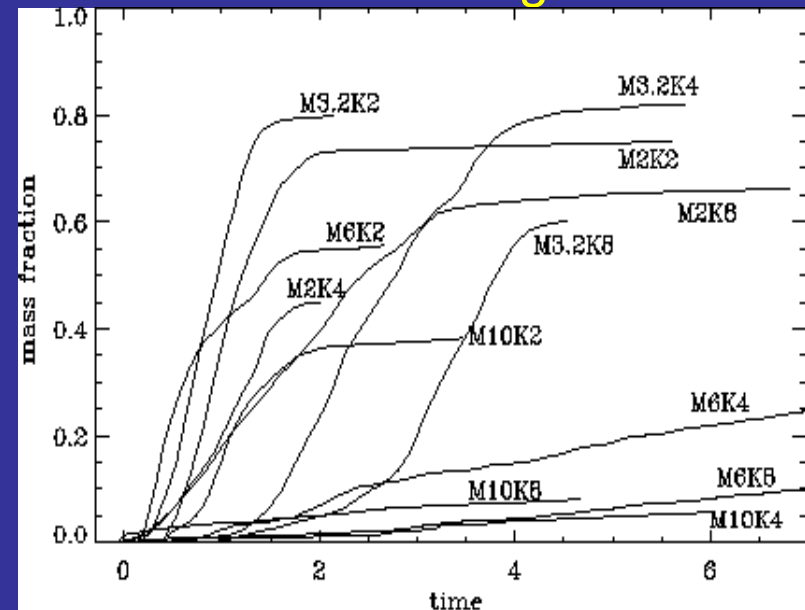
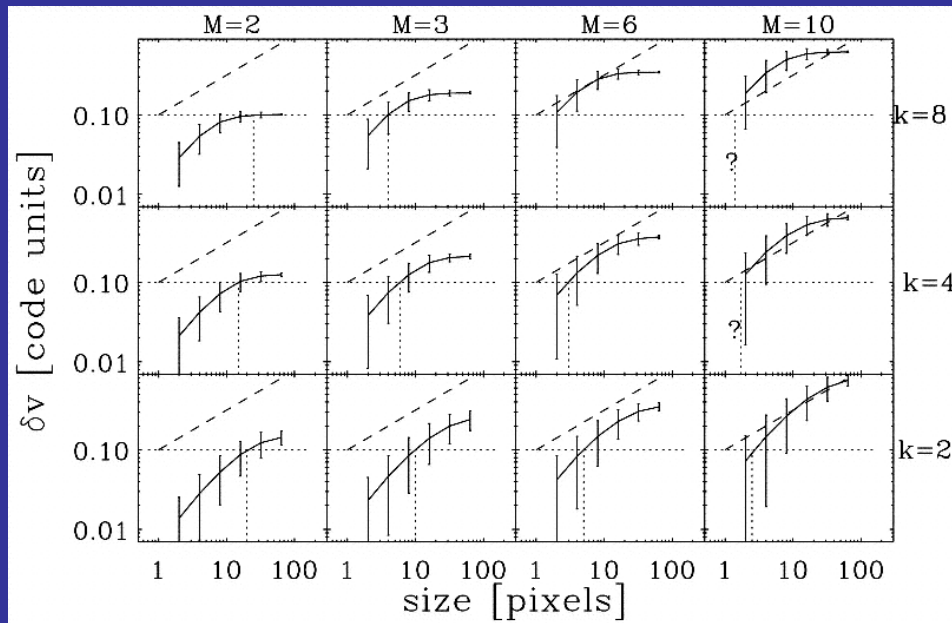
Heyer & Brunt 2004

- Below λ_s :
 - Turbulent subfragmentation becomes weaker ($\delta\rho/\rho \sim M_s^2 < 1$) (or $\sim M_a$ for MHD turbulence – Padoan & Nordlund 2002)
 - Turbulent support becomes subdominant ($\delta u_{\text{turb}} < c_s$).

→ Maybe SFE related to fraction of mass deposited by turbulence in Jeans-unstable cores of size $< \lambda_s$? (i.e., “super-Jeans”, subsonic cores).

- Supported by simulations of varying M_s and driving scale at constant $J=L/L_j=4$ (Vázquez-Semadeni, Ballesteros-Paredes & Klessen 2003, ApJ 585, L131).
 - Sonic scale and SFE measured in the simulations:

SFE measured as collapsed mass fraction after 2 crossing times.



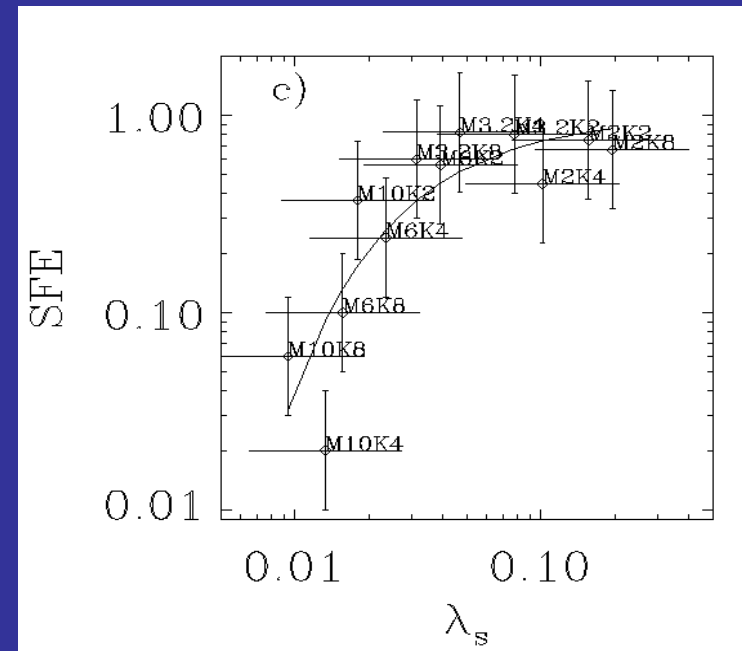
λ_s decreases with increasing levels of turbulence at given T .

