

Transient dust in warm debris disks

Detection of Fe-rich olivine grains

J. Olofsson¹, A. Juhász², Th. Henning¹, H. Mutschke³, A. Tamanai⁴, A. Moór⁵, P. Ábrahám⁵, and M. Min⁶

¹ Max Planck Institut für Astronomie, Königstuhl 17, 69117 Heidelberg, Germany
e-mail: olofsson@mpia.de

² Leiden Observatory, Leiden University, PO Box 9513, NL-2300 RA Leiden, The Netherlands

³ Astrophysikalisches Institut und Universitäts-Sternwarte (AIU), Schillergäichen 2-3, 07745 Jena, Germany

⁴ University Heidelberg, Kirchoff-Institut for Physics, 69120 Heidelberg, Germany

⁵ Konkoly Observatory, Research Centre for Astronomy and Earth Sciences, Hungarian Academy of Sciences, P.O. Box 67, H-1525 Budapest, Hungary

⁶ Astronomical Institute, University of Utrecht, PO Box 80000, 3508 TA, Utrecht, The Netherlands

Received December 23, 2011; accepted

ABSTRACT

Context. Debris disks trace remnant reservoirs of leftover planetesimals in planetary systems. A handful of “warm” debris disks have been discovered in the last years, where emission in excess starts in the mid-infrared. An interesting subset within these warm debris disks are the ones where emission features are detected in mid-IR spectra, which points towards the presence of μm -sized dust grains, with temperatures above ~ 100 K. Given the ages of the host stars, the presence of these small grains is puzzling, and questions their origin and survival in time.

Aims. This study focuses on determining the mineralogy of the dust in these warm debris disks, based on Spitzer/Irs spectroscopic data, in order to provide new insights into the origin of the dust grains.

Methods. We develop and present a new radiative transfer code (DEBRA) dedicated to SED modeling of optically thin disks. The DEBRA code is designed such as it can determine dust composition and disk properties simultaneously. We make use of this code on the SEDs of eight warm debris disks, in combination with recent laboratory experiments on dust optical properties.

Results. We find that most, if not all, debris disks in our sample are experiencing a transient phase, suggesting a production of small dust grains on relatively short timescales. Dust replenishment should be efficient on timescales of months for at least three sources. From a mineralogical point of view, we find that crystalline pyroxene grains (enstatite) have small abundances compared to crystalline olivine grains. The main result of our study is that we find evidences for Fe-rich olivine crystalline grains ($\text{Mg} / [\text{Mg} + \text{Fe}] \sim 80\%$) for several debris disks. This finding contrasts with studies of gas-rich protoplanetary disks, where Fe-bearing crystalline grains are usually not observed.

Conclusions. The presence of Fe-rich olivine grains, and the overall differences between the mineralogy of dust in Class II compared to debris disks suggest that the transient crystalline dust in warm debris disk is of a new generation. We discuss possible crystallization routes to explain our results, and also comment on the mechanisms that may be responsible for the production of small dust grains. We suggest that one object may be very similar to our young Sun during its Late Heavy Bombardment phase. Finally, we underline the unfortunate limitations of unresolved spectroscopy, and discuss possible future improvements on our modeling.

Key words. Stars: individual: HD 113766 A, HD 69830, BD+20 307, HD 15407 A, HD 169666, HD 98800 B, [GBR2007] ID 8, [GBR2007] ID 9 – planetary systems: Zodiacal dust – circumstellar matter – Infrared: stars – Techniques: spectroscopic

1. Introduction

Several hundred stars are known to harbor debris disks. The first one was discovered around Vega (Aumann et al. 1984), using the Infrared Astronomical Satellite (IRAS). As detailed in the review by Wyatt (2008), debris disks are characterized by photospheric emission down to mid-infrared (IR) wavelengths. Emission is then seen in excess at longer wavelengths, caused by thermal emission arising from dust, located in circumstellar belts, heated by the stellar radiation. One key characteristic of debris disks is that they are optically thin at all wavelengths. A comparison with Solar System’s Kuiper belt or asteroid belt is often made in order to explain such systems. Planetesimals and km-sized bodies contribute marginally to the thermal emission, smaller dust grains that radiate more efficiently are therefore re-

quired to reproduce the observed emission in excess. However, the smallest dust grains are also subjected to various effects, such as Poynting-Robertson drag and radiation pressure, that can evacuate them from the system on short timescales. This suggests that the population of small dust grains is continuously replenished. Such replenishment can for example happen via destructive collisions of km-sized bodies and will eventually result in a collisional cascade that will produce the observed excess emission, and may converge to a steady-state evolution of the debris disk.

