

# A Unified Description of Screened Modified

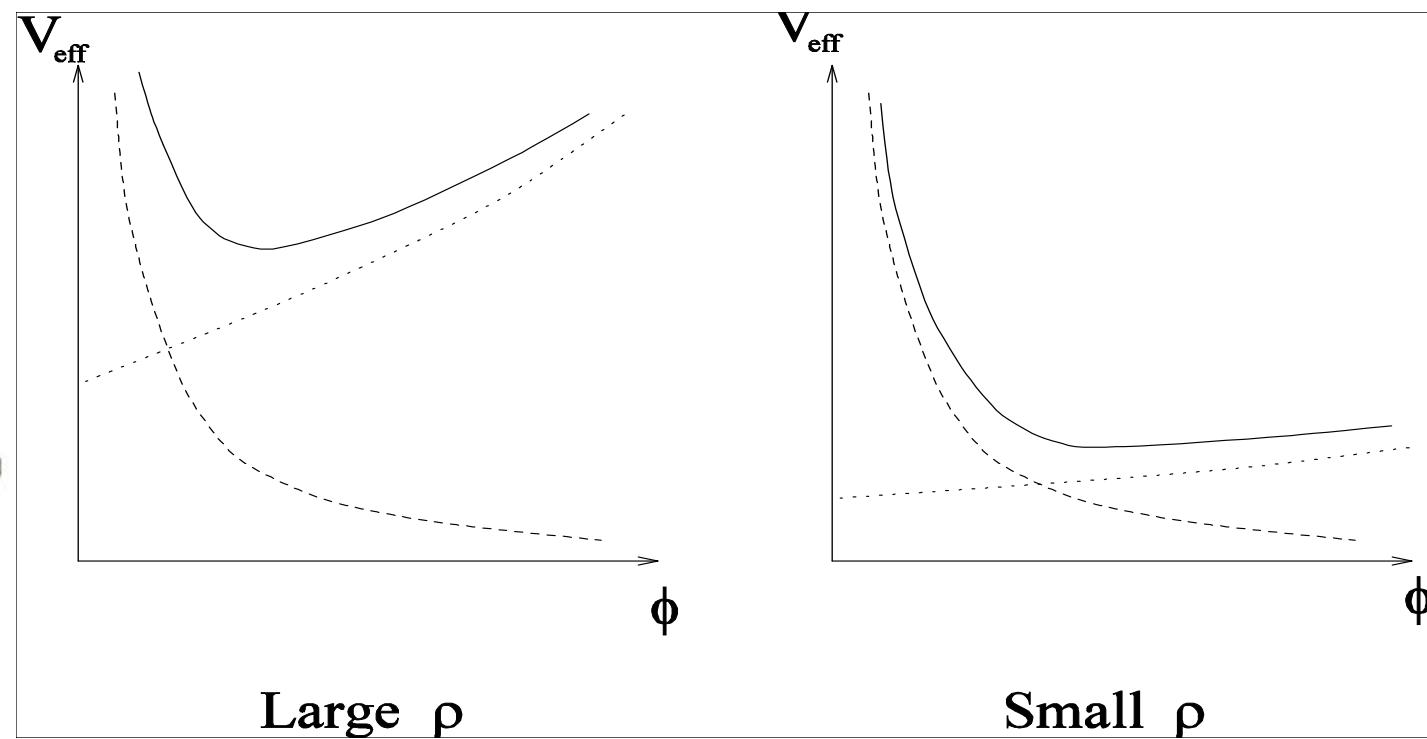
with P. Brax, B. Li, H. Winter and G. Zhao  
1203.4812; 1206.3568

# Outline

- Introduction
- Screened Modified Gravity
- Unified Description - chameleons, symmetrons and dilatons
- Linear and non-linear regime
- N-body plots for symmetron and dilaton

$$S = \int d^4x \sqrt{-g} \left( \frac{1}{16\pi G_N} R - \frac{1}{2} (\partial\phi)^2 - V(\phi) \right) + \mathcal{S}_m(\psi_m, A^2(\phi) g_{\mu\nu})$$

$$V_{eff}(\phi) = V(\phi) + \rho_m A(\phi)$$



# Dilaton

$$V(\phi) = V_0 e^{-\alpha\phi}$$

where the parameters are derived from string theory. The coupling to matter is given by

$$A(\phi) = 1 + \frac{A_2}{2} (\phi - \phi_\star)^2$$

see (Brax et al 1005.3735) for full cosmological behaviour, local constraints and linear perturbation

# Screening

since the field has to be at the min of the effective potential by BBN we deduce

$$\begin{aligned}\phi(a) &= \phi_i + \frac{3}{m_{\text{Pl}}} \int_{a_i}^a da \frac{\beta(a)\rho(a)}{am^2(a)} \\ V(a) &= V_i - \frac{3}{m_{\text{Pl}}} \int_{a_i}^a da \frac{\beta^2(a)\rho^2(a)}{am^2(a)}\end{aligned}$$

the full non-linear dynamics can be constructed from this the mass and coupling parameter as a function of redshift

the screening condition becomes

$$\int_{a_{\text{in}}}^{a_{\text{out}}} \frac{\beta(a)}{am^2(a)} \rho_m(a) da \ll \beta_{\text{out}} m_{\text{Pl}}^2 \Phi_N,$$

we can apply this to all the models

in addition we can study the growth of  
structures in the linear and nonlinear  
regime

# Perturbations

linear regime

$$\Delta_m'' + \frac{a'}{a} \Delta_m' - \frac{1}{2} \frac{\rho_m}{m_{\text{Pl}}^2} a^2 \Delta_m \left[ 1 + \frac{2\beta^2(a)}{1 + \frac{a^2 m^2(a)}{k^2}} \right] = 0,$$