M_s	2	3.2	6	10
$2\tau_c/\tau_0$	12.5	10	6.7	2

$$\tau_0 \sim 2/3 \tau_{ff}$$

SFE depends monotonically on λ_s
(regardless of driving length)

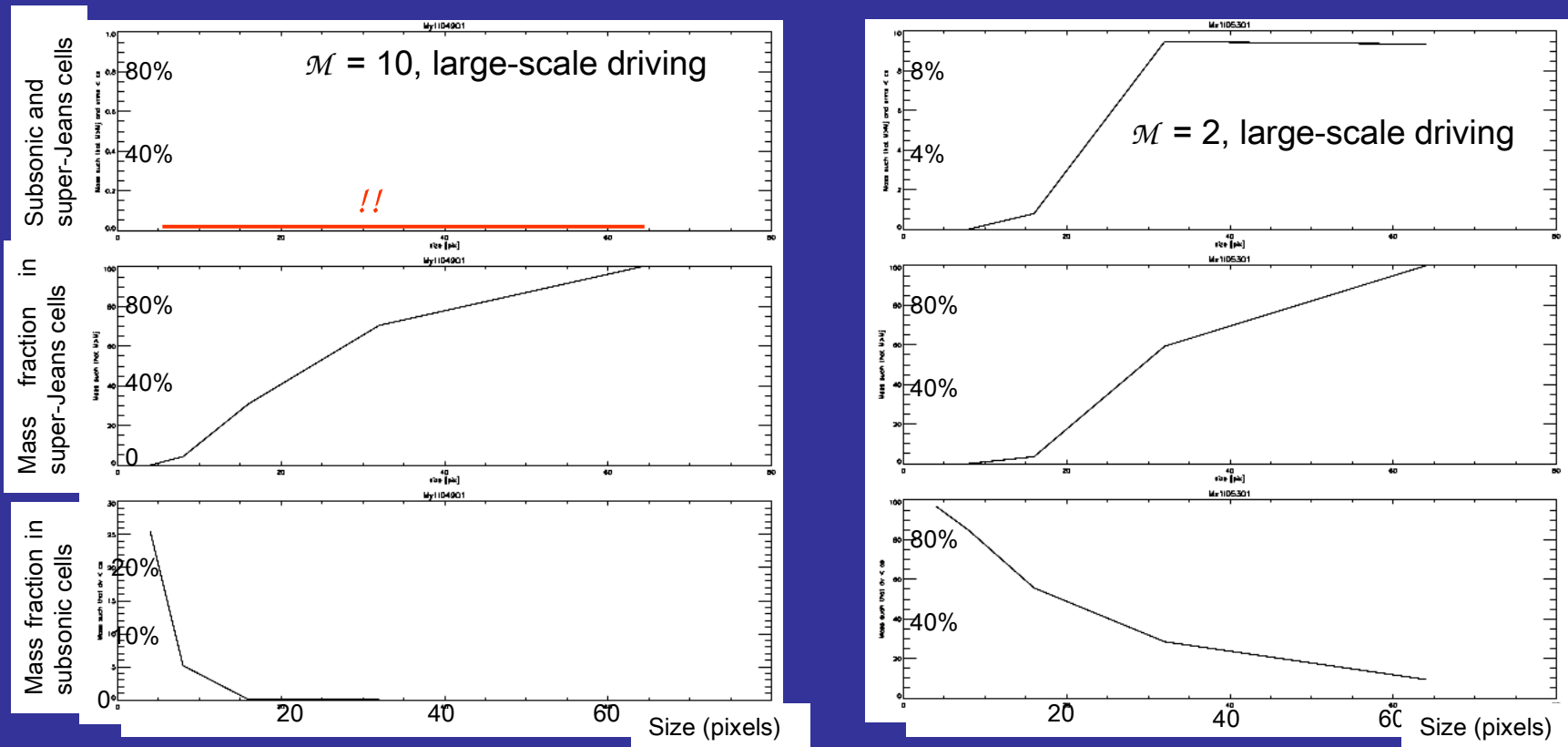
$$\text{SFE}(\lambda_s) \propto \exp(-\lambda_0 / \lambda_s)$$



Vázquez-Semadeni et al. 2003

- The model has been extended by [Krumholz & McKee \(2005\)](#) to use the ratio of λ_s to the Jeans length L_J as the criterion for gravitational collapse.

- Caveat** : Fraction of mass in subsonic, super-Jeans cells as function of cell size may be lower than mass in collapsed objects, even zero at large Mach numbers (Vázquez-Semadeni & Ballesteros-Paredes, in prep.).



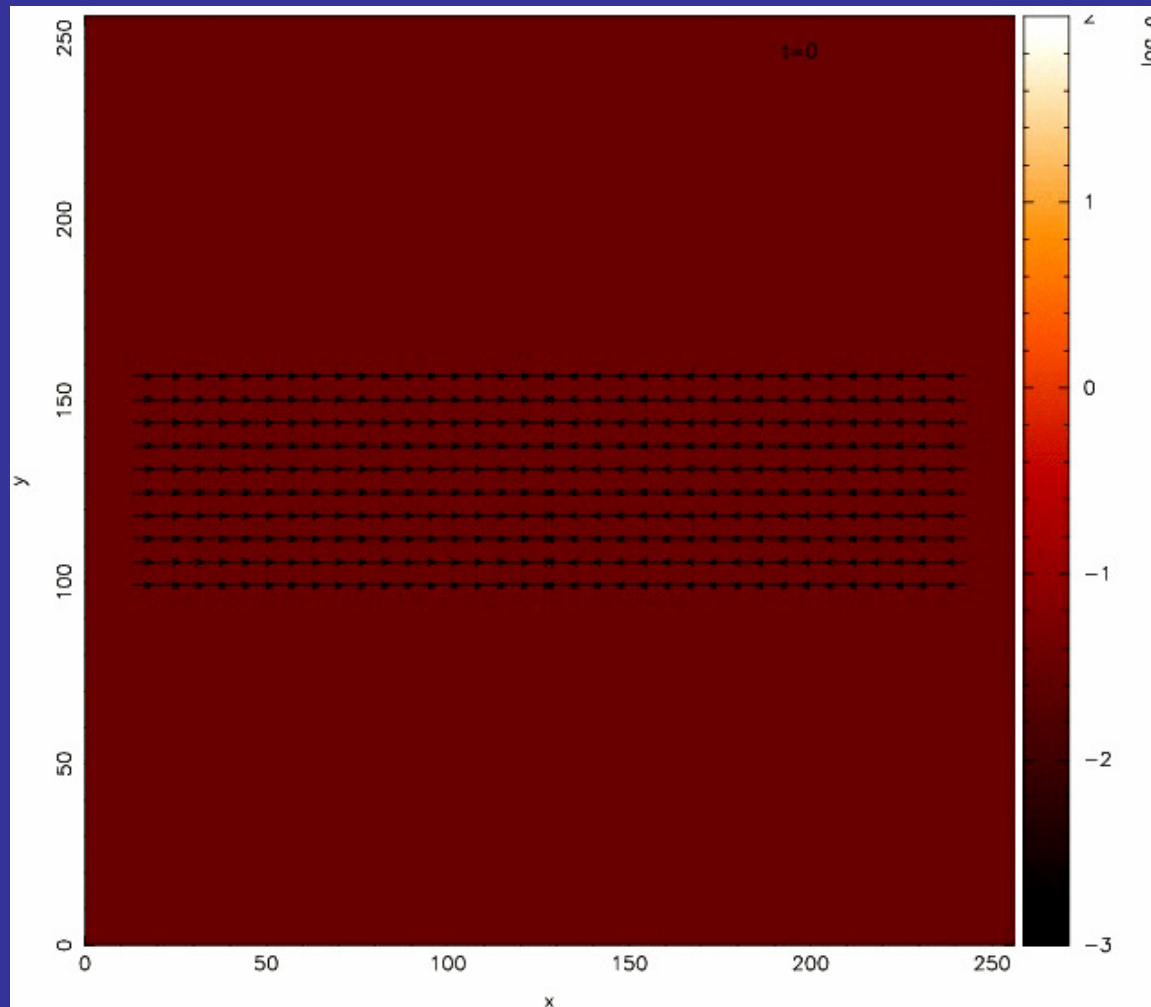
Subcells in simulation, not clumps.