Most of the known debris disks are Kuiper belt-type disks (Carpenter et al. 2009), showing no evidences for warm dust grains in their inner regions. However, thanks to mid-IR observations, several unusual objects emerged during the last years: warm debris disks. For these systems, the emission in excess starts in the mid-IR ($\sim 5\text{--}10\mu\text{m}$), suggesting the presence a ma-

terial much closer to the central star. Our study will focus on a subset of this group: warm debris disks with emission features. These very rare objects ($\sim 2\%$ of all debris disks, Bryden et al. 2006, Chen et al. 2006) are F-, G-, K-stars, whose mid-IR spectra ($5\text{--}35\ \mu\text{m}$) show strong emission features. Such emission features are usually seen in primordial disks around young Class II objects, and are associated with μm -sized silicate dust grains. In addition to the relatively small sizes of such dust grains, the presence of a strong excess at short wavelengths ($\sim 5\ \mu\text{m}$) is indicative of a high temperature for the dust grains, suggesting the dust is located close to the star. The presence of these small dust grains close to the star is intriguing since the solution of a steady-state evolution may not be suitable anymore. At the ages of the systems, the smallest grains could have been depleted, either by radiation pressure or Poynting-Robertson drag. This means the disks could be experiencing a transient phase (e.g., Wyatt et al. 2007). In that case, two scenarios are possible, either (i) a recent collision between two or more bodies (planetesimals, comets, asteroids), or (ii) an outer belt feeds the inner regions, in a very similar scenario to the Late Heavy Bombardment (LHB), that happened in the Solar System. To have a better understanding of the origin of this dust, and assess whether it is transient or not, one way is to study the Spectral Energy Distribution (SED) of these rare debris disks, and determine the temperature of the dust responsible for the emission, and thus its distance from the star.

A powerful diagnostic to complement such study is to make use of the emission features that arise from (sub-) μm -sized grains. The width, shape and peak positions of the emission features provide information on the composition of the dust grains (see Henning 2010 for a recent review). The most common dust species are from the silicate class, which have been detected in various environments, e.g., in the interstellar medium, winds of AGB stars, and objects from the Solar System. The main building blocks of the silicates are Si, O, Mg and Fe (then followed by Ca and Al). When modeling the dust mineralogy via spectral decomposition, one should therefore focus primarily on dust species that contained these four main ingredients. Silicates are an association of $[\text{SiO}_4]^{4-}$ tetrahedra, with inclusions of Mg or Fe (following cosmic abundances). Not only the chemical composition of the grains has an influence on the emission features; the internal periodic arrangement of the SiO_4 tetrahedra will have a significant impact on the emission features. Grains with long-range order are in their crystalline form, as opposed to the amorphous form. As explained in detail in Henning (2010), optical properties of crystalline and amorphous grains are significantly different, with multiple sharp emission features for the former one, while emission features are much broader for amorphous grains. It is therefore possible to distinguish between amorphous and crystalline grains, as well as different dust species. Modeling the dust mineralogy can therefore provide useful insights into dust composition and crystallization processes at stake in the debris disks, which in turn provide information on the origin of the dust. To assess if the grains trace a second generation or are the remnants of the primordial, gas-rich disks one has to search for possible differences between the mineralogy of dust in Class II disks compared to the dust we observe in debris disks. A result that appears to be constant for the petrology of crystalline olivine grains in protoplanetary disks is that they are Mg-rich (and hence Fe-poor, see Bouwman et al. 2008, Juhász et al. 2010, Olofsson et al. 2010). The lack of Fe-bearing crystalline grains has been studied by several authors and can be explained by crystallization processes that take place in massive disks (gas-phase condensation, Gail 2010 or thermal annealing,

Nuth & Johnson 2006). One may wonder if we observe a similar result for the dust in warm debris disks. If not, we may be witnessing the production of a second generation of dust, where crystalline grains are produced by different mechanisms.