for chameleons

$$m = m_0 a^{-r}, \beta = \beta_0 a^{-s}$$

In general

$r < 3/2$  -- dilaton

$3/2 < r < 3$  -- chameleon

$3 < r$  - large  $R$   $f(R)$

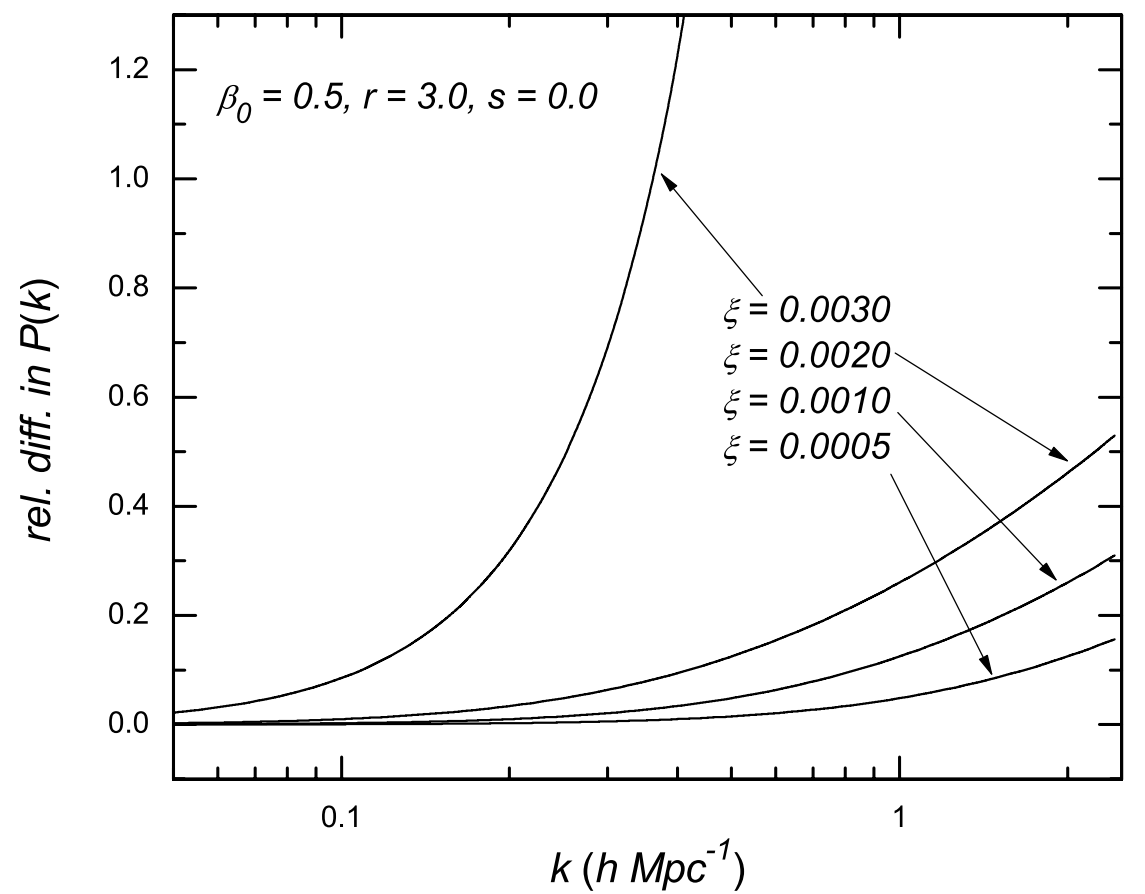
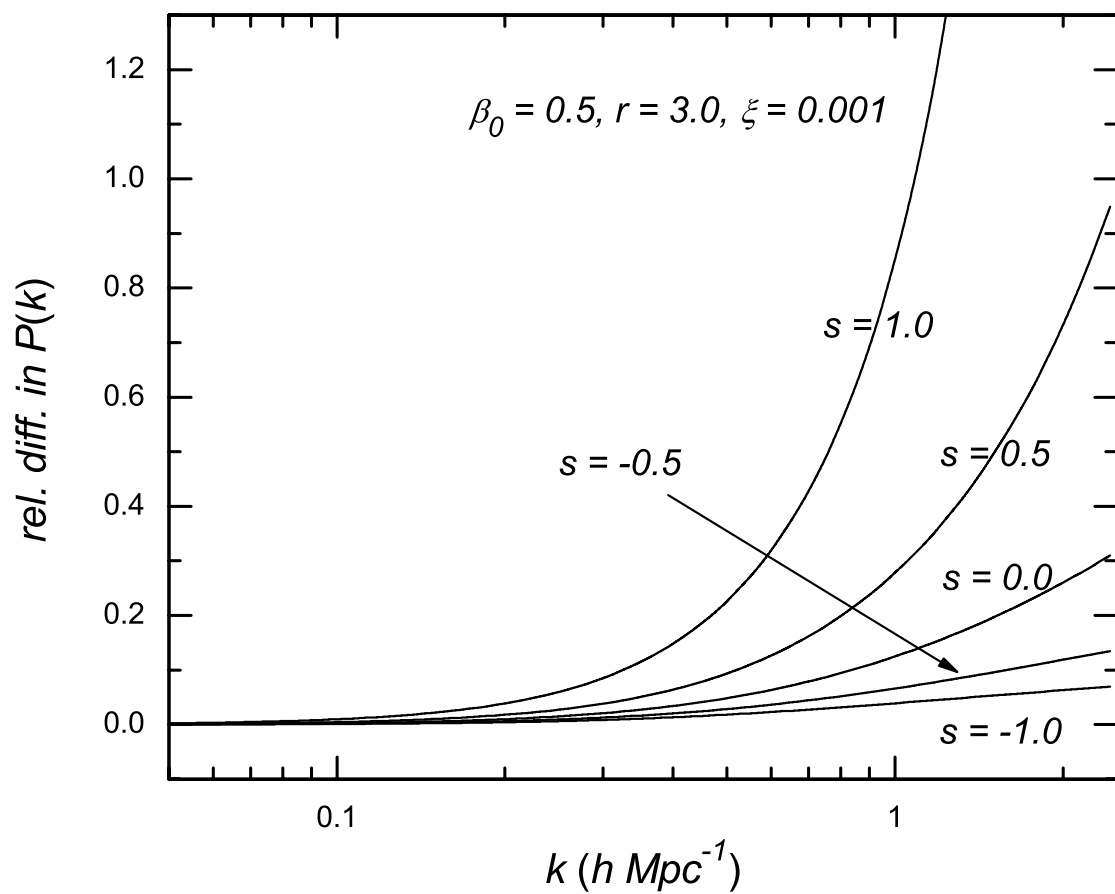
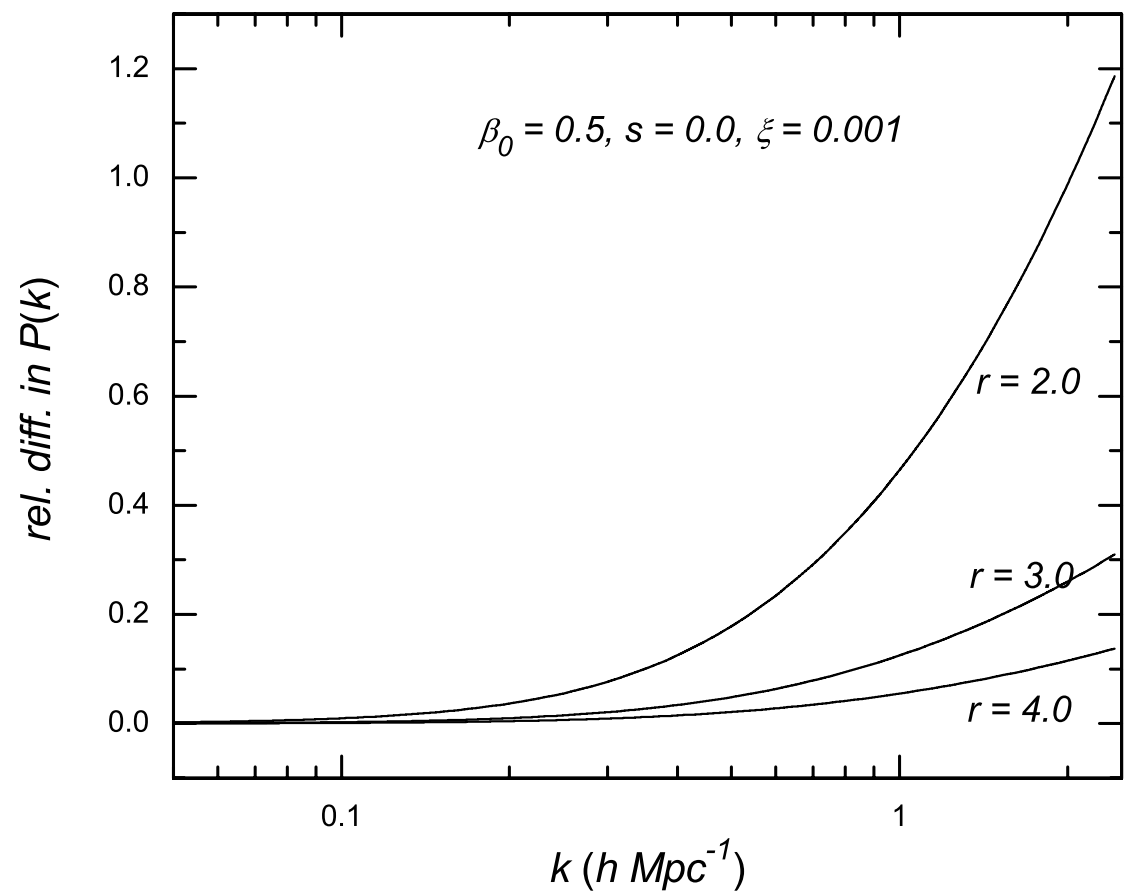
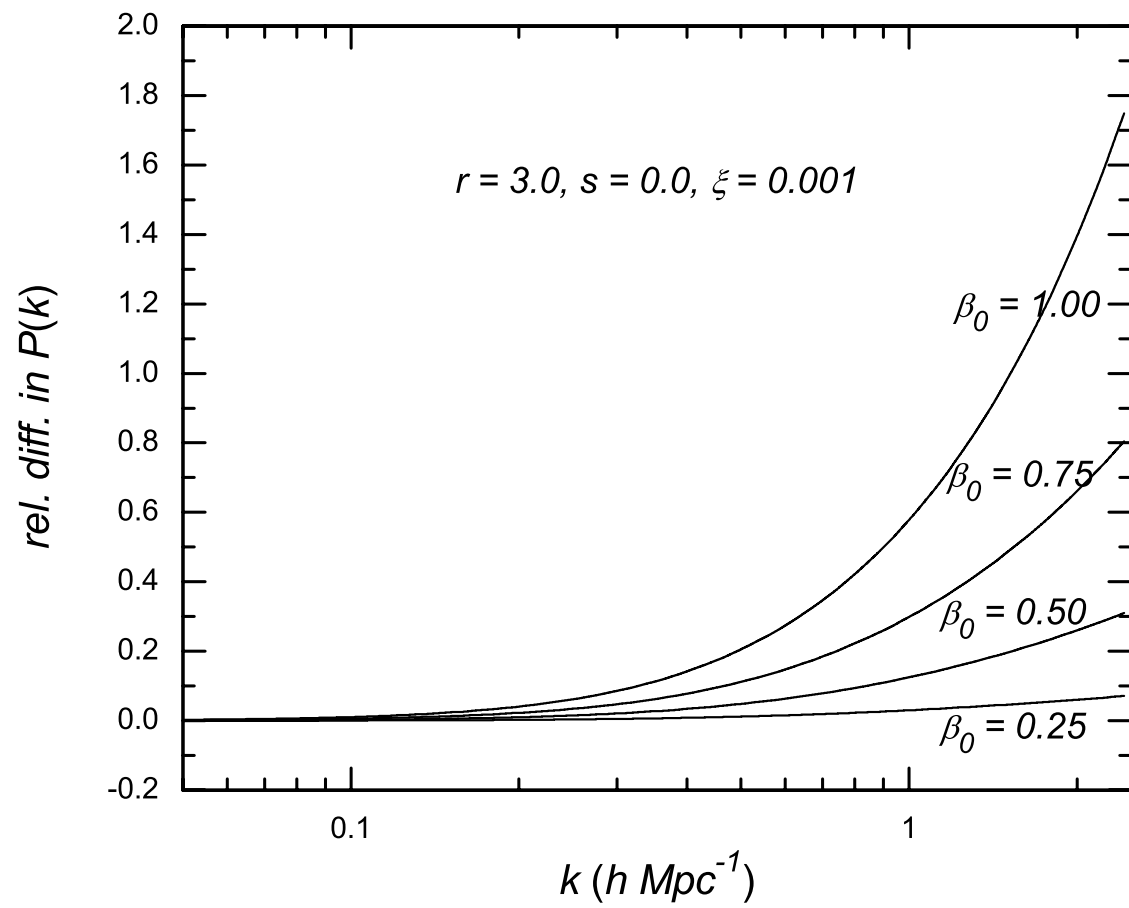
what about the symmetron?