- Conclude:
 - Not all collapsing mass may come from subsonic, super-Jeans structures (Bate, Bonnell et al...)
 - Supersonic regions may also be involved in collapse to form stars.
 - Must *flow* into the collapsed object.
 - What is the nature of the “turbulence” in these regions?
 - Support against collapse (e.g., Matzner, McKee, Krumholz, Tan, Li & Nakamura), or *driven* by gravity? (Goldreich & Kwan 1974; Burkert & Hartmann 2004; Hartmann & Burkert 2007; Peretto et al. 2007; Field et al. 2006) (“Chicken or egg?”).
 - Need to study the *formation* of the cores in clouds with “natural” turbulence (Heitsch’s talk).

→ Cloud formation and evolution studies

III. Search for massive-SF-like regions in numerical simulations of MC formation (Vázquez-Semadeni, Ballesteros-Paredes, Jappsen, Klessen & Price 2007, in prep.)

- *Preliminary and tentative results!*
- Use simulations of MC formation by transonic compressions in diffuse WNM (Vázquez-Semadeni et al. 2007, ApJ 657, 870).

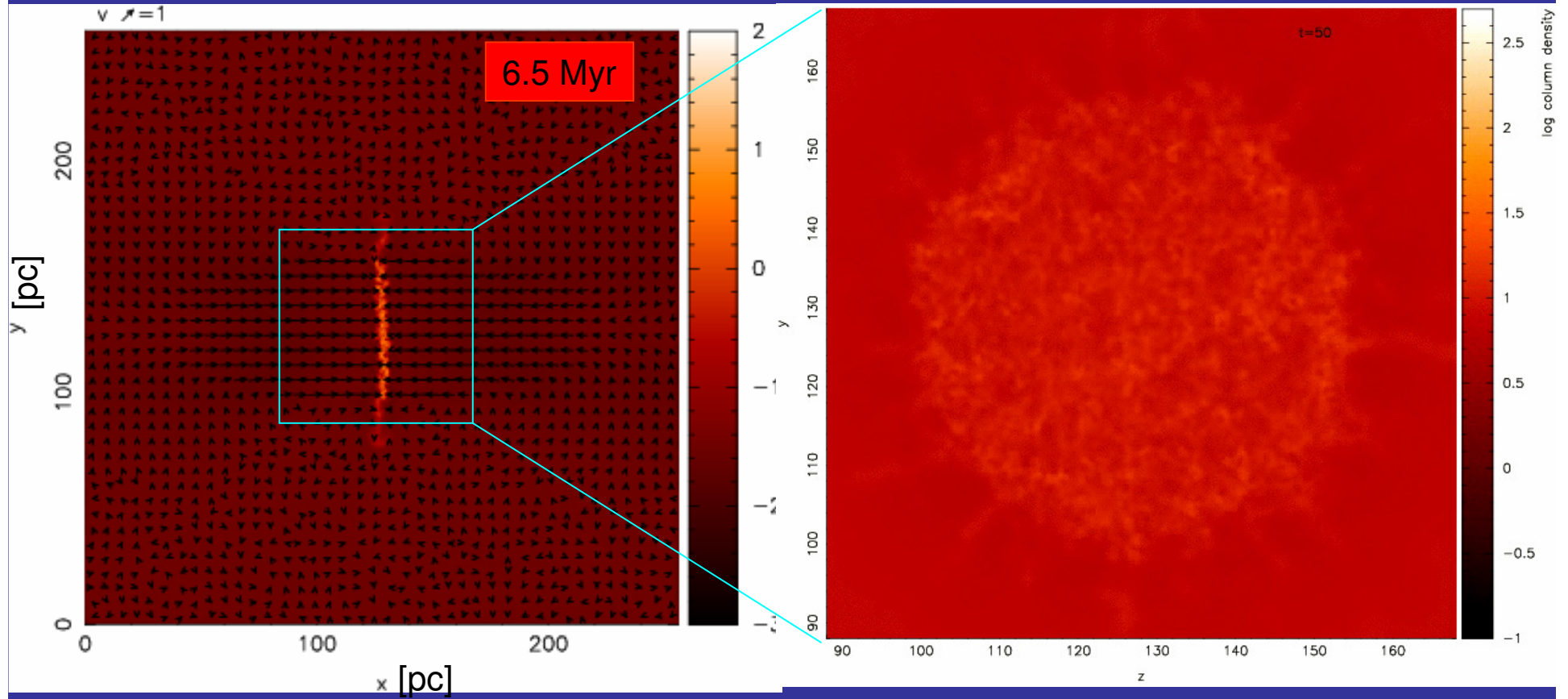


SPH simulation includes cooling (leading to TI) and self-gravity.

$L = 256 \text{ pc}$
 $Dt = 39 \text{ Myr}$
 $\langle n \rangle = 1 \text{ cm}^{-3}$
 $v_{\text{inf}} = 9.2 \text{ km s}^{-1}$
 $T_{\text{ini}} = 5000 \text{ K}$