In the following, we present in Section 2 the stellar sample of warm debris disks and briefly summarize the Spitzer/Irs data reduction. The new radiative transfer code we use is described in details in Sect. 3. We also justify which dust species we choose to model the mid-IR spectra. Section 4 presents two early results that we find relevant to SED modeling of debris disks, and we present our best fit models for individual sources in Sect. 5. The discussion of our results is then separated in three parts. In Section 6 we investigate the survival and transient nature of the dust grains as well as the time variability for some of the sources in our sample, while in Sect. 7 we emphasize on the origin of the crystalline grains that we detect in the dust belts. Finally, we briefly compare our results to what is known for objects in our Solar System in Sect. 8, and we summarize our findings and possible further improvements in Sect. 9.

2. Spitzer/Irs data and stellar parameters

The stellar sample (Table 1) was gathered from the literature and investigating the Spitzer archive, looking for warm debris disks around solar analogs (F-, G-, K-type stars). The Spitzer/Irs spectra were retrieved from the Spitzer archive and reduced using the FEPS pipeline (see Bouwman et al. 2008).

All stellar parameters, namely T_{eff} , L_{\star} and d_{\star} (effective temperature, luminosity and distance, respectively), were taken from literature. To reproduce the photospheric emission we use Kurucz model that best matches the photometric near-IR data and will minimize the excess emission in the $5\text{--}8\ \mu\text{m}$ range of the Irs data. The spectral decomposition procedure is extremely sensitive to the photospheric emission model used, especially for objects that display small mid-IR excess (e.g., HD 69830, see Fig. 12).

As several AORs contained both high and low spectral resolution observations at different wavelengths (e.g. Short-Low and Long-High), we rebin the spectra to the minimal spectral resolution in order to obtain a uniform spectral resolution over the entire wavelength range. Doing so, the fitting process do not favor spectral regions with higher spectral resolution.

Four objects in our sample are known multiple systems: HD 113766, HD 98800, BD+20 307 and HD 15407. For the first two objects, given the small separation of the different components and the slit size of the Irs instrument, the photospheric emission of the dust-free component has to be removed in order to consistently determine the dust location around the star. For HD 113766 we use the luminosities as reported in Lisse et al. (2008) for HD 113766 A (around which the dust is located). For HD 98800 we use the Keck diffraction-limited observations of the quadruple system by Prato et al. (2001) to model photospheric emission of the two components. Prato et al. (2001) indeed provide the individual fluxes for the two spectroscopic binaries A & B (the dust being located around the spectroscopic binary HD 98800 B). For these two objects we will now refer to HD 113766 A and HD 98800 B. BD+20 307 is a spectroscopic binary (3.5 days period, Weinberger 2008) of two, almost identical, late F-type stars and the disk is circumbinary, meaning no correction has to be applied to the Irs spectrum. HD 15407 is a binary system with a $21.21''$ separation. The dust is located around the primary HD 15407 A (F5) and the secondary HD 15407 B (K2, Melis et al. 2010) does not appear in the Irs

Table 1. Stellar sample

Starname	Alias	SpT	T_{eff} [K]	L_{\star} [L_{\odot}]	d_{\star} [pc]	Age [Myr]	AOR	Epoch [d-m-y]	Ref
HD 113766 A		F4V	6800	4.4	123	10-16	3579904	01-03-2004	1, 2, 3
HD 69830		K0V	5650	0.67	12.6	4000	28830720	14-01-2009	4, 5, 6
BD+20 307	HIP 8920	G0	5750	1.8	92 ± 11	~ 1000	14416384	15-01-2006	7, 8, 9
HD 15407 A	-	F5V	6500	4.5	54.7	80	26122496	09-10-2008	10
HD 169666		F5	6750	5.16	53.2	2100	15016960	02-07-2005	11
HD 98800 B	TWA4	K5	4600	0.62	44.9	8-10	3571968	25-06-2004	12, 13, 14, 15
[GBR2007] ID 9	2MASS J08101691-4856291	K	4250	0.39	361	30	21755392	10-06-2007	16
[GBR2007] ID 8	2MASS J08090250-4858172	G	5750	0.8	361	30	21755136	16-06-2007	16