$$m(a) = m_{\star} \sqrt{1 - \left(\frac{a_{\star}}{a}\right)^3}, \quad \beta(a) = \beta_{\star} \sqrt{1 - \left(\frac{a_{\star}}{a}\right)^3},$$

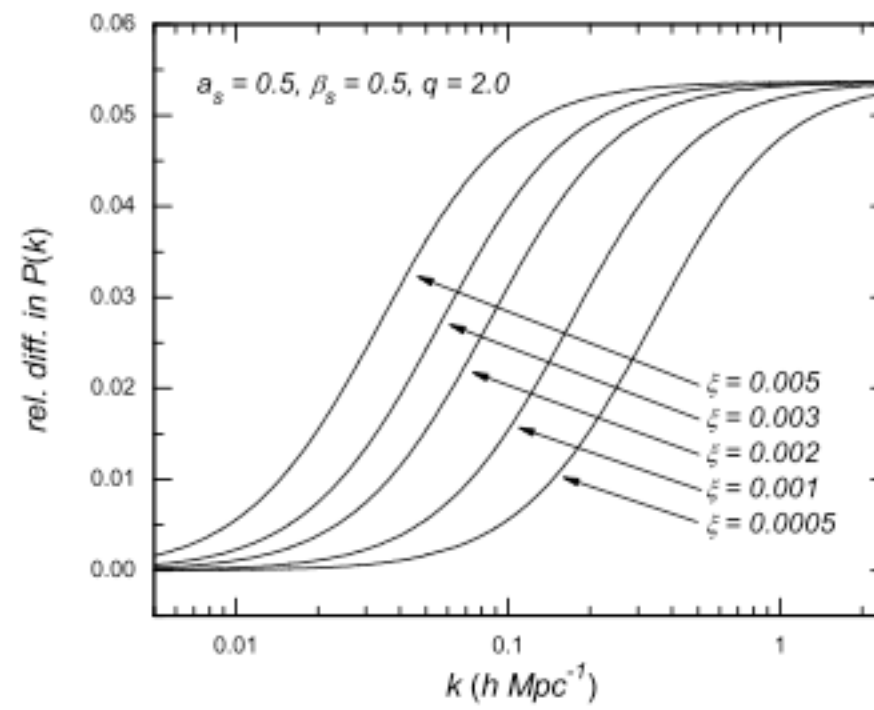
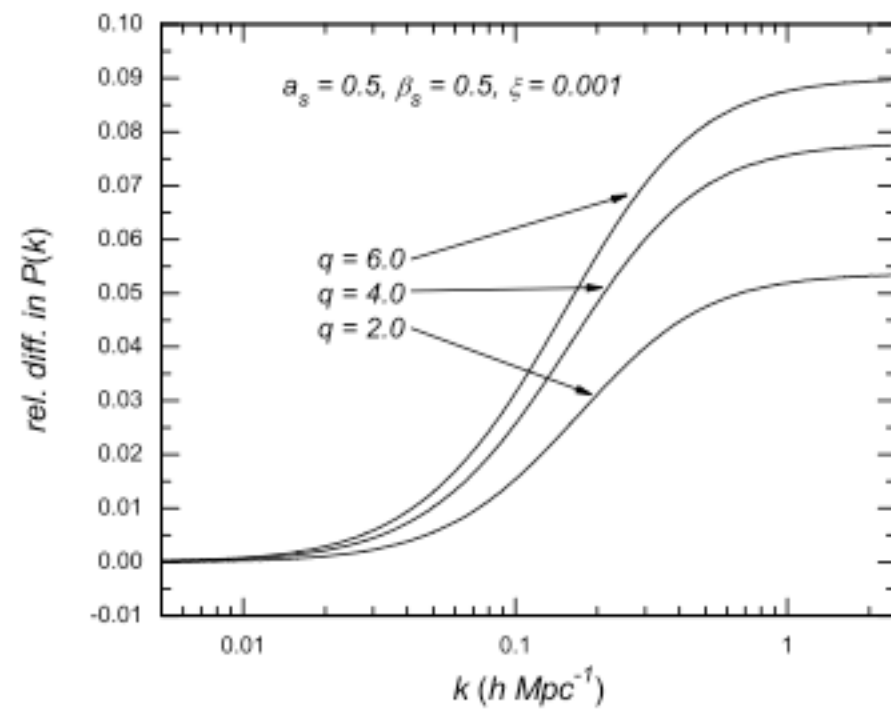
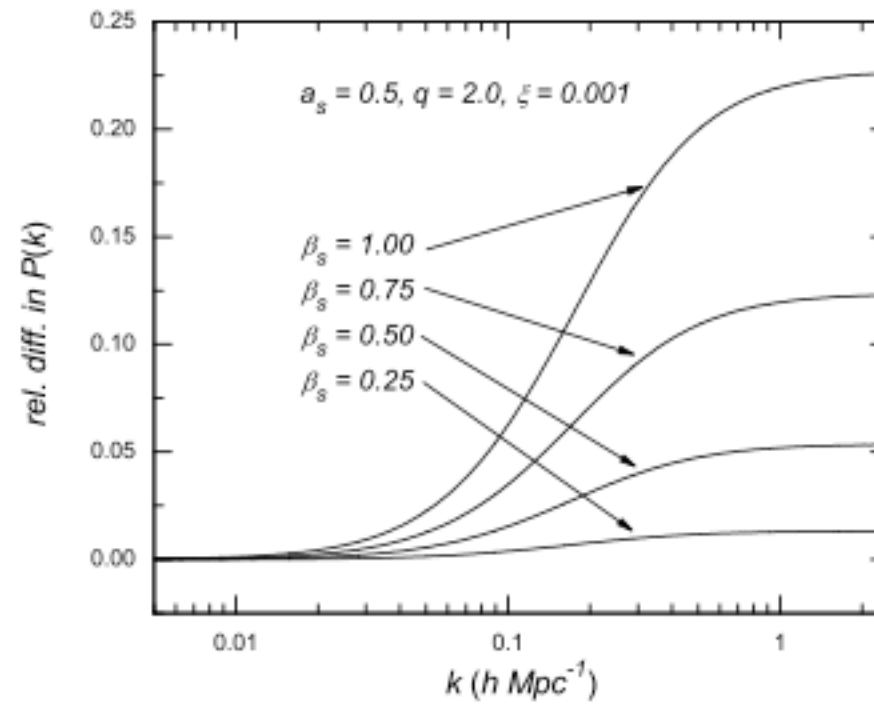
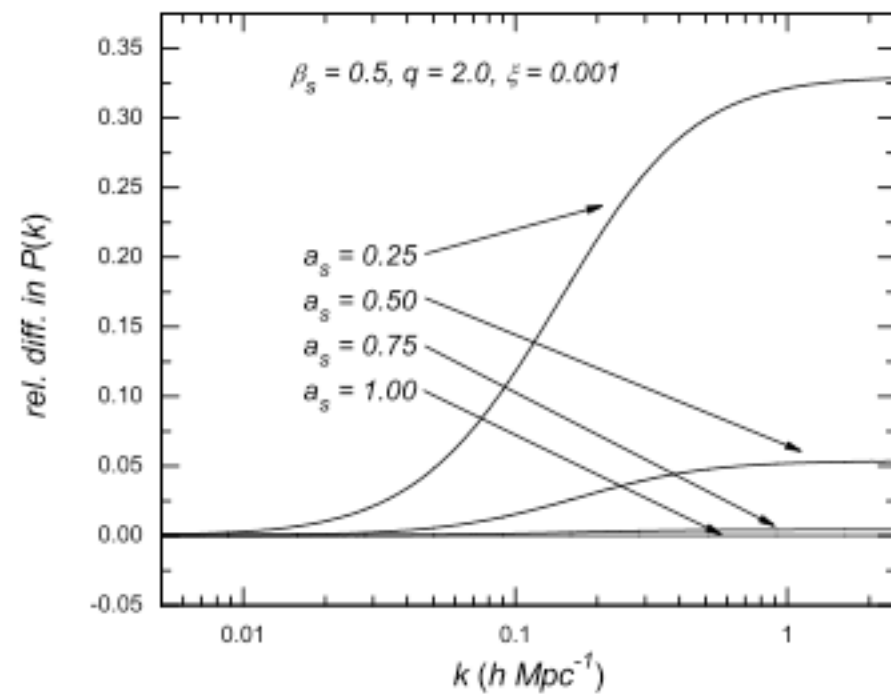
generalised symmetron

$$m(a) = m_{\star} \left[1 - \left(\frac{a_{\star}}{a}\right)^3\right]^{1/p}, \quad \beta(a) = \beta_{\star} \left[1 - \left(\frac{a_{\star}}{a}\right)^3\right]^{1/q},$$