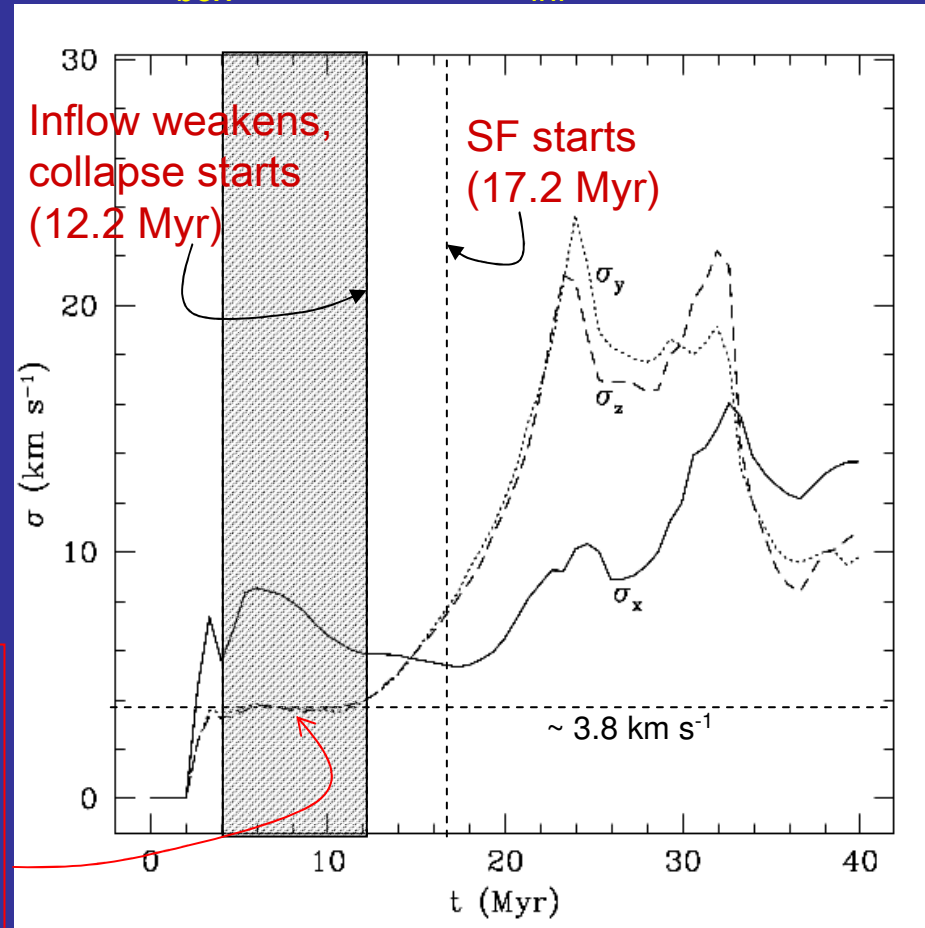
Cloud formation and turbulence generation proceed by TI, KHI, and NTSI as described by Fabian Heitsch. 11

6.5 – 39 Myr



- E_{kin} driven first by inflow, then by gravitational contraction.

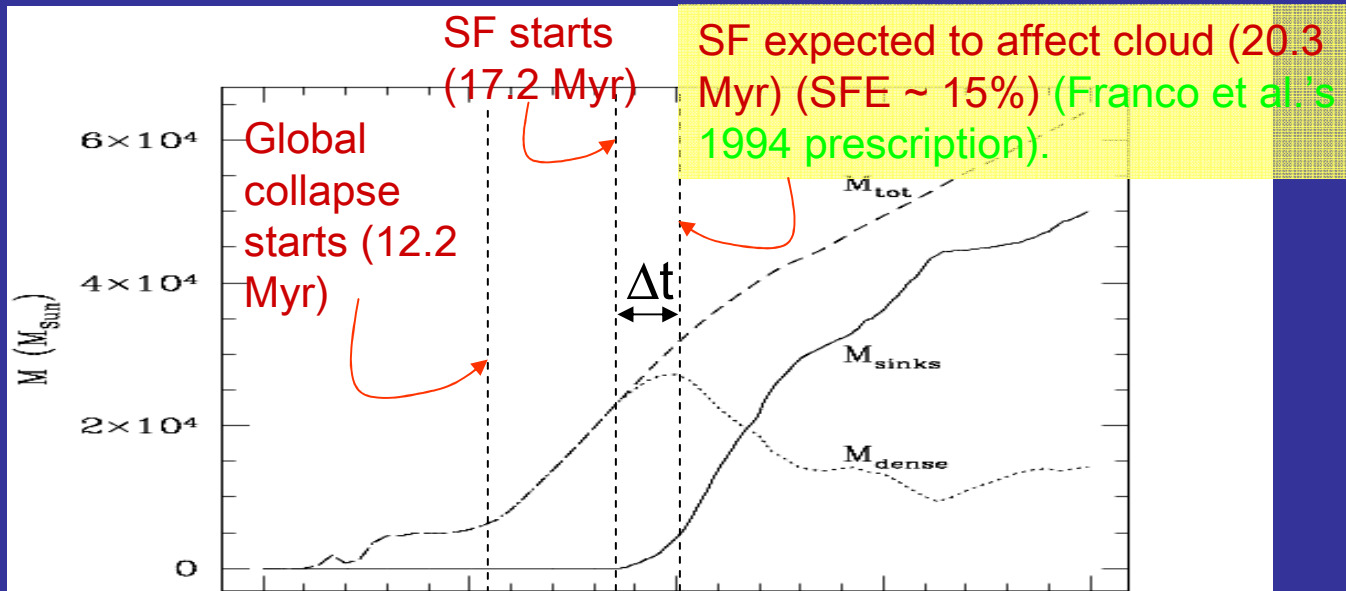
$$L_{\text{box}} = 256 \text{ pc} , L_{\text{inf}} = 112 \text{ pc}$$



Turbulence driven by compression, through NTSI, TI and KHI

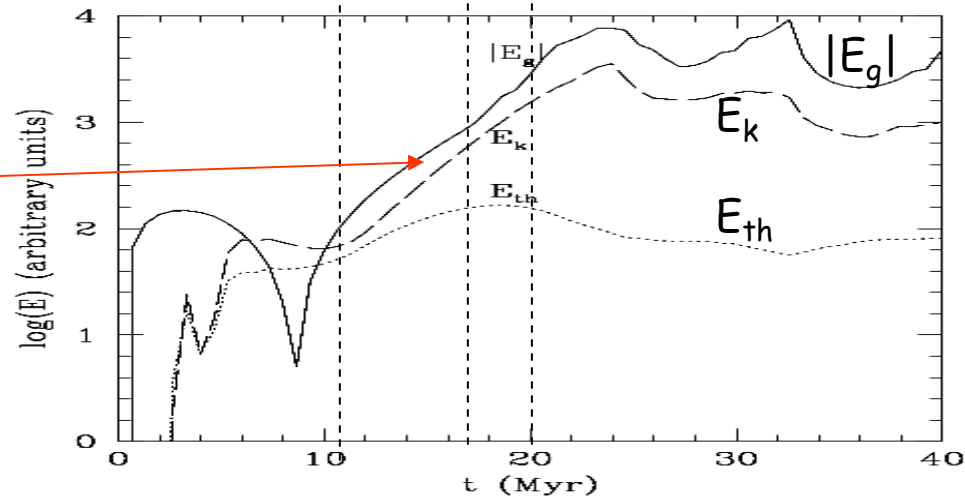
(Walder & Folini 1998;
Koyama & Inutsuka 2002;
Audit & Hennebelle 2005;
Heitsch et al. 2005, 2006;
Vázquez-Semadeni et al 2006)

(Vázquez-Semadeni et al. 2007)



- Turbulent E_{kin} fed by collision first, then by gravitational contraction.