References: (1) Chen et al. (2005), (2) Chen et al. (2006), (3) Lisse et al. (2008), (4) Beichman et al. (2005), (5) Lisse et al. (2007), (6) Beichman et al. (2011), (7) Weinberger (2008), (8) Zuckerman et al. (2008), (9) Weinberger et al. (2011), (10) Melis et al. (2010), (11) Moór et al. (2009), (12) Koerner et al. (2000), (13) Prato et al. (2001), (14) Furlan et al. (2007), (15) Verrier & Evans (2008), (16) Gorlova et al. (2007)

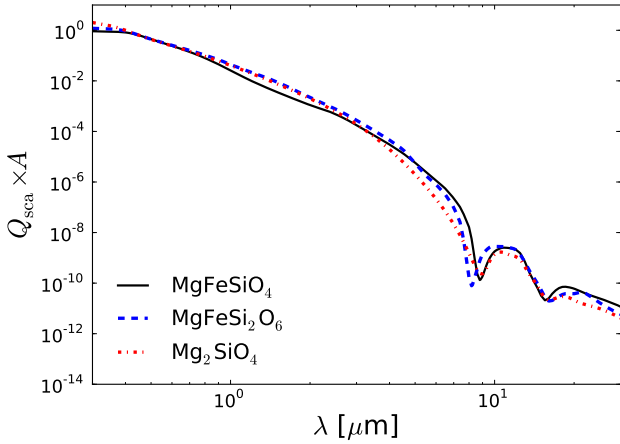


Fig. 1. $Q_{\text{sca}} \times A$ for $0.1 \mu\text{m}$ -sized grains as a function of wavelength for three dust species. Q_{sca} being the mass scattering coefficients and A the albedo of the grain (see Eqn 1).

slit, therefore no correction to the spectrum has to be performed for this object.

3. Modeling the optically thin emission

To model the SEDs including the Irs spectra, we use the “DEBRIS disk RAdiative transfer” code (DEBRA hereafter), a radiative transfer code dedicated to optically thin disks. Originally adapted from the “Debris disk radiative transfer simulation tool” (Wolf & Hillenbrand 2005), the code has been improved in order to determine the dust mineralogy as probed by the Irs spectra.

3.1. The disk setup

The DEBRA code is designed to compute the SED of a “star + disk” system in the optically thin regime. The star is defined by its Kurucz model and distance d_{\star} , and the radiative transfer code can then compute the reprocessed (thermal and scattered) emission arising from the circumstellar dust. The disk is defined by the following parameters: inner and outer radii r_{in} and r_{out} , mass of dust M_{dust} and a volume density power exponent α (≤ 0). The density profile is assumed to be a power-law ($\Sigma(r) = \Sigma(r/r_0)^\alpha$) between r_{in} and r_{out} (sampled on $nr = 80$ radial points). For a given set of parameters, one can then compute

the total emission arising from the disk over a large wavelength range (for more details, see Wolf & Hillenbrand 2005). For all the sources presented in this study, we choose $\lambda_{\text{min}} = 0.3 \mu\text{m}$ and $\lambda_{\text{max}} = 100 \text{ mm}$. The broad wavelength range is required to sample the heating (stellar) and cooling (dust) radiation field to calculate the dust temperature correctly.