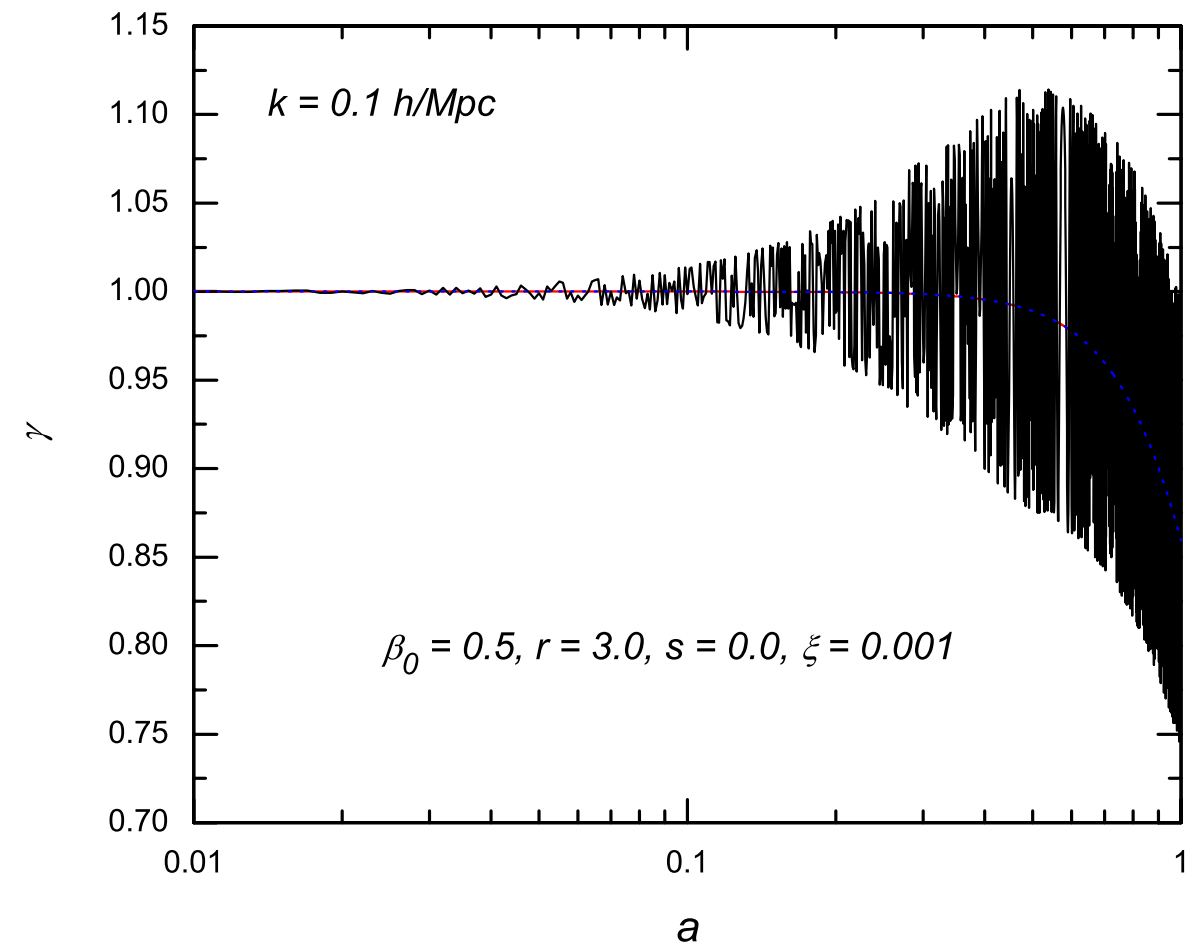
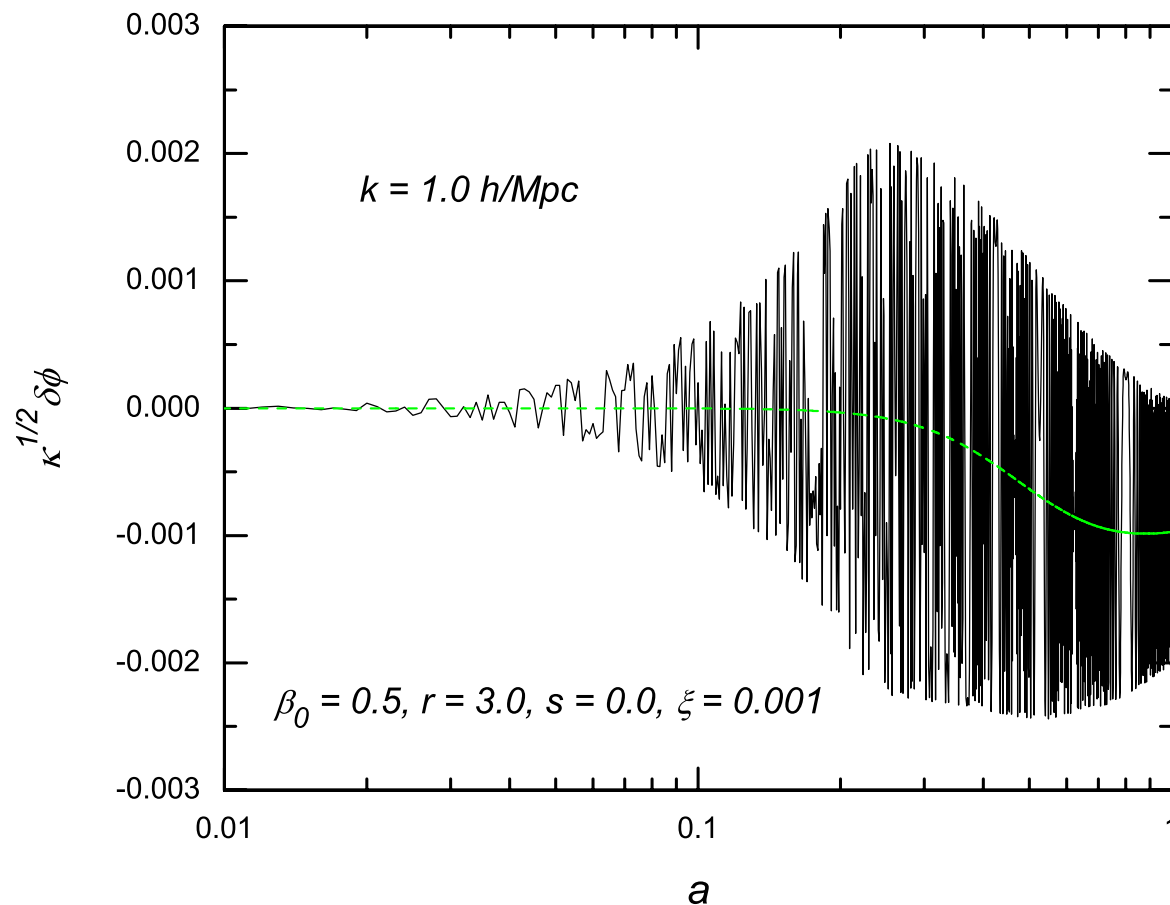
# symmetron