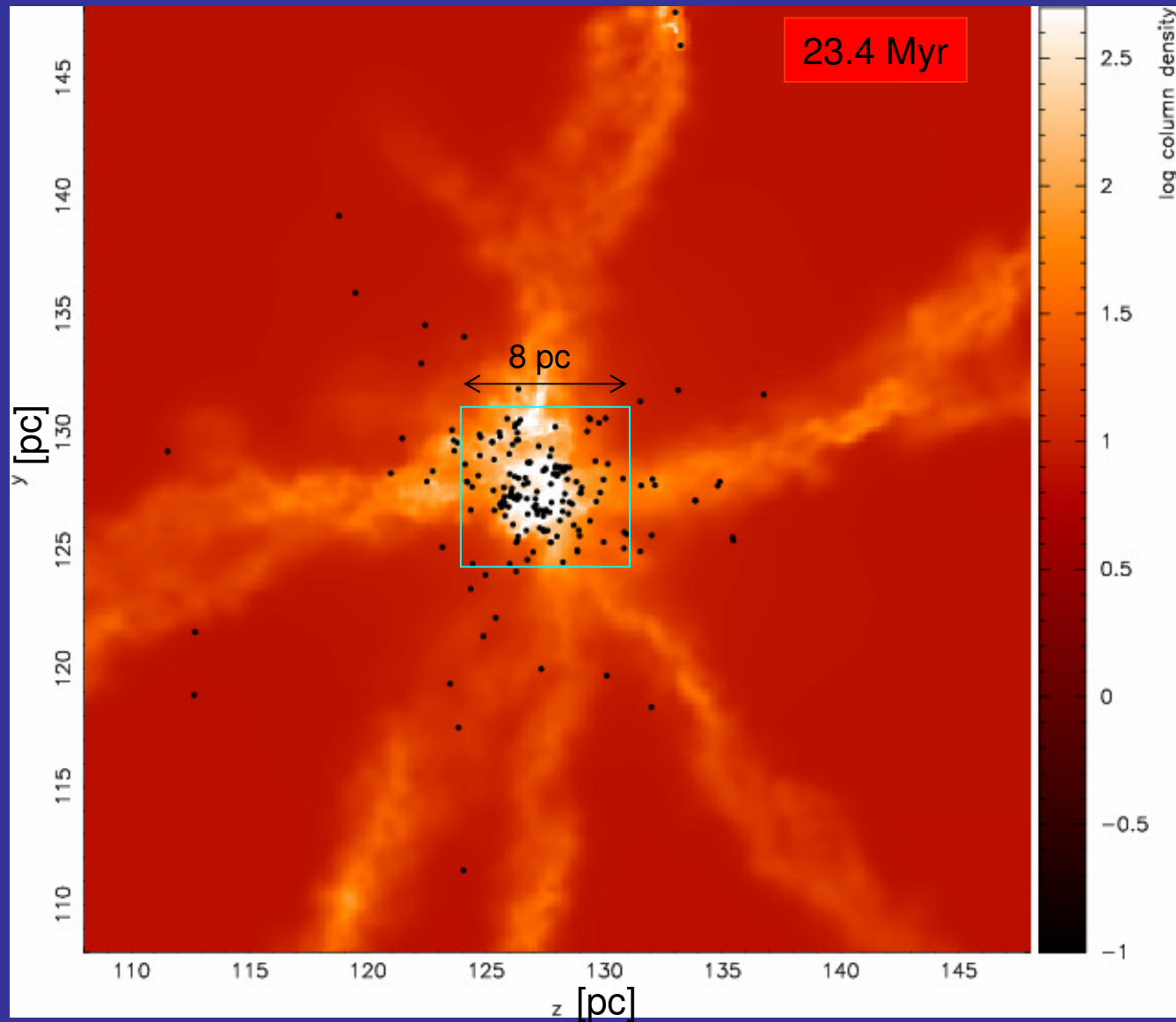
- $|E_g| \sim 2 E_k$



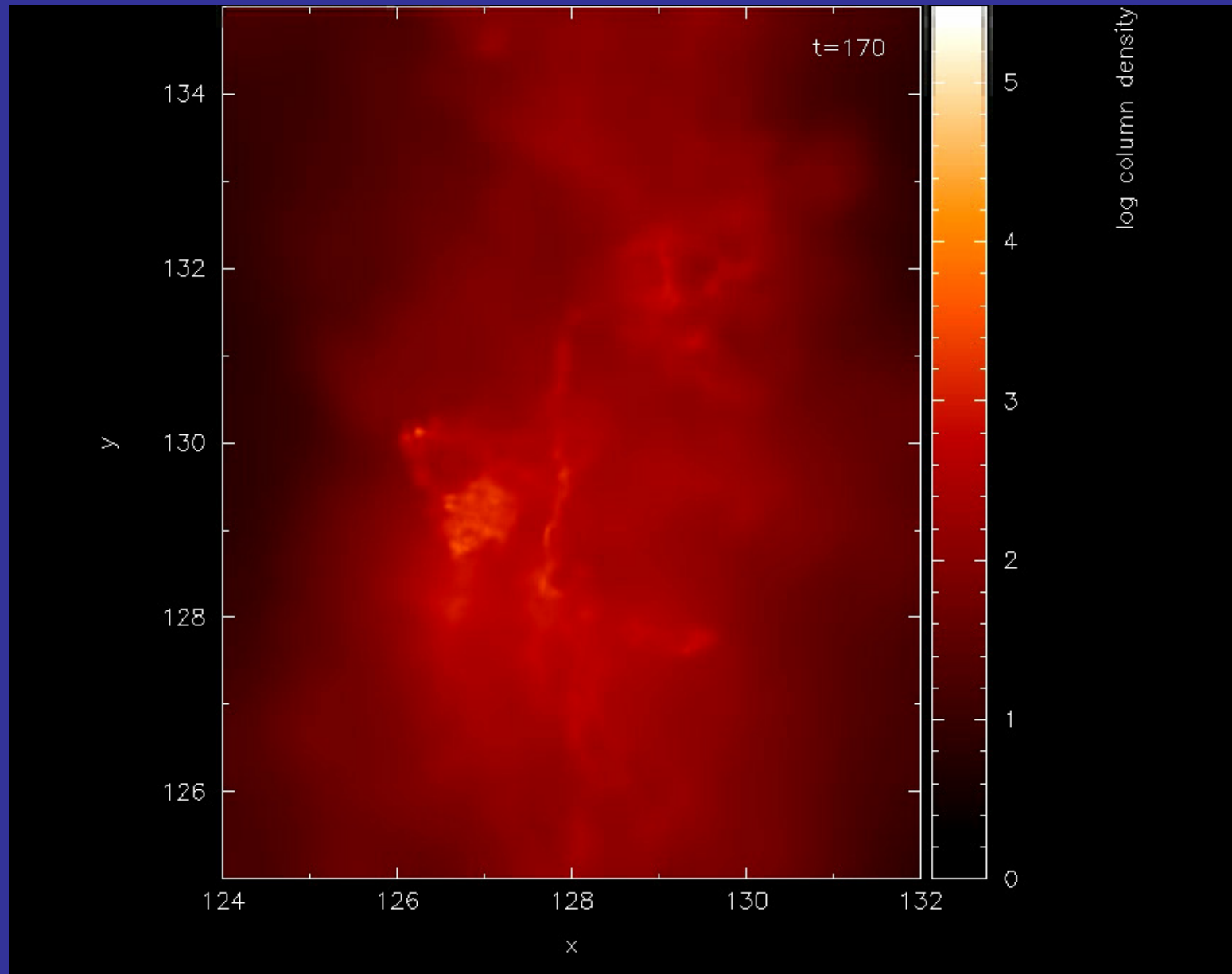
Run with:
 $L_{box} = 256 \text{ pc}$,
 $L_{inf} = 112 \text{ pc}$