We can question whether scattered light emission should be considered when trying to model mid-IR (5–35 μm) spectra. If the albedo A of a single particle of size s is defined as follows:

$$A = \frac{Q_{\text{sca}}(\lambda, s)}{Q_{\text{abs}}(\lambda, s) + Q_{\text{sca}}(\lambda, s)}, \quad (1)$$

where $Q_{\text{abs}}(\lambda, s)$ and $Q_{\text{sca}}(\lambda, s)$ are the mass absorption and scattering coefficients (see Sect. 3.2 for more details), then the scattered light emission of this single dust grain is proportional to $Q_{\text{sca}}(\lambda, s) \times A$ (Wolf & Hillenbrand 2005). Figure 1 shows how this quantity decreases dramatically with increasing wavelengths, for three different amorphous dust species, for $s = 0.1 \mu\text{m}$. At $5 \mu\text{m}$ where Irs spectra begin, the contribution of scattered light is ~ 4 orders of magnitude smaller than at $1 \mu\text{m}$. This overall means that the fitting process, which focuses on the Irs spectra, can hardly handle the scattered light emission as its contribution is only significant at shorter wavelengths.

The last free parameter of our model is the grain size distribution, which we assume to be a differential power-law of index p ($dn(s) \propto s^p ds$, $p \leq 0$) between s_{min} and s_{max} . One given model is therefore the sum of the stellar flux and thermal emission from several dust species of different sizes and is defined by four parameters r_{in} , r_{out} , α and p . The mass of dust M_{dust} within the disk is determined when fitting the dust composition over the Irs spectral range. The fitting is performed via the Levenberg-Marquardt algorithm on a linear combination of all the thermal emission contributions from the considered dust species. For a given grain size and dust composition, the thermal emission flux F_{emis} , received by an observer at distance d_{\star} , is computed as follows:

$$F_{\text{emis}}(r, \lambda, s) = \pi \times \left(\frac{s}{d_{\star}}\right)^2 \times Q_{\text{abs}}(\lambda, s) \times B_{\lambda}(\lambda, T_d), \quad (2)$$

where $B_{\lambda}(\lambda, T_d)$ is the Planck function at temperature T_d , which depends on grain sizes, compositions and distance from the star r (see Sect. 3.2.2). Finally, to determine the best fit over the entire SED (i.e., the disk geometry as well as the dust composition) we use a genetic algorithm (pikaia, Charbonneau 1995) over the four free parameters mentioned above. The code iterates over 100 generations, with 40 individuals per generation, and fits the

dust composition at every iteration. Each time, the goodness of the fit is estimated with a reduced χ_r^2 and the best fit is the one with the lowest χ_r^2 .

3.2. The dust setup

The Irs spectra of the objects in our sample show prominent emission features that are associated with warm (sub-) μm silicate dust grains. In order to reproduce the data we use several dust species, from amorphous (with olivine and pyroxene stoichiometries) and crystalline silicates (forsterite and enstatite) to crystalline β -cristobalite silica and amorphous carbon. These amorphous or crystalline dust grains, or their associated emission features, have been observed numerous times in disks around young Class II objects (e.g., Bouwman et al. 2001, 2008, Juhász et al. 2010, Oliveira et al. 2011, Olofsson et al. 2009, 2010, Sargent et al. 2009b), and disks around evolved Class III objects (e.g., Beichman et al. 2011, Lisse et al. 2007, 2008, 2009, Weinberger et al. 2011). Amorphous grains of olivine ($\text{Mg}_x\text{Fe}_{1-x}\text{SiO}_4$) and pyroxene ($\text{Mg}_x\text{Fe}_{1-x}\text{Si}_2\text{O}_6$) stoichiometries are also the dominant species in the interstellar medium (~98%, Kemper et al. 2005, Min et al. 2007). They are the major building blocks for the dust content in circumstellar environments. And yet, many uncertainties remain on the optical properties of these grains. For instance Tamanai et al. (2006) show the differences in band positions and band widths of the emissivity profiles for several species (crystalline forsterite and enstatite or amorphous Mg_2SiO_4) when using two different laboratory techniques (KBr pellets versus aerosol measurements). These measured emissivity profiles are also different from mass absorption coefficients calculated using optical constant measurements. One must be aware of these unfortunate limitations, and interpret carefully the results from spectral decomposition when using several dust constituents. This is the main reason why we focus our study on silicate dust grains and choose not to include additional, more complex dust species, to limit degeneracies in the fitting process. We would only have included other dust components when needed to reproduce the data. All the dust components used in this study are listed in Table 2.