In Jordan frame the usual param is

$$\Phi = \gamma(k, a)\Psi$$

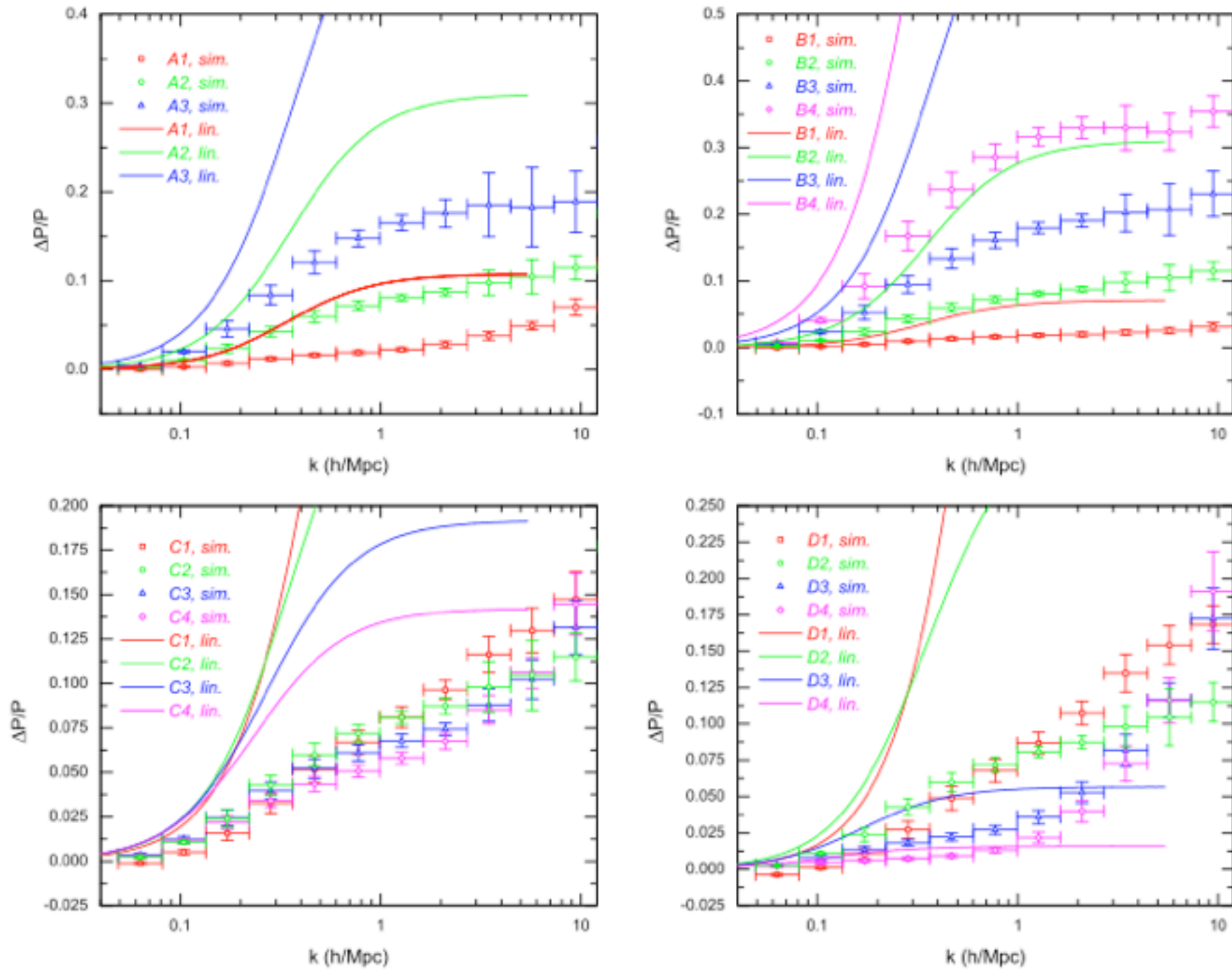
for comparison, this param is fine in linear regime, but doesn't capture the full nonlinear dynamics of screened models