(Vázquez-Semadeni et al. 2007)

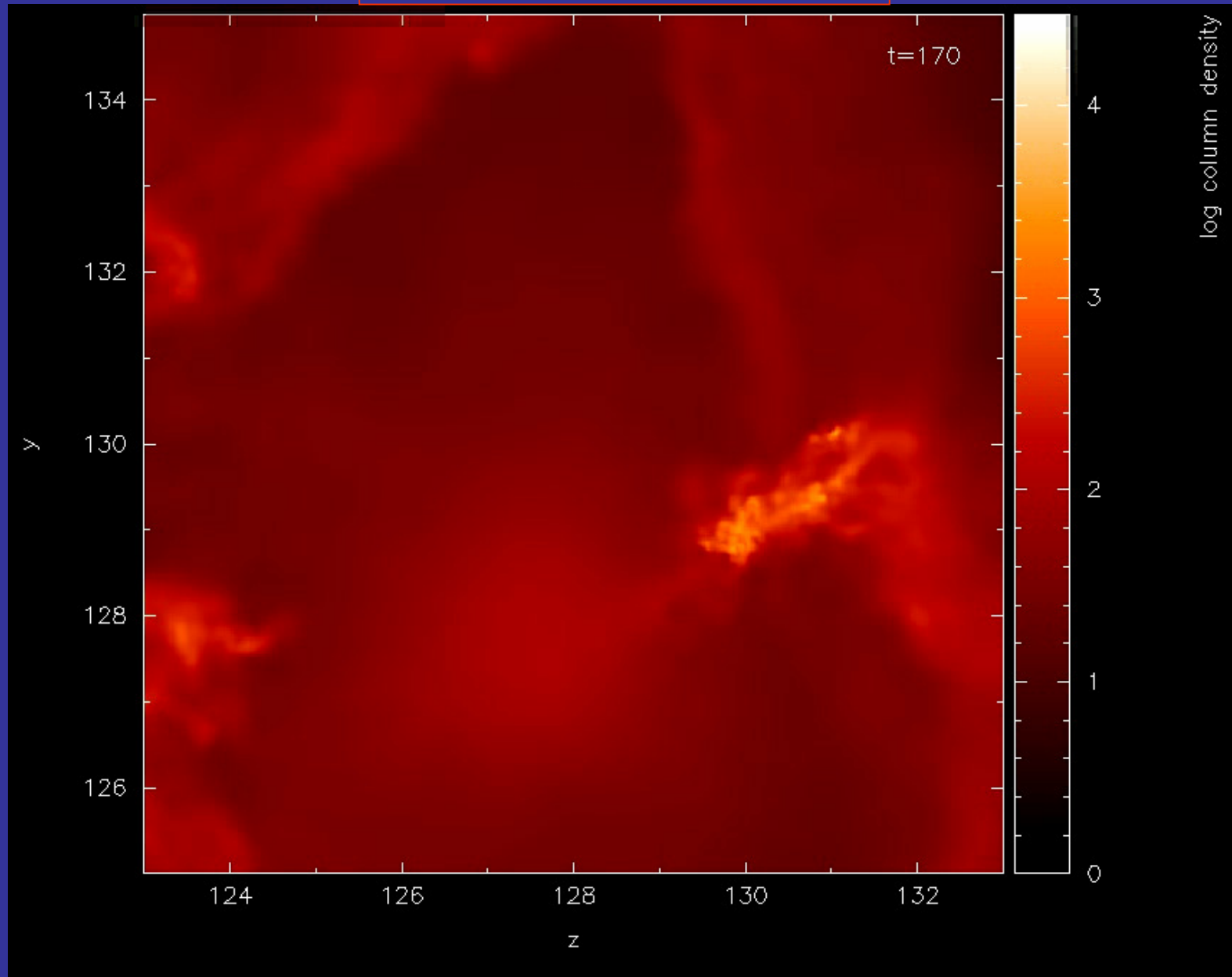
- Focus on time and place of central collision