3.2.1. Mass absorption coefficients

There are several ways to obtain the mass absorption coefficients, such as optical constant combined with a scattering theory or using laboratory extinction measurements (Henning & Mutschke 2010). When possible (and relevant, see Sect. 3.2.3) we opt for optical constant as it enables us to separate the contributions from different grain sizes, while aerosol measurements do not disentangle such contributions as the sample is a mixture of grains with different sizes. For the amorphous silicate grains we use the Distribution of Hollow Spheres (DHS) scattering theory, with a filling factor f_{max} of 0.7. Min et al. (2007) have shown these mass absorption coefficients best match the extinction profile toward the galactic center at 10 and 20 μm . For the crystalline species we also use the DHS scattering theory, but with a filling factor f_{max} of 1.0, for which the mass absorption coefficients usually best match the observed band positions and compare well with other laboratory measurements (see Juhász et al. 2010). For the amorphous carbon grains we use the DHS theory with a filling factor f_{max} of 1.

We compute mass absorption coefficients for ten discrete grain sizes, logarithmically spaced between $s_{\text{min}} = 0.1 \mu\text{m}$ and s_{max} . For the crystalline silicate and silica grains, we limit s_{max}

to 1 μm . According to crystallization models, one does not expect to produce pure crystals of several μm via thermal annealing. Grains are more likely to grow via collisional aggregation of both small crystalline and amorphous material (e.g. Min et al. 2008). Additionally, optical properties of large crystalline grains can become degenerate with optical properties of other dust species, since the recognizable emission features of small crystals fade away when increasing s . Limiting s_{max} to 1 μm is in line with previous studies of dust in protoplanetary disks (e.g., Bouwman et al. 2008, Juhász et al. 2010, Olofsson et al. 2010) where crystalline grains are found to be much smaller than amorphous grains. Concerning the maximum grain size for amorphous grain, we tested several values for s_{max} and the outcomes of this exercise are discussed in Sect. 4.2.

3.2.2. Temperature determination

Young massive Class II circumstellar disks contain large amount of dust, but also contains gas. When modeling Irs spectra for these objects, one can therefore assume all the dust grains, located in the optically thin regions of the disks, are in thermal contact and have the same temperature at a given distance r . However debris disks are devoid of gas and the previous assumption cannot hold anymore as no medium can efficiently support thermal exchange between dust grains. Consequently, we have to compute the temperature of the grains individually. The temperature T_d of a given dust grain does not only depend on its distance r to the star, but also on how efficiently it can absorb stellar flux, which is traced by $Q_{\text{abs}}(\lambda, s)$.

$$r(s, T_d) = \frac{d_\star}{2} \times \sqrt{\frac{\int_\lambda F_\star(\lambda) \times Q_{\text{abs}}(\lambda, s) d\lambda}{\int_\lambda \pi \times B_\lambda(\lambda, T_d) \times Q_{\text{abs}}(\lambda, s) d\lambda}}, \quad (3)$$

where $F_\star(\lambda)$ is the stellar flux at distance d_\star from Earth. Equation 3 shows how r can be related to T_d , via the corresponding Planck function $B_\lambda(\lambda, T_d)$. But it also shows that r depends on the mass absorption coefficients used, especially in the near-IR where $F_\star(\lambda)$ peaks. For the F-, G-, K-type stars of our sample F_\star peaks around $\sim 0.5\text{--}0.6 \mu\text{m}$, a wavelength included in our spectral range ($\lambda_{\text{min}} = 0.3 \mu\text{m}$). Reversing numerically Eqn. 3, for each dust species and for each grain sizes we are able to consistently compute their temperature depending on their distance to the star.

3.2.3. Forsterite and crystalline olivine grains

Mass absorption coefficients are the tools used to interpret observations, and one should always make sure to use the most relevant ones. As mentioned previously, mass absorption coefficients suffer from uncertainties at different levels. Laboratory measurements are unfortunately subjected to various environmental effects such as electromagnetic polarization of the embedding medium (KBr, CsI or PE) when using the pellet technique (see Tamanai et al. 2006, 2009). Additionally when using optical constant, scattering theories have their own caveats, as discussed for example in Min et al. (2007) for both Mie and DHS theories. By inspecting the emission features associated with forsterite grains, we can determine which mass absorption coefficients are best suited to reproduce the observations. Figure 2 displays Q_{abs} values from different laboratory measurements for two wavelength ranges. From bottom to top, the green line shows mass absorption coefficients calculated with the DHS theory from optical constant of Sogawa et al. (2006), blue and red

Table 2. Dust species used. State refers to the amorphous (“A”) or crystalline (“C”) form.