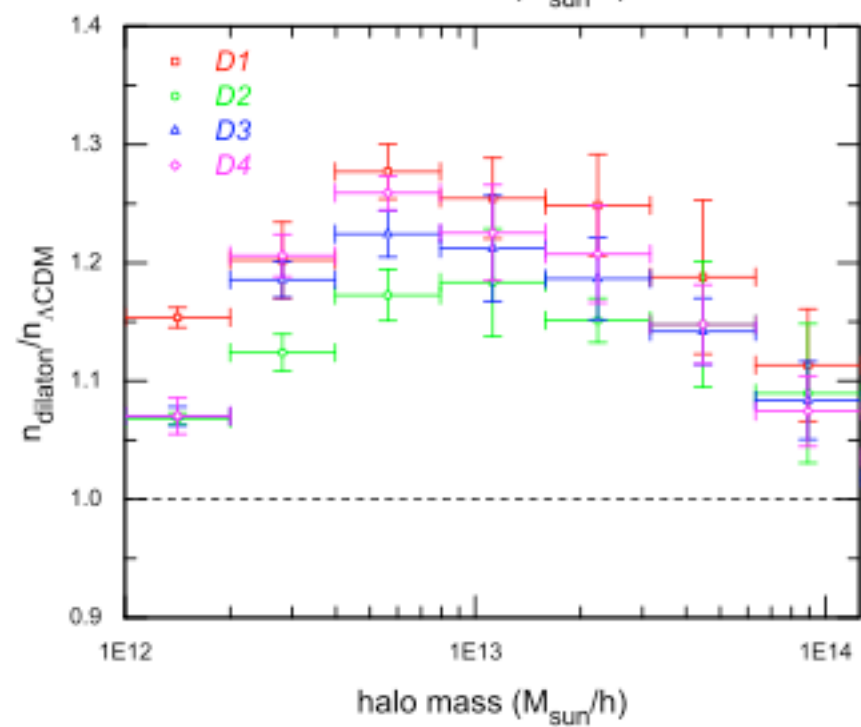
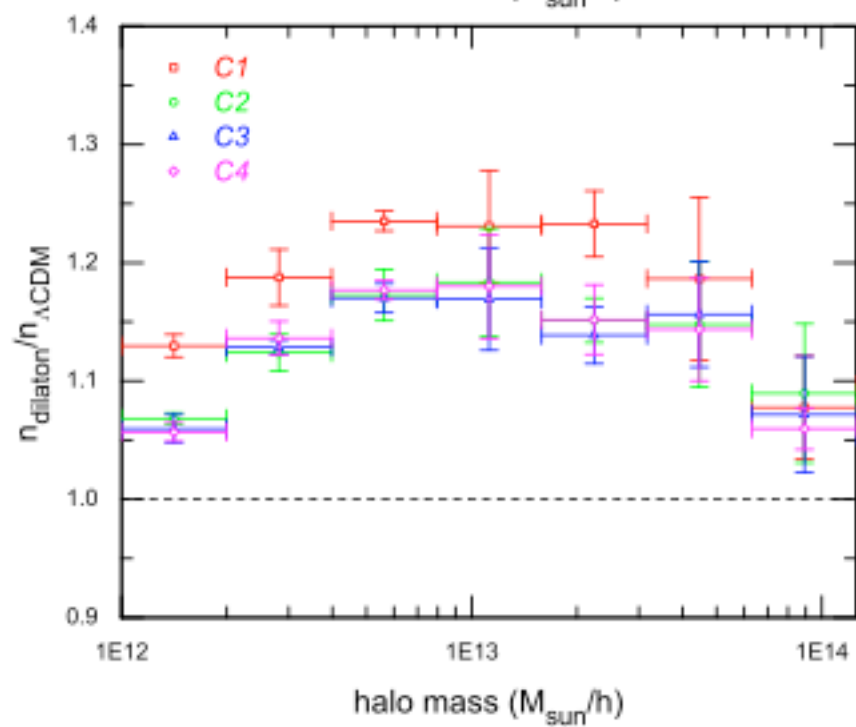
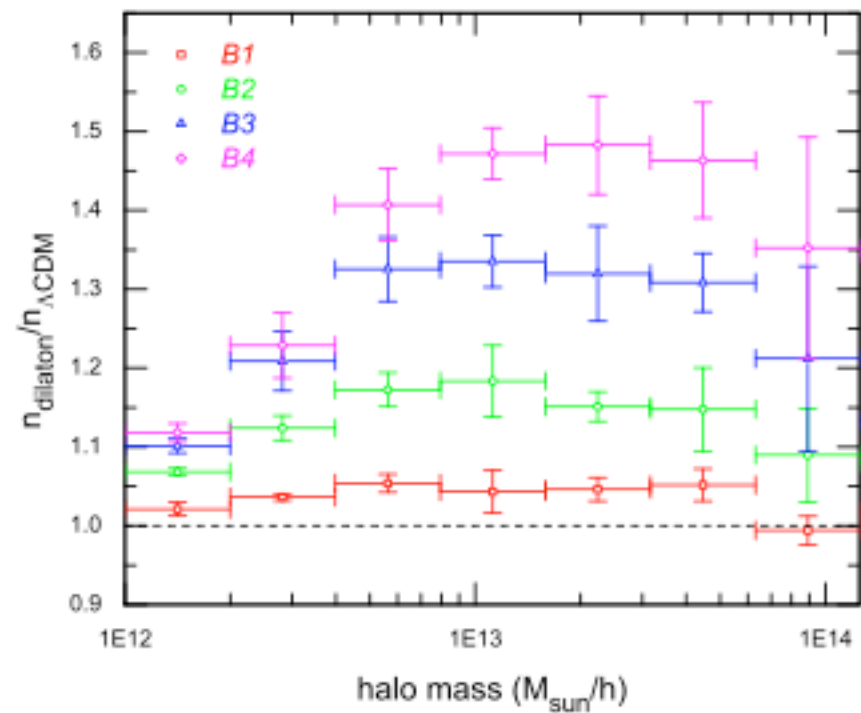
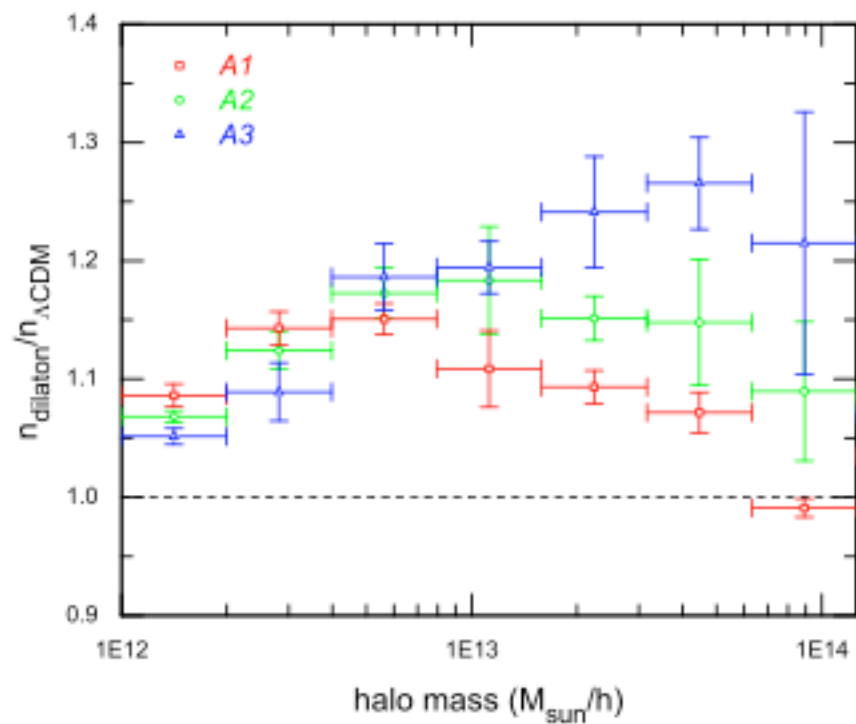