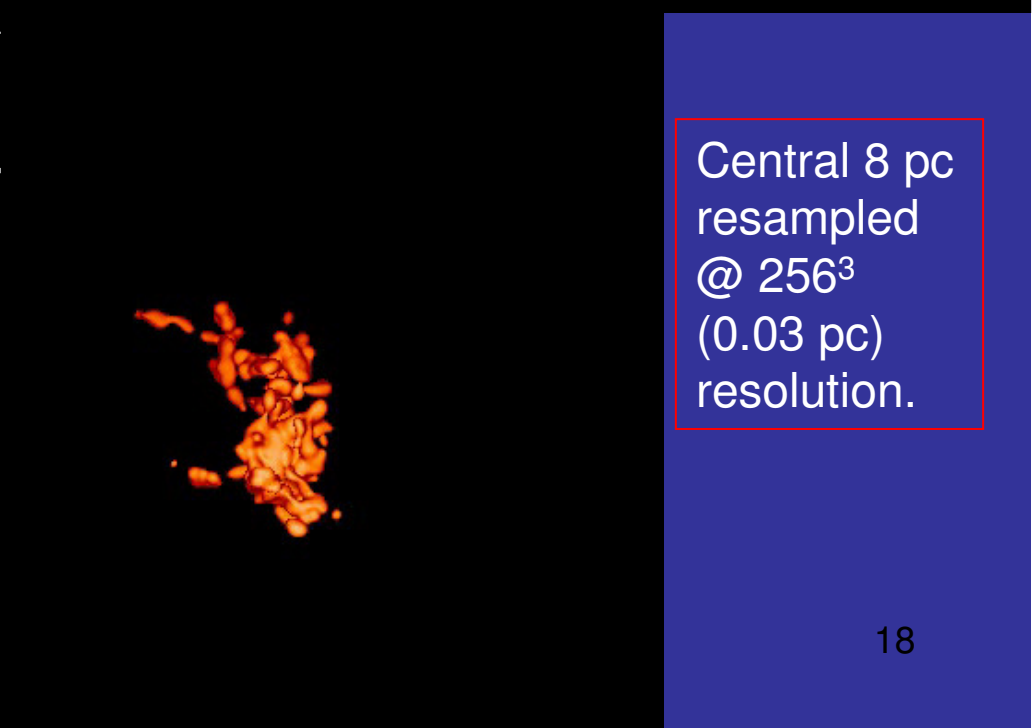
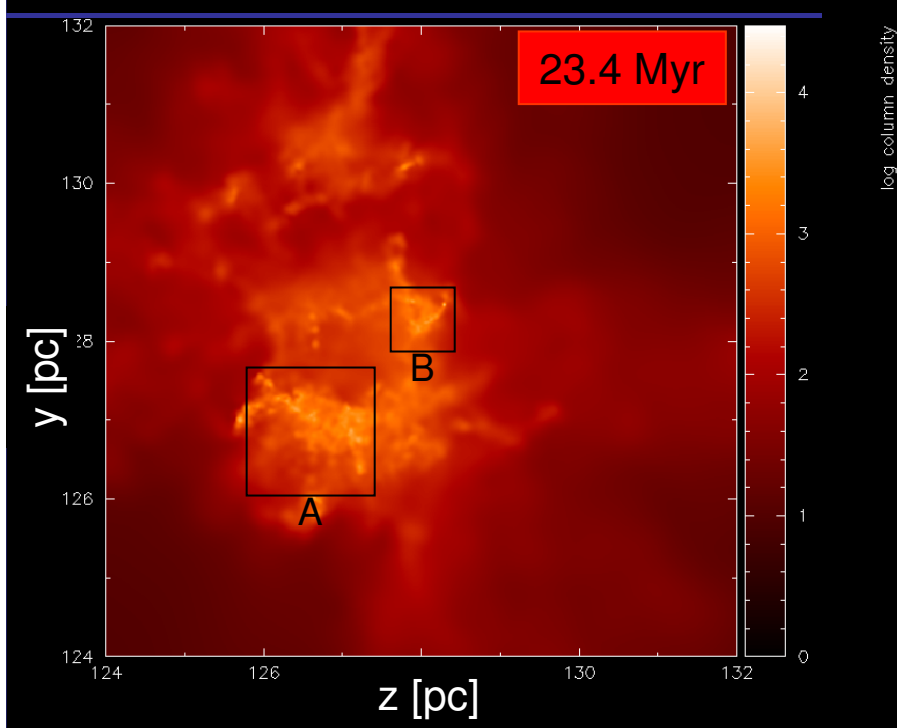
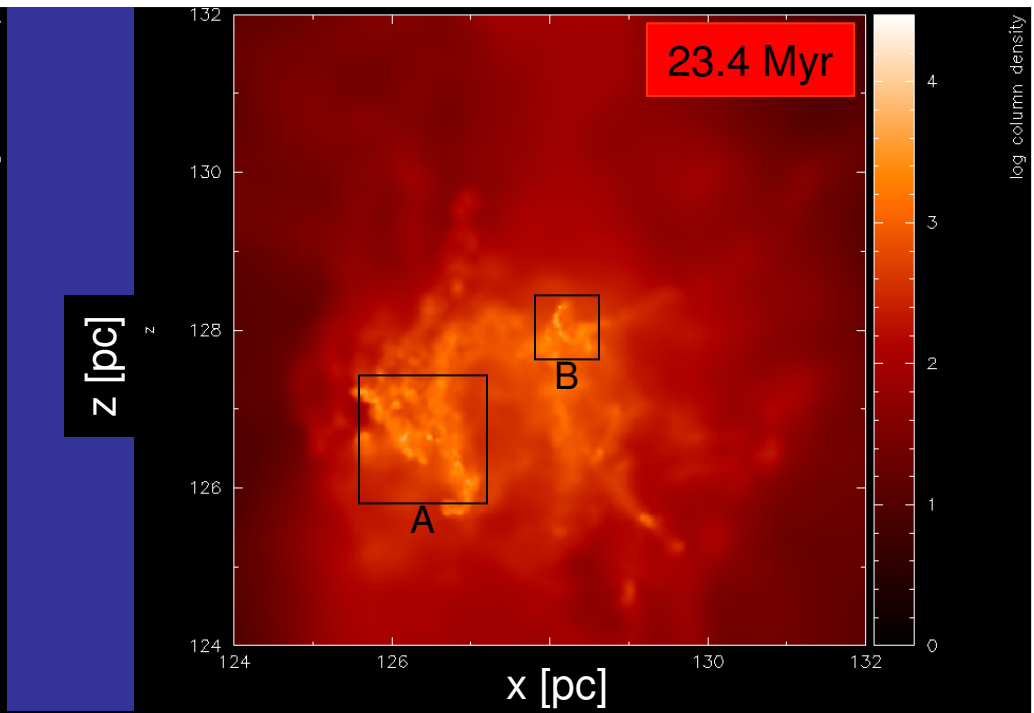
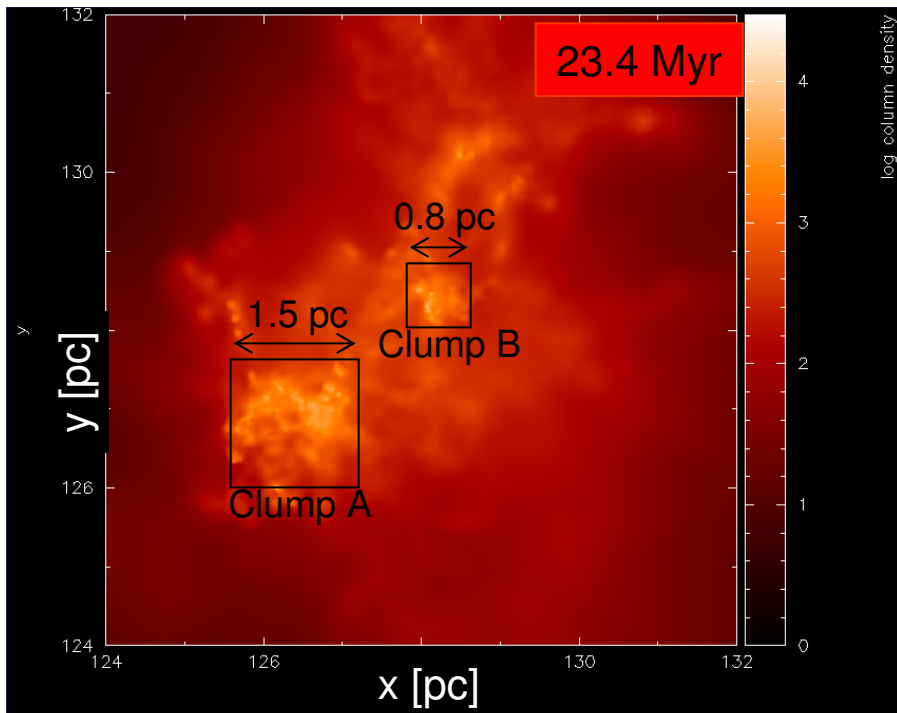


22.1 – 24.7 Myr ($\Delta t = 2.6$ Myr)