Species	State	Chemical formula	Density [g.cm ⁻³]	Mg / [Mg + Fe] [%]	Shape	Ref
Amorphous silicate (olivine stoichiometry)	A	MgFeSiO ₄	3.71	50	DHS	1
Amorphous silicate (olivine stoichiometry)	A	Mg ₂ SiO ₄	3.88	100	DHS	2
Amorphous silicate (pyroxene stoichiometry)	A	MgFeSi ₂ O ₆	3.20	50	DHS	1
Amorphous silicate (pyroxene stoichiometry)	A	MgSiO ₃	2.71	100	DHS	2
Mg-rich olivine (San Carlos sample)	C	Mg _{1.9} Fe _{0.1} SiO ₄	3.33	95	Aerosol	3
Fe-rich olivine (Sri Lanka sample)	C	Mg _{1.6} Fe _{0.4} SiO ₄	3.33	80	Aerosol	3
Ortho-enstatite	C	MgSiO ₃	2.78	100	DHS	4
Silica (β -cristobalite)	C	SiO ₂	2.27	-	Aerosol	5
Carbon	A	C	1.67	-	DHS	6

References: (1) Dorschner et al. (1995), (2) Jäger et al. (2003), (3) Tamanai & Mutschke (2010), (4) Jäger et al. (1998a), (5) Tamanai (2010), (6) Jäger et al. (1998b, “cel600”, “cel800”, “cel1000”: cellulose pyrolyzed at 600, 800 and 1000 °C, respectively)

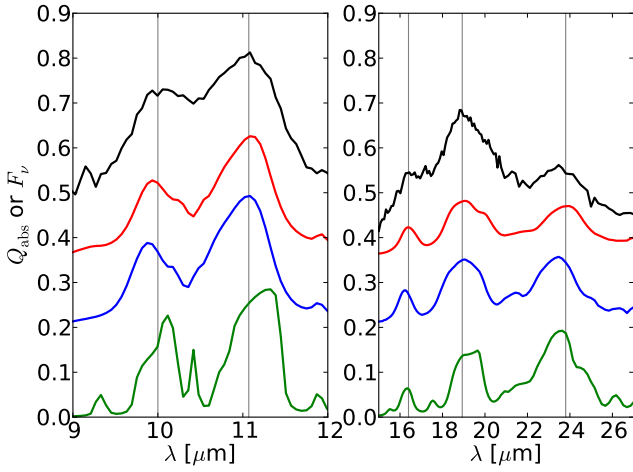


Fig. 2. From bottom to top: Q_{abs} values for forsterite (green, Sogawa et al. 2006), Mg-rich (blue), Fe-rich olivine grains (red) and stellar subtracted spectrum of HD 69830 in Jy (black). Vertical lines show peak positions of the main emission features in the spectrum. To ease readability, Q_{abs} values and the spectrum are scaled and offset

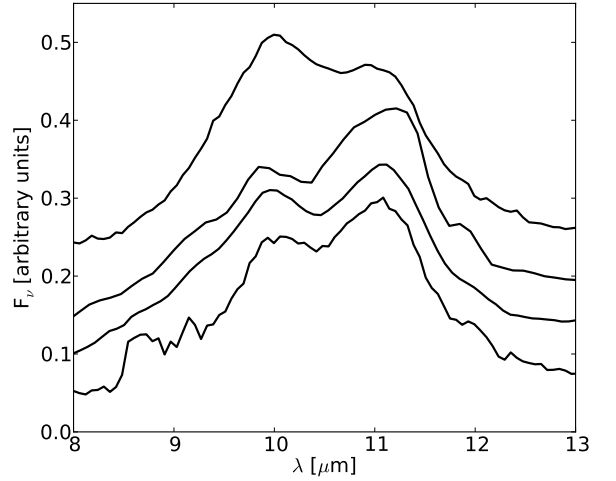


Fig. 3. 10 μm features for photosphere subtracted spectra showing the most prominent crystalline olivine features (from bottom to top: HD 69830, HD 113766 A, [GBR2007] ID 9 and [GBR2007] ID 8). Spectra have been scaled and offset