# N-body Simulations

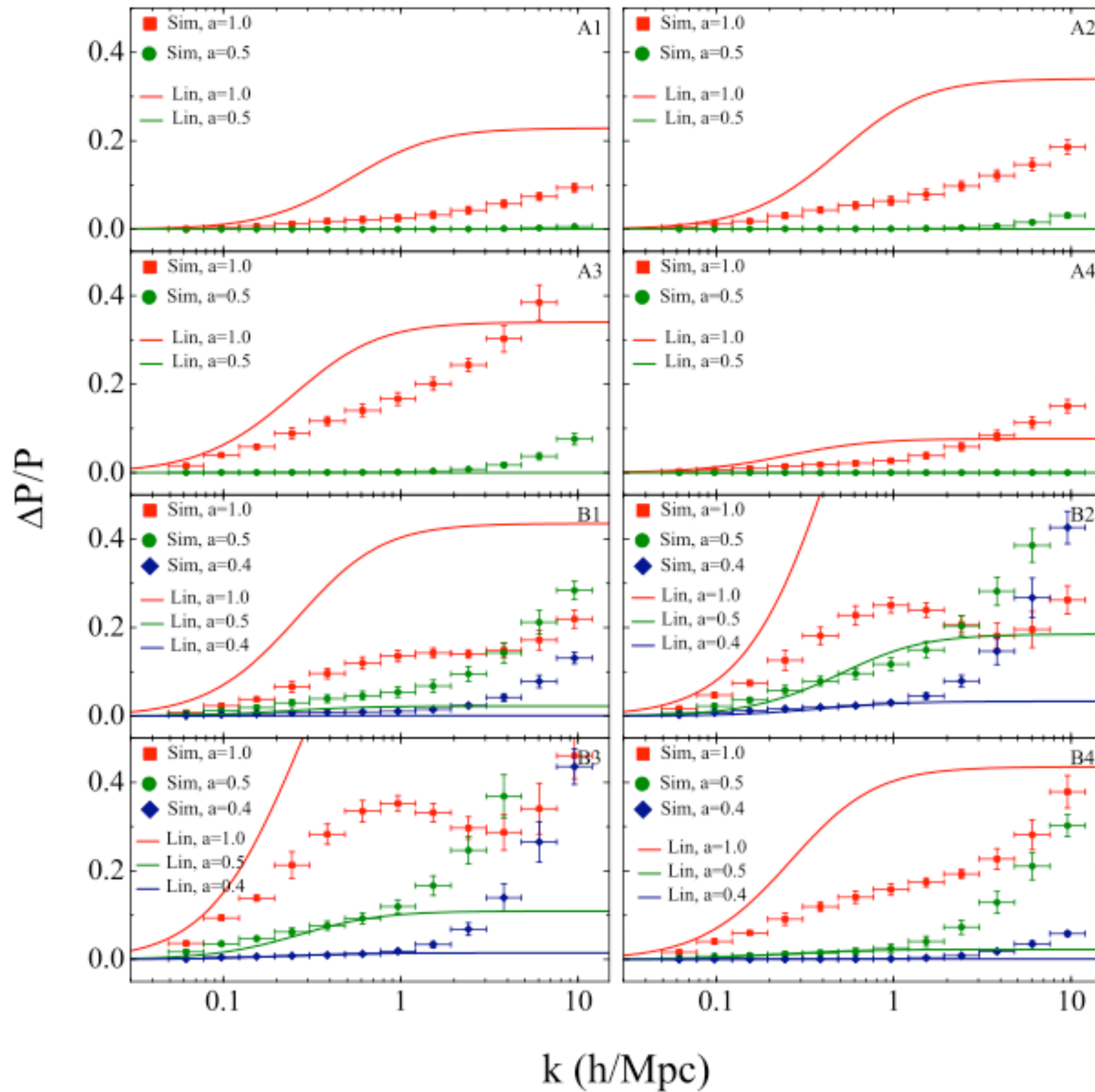
## dilatons

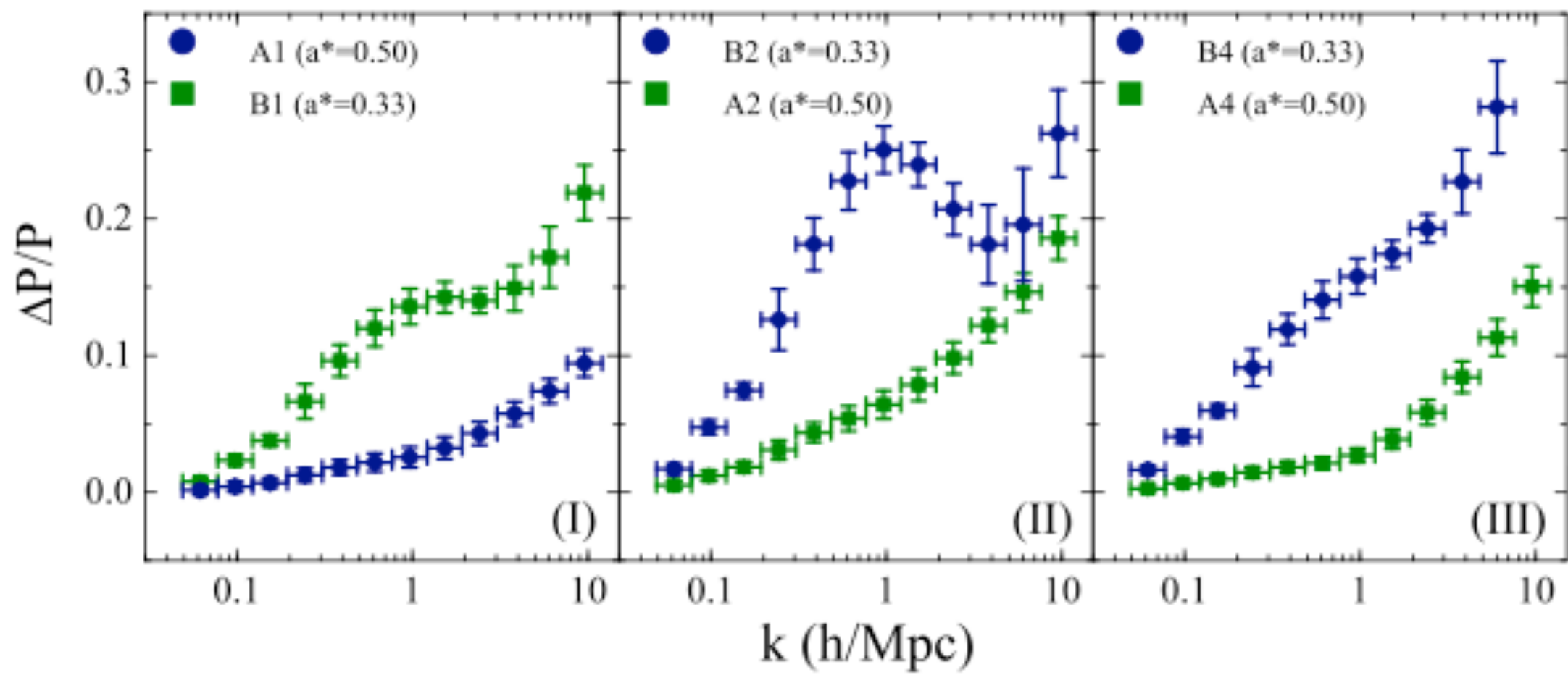
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# symmetron





One can study the variation in the fine structure constant and variation in particle masses using this parametrisation for all models.

We can use the same parametrisation to investigate the effect of modified gravity on 2lcm predictions -- to appear soon (P.Brax and S. Clesse)  
and BAO -- to appear shortly (as above + C. van de Bruck and G. Sculthorpe)



- We have presented a simple parametrisation which fully encapsulates the dynamics of screened, modified gravity models (chameleons, symmetrons and dilatons)
- Reconstruct the potential and coupling
- Studied the growth of structures in the linear and nonlinear regime.
- Apply to  $\Lambda$ CDM cosmology and BAO