22.1 – 24.7 Myr ($\Delta t = 2.6$ Myr)





- Physical properties:

- Whole 8-pc region:

- $\langle n \rangle = 450 \text{ cm}^{-3}$
- $\sigma_{3D} = 5.0 \text{ km s}^{-1}$; $\sigma_x = 2.3 \text{ km s}^{-1}$; $\sigma_y, \sigma_z \sim 3.1 \text{ km s}^{-1}$
- $M \sim 7000 M_{\text{sun}}$

- Clump A ($L = 1.5 \text{ pc}$):

- $\langle n \rangle = 1.27 \times 10^4 \text{ cm}^{-3}$
- $\sigma_{3D} = 3.6 \text{ km s}^{-1}$
- $M \sim 1400 M_{\text{sun}}$

- Clump B ($L = 0.8 \text{ pc}$):

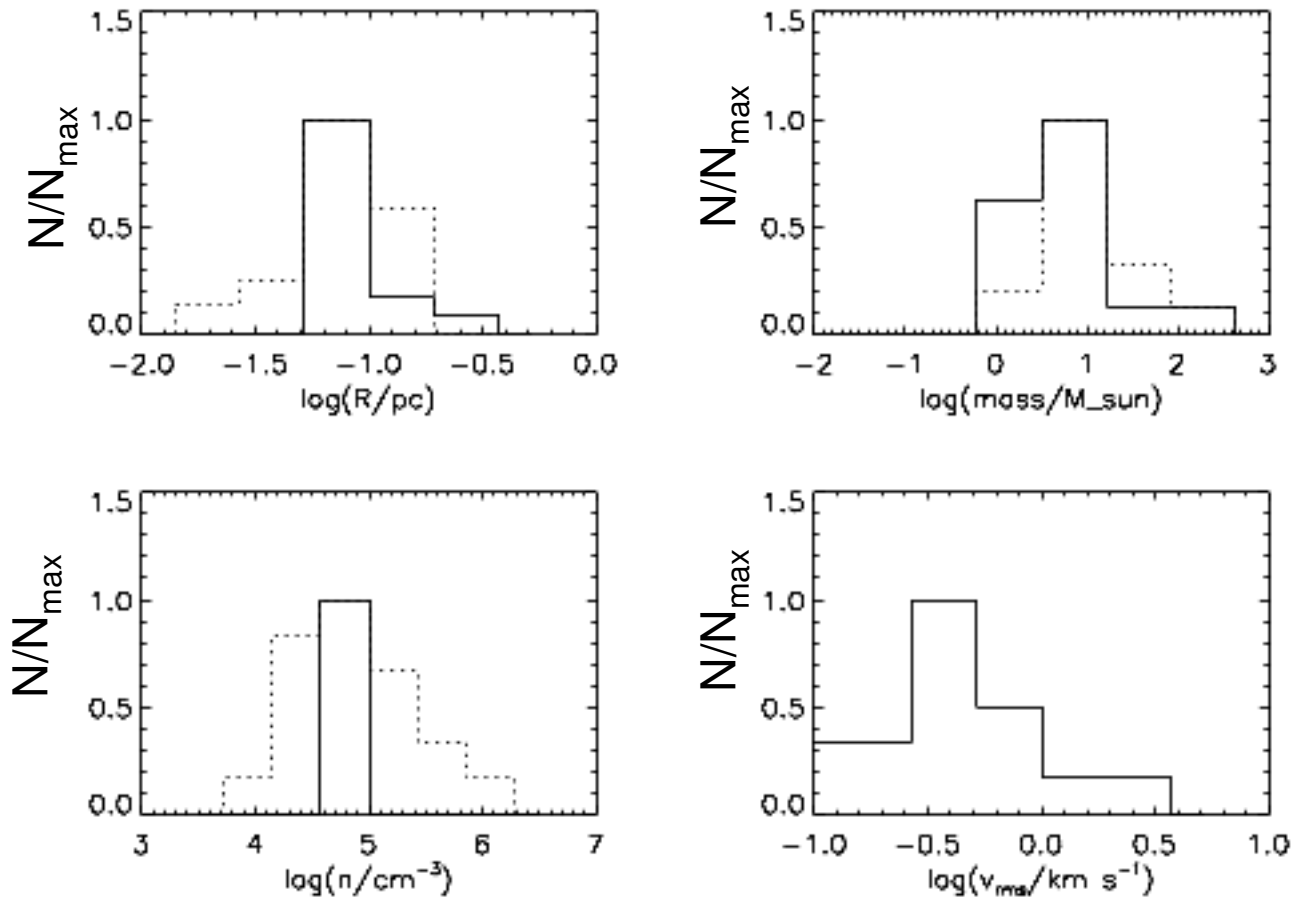
- $\langle n \rangle = 1.72 \times 10^4 \text{ cm}^{-3}$
- $\sigma_{3D} = 2.8 \text{ km s}^{-1}$
- $M = 300 M_{\text{sun}}$

“Typical” Motte et al. (2007) clump:
 $L \sim 0.8 \text{ pc}$
 $n \sim 7000 \text{ cm}^{-3}$

- High-density cores: (simple density threshold criterion, $n > 5 \times 10^4 \text{ cm}^{-3}$, $M > 4 M_{\text{sun}}$).

- Found 15 cores with
 - $n_{\text{max}} \sim 10^{5-6} \text{ cm}^{-3}$.
 - Lifetimes $\ll 1.3 \times 10^5 \text{ yr}$ (appear and disappear in $\ll dt$ between frames). Compare to Motte’s estimate: $\sim 10^3 \text{ yr}$.

- Core statistics:
 - (Zeroth order confrontation with observations.)
- Simulation
⋯ Cygnus X-North (57 cores)
 (Motte et al. 2007 [arXiv:0708.2774]).



Conclude:

The central region of collapse exhibits similar statistical properties to regions of massive SF.

Note: Velocity field has a large infall component, not just random turbulence.

- **Conclusions:**

- Random turbulence provides an effective filter for the mass that can collapse in a MC.

- Super Jeans-, subsonic-fraction model of “mass filtering” for collapse explains low SFE.

- **However:**

- Subsonic, super-Jeans model may possibly miss part of the total mass involved in collapse.
- Numerical simulations of molecular cloud formation with self-gravity (Vázquez-Semadeni et al. 2007) with global cloud contraction (Goldreich & Kwan 1974; Hartmann & Burkert 2007 [Orion]; Andre et al 2007 [Oph]) suggest that
 - Clouds may follow a secular evolutionary path, without equilibrium.
 - » *Appear* virialized, though, due to gravitational contraction.
 - “Turbulence” (*at all scales*) may contain a significant *infall* component.
 - Cores in center of global collapse resemble high-mass SF regions.
 - SFE probably regulated by stellar feedback in this case.
 - » Equilibrate the cloud or disperse it??
 - Work in progress: magnetic fields (Banerjee); stellar feedback (Gómez).

The End