lines represent aerosol measurements for Mg-rich and Fe-rich olivine grains, with Mg fractions (= Mg/[Mg + Fe]) of 95 and 80 %, respectively (Tamanai & Mutschke 2010) and the black line is the photosphere subtracted spectrum of HD 69830 (Q_{abs} values and spectrum have been scaled and offset to ease readability). Figure 3 shows the photosphere subtracted spectra (with different scalings and offsets) for four objects showing prominent emission features associated with crystalline olivine grains (from bottom to top: HD 69830, HD 113766 A, [GBR2007] ID 9 and [GBR2007] ID 8). Band positions and shapes are similar for the four sources and compare very well with the two aerosol measurements shown in Fig. 2. The great advantage of the aerosol technique is that it minimizes the contamination by environmental effects. As the mass absorption coefficients computed with the optical constant from Sogawa et al. (2006) in combination with the DHS scattering theory do not match the observed band positions, we choose not to use them. According

to Mutschke et al. (2009), such differences around 10 μm arise from a simplified treatment of the crystal anisotropy in scattering theories such as DHS. Therefore, we instead use the aerosol mass absorption coefficients for Mg- and Fe-rich olivines to represent for crystalline olivine grains in our modeling approach.

The main difference between these two samples is the iron content. The Mg fractions are ~ 95 and 80% for both samples (produced in San Carlos, USA and Ratnapura, Sri Lanka, respectively). The iron content mostly makes a difference at wavelengths longer than 20 μm (Fig. 2, right panel). For the Fe-rich olivine sample, the 19 μm feature is slightly broader and the feature at 23–24 μm is shifted toward longer wavelengths, a result consistent with the measurements from Koike et al. (2003). These differences will enable us to assess to a fine level the iron content in crystalline olivine grains.

The disadvantages of using the aerosol measurements instead of optical constant are: (i) the overall level of Q_{abs} values is difficult to constrain. The experiment is non-quantitative, in

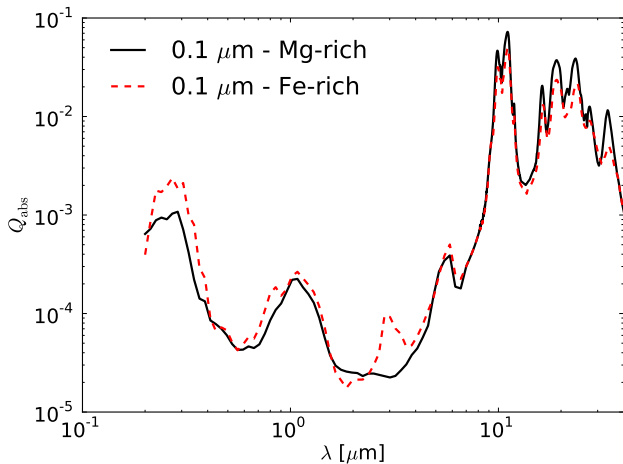


Fig. 4. Q_{abs} values for $0.1 \mu\text{m}$ -sized Mg- and Fe-rich olivine grains (solid black line and red dashed line, respectively), using the near-IR measurements from Zeidler et al. (2011) combined with aerosol measurements from Tamanai & Mutschke (2010)

the sense that it is not possible to determine how many particles have entered the aerosol chamber, and thus the total dust mass is uncertain. And (ii), the information about the contribution of different grain sizes is lost. It is unfortunately not possible to select one finite grain size while injecting them in the aerosol apparatus. However, it is still possible to select a maximum grain size, which is $\sim 1 \mu\text{m}$ in the case of the two samples discussed here.

In a recent study, Zeidler et al. (2011) determined the optical constant for several dust species in the near-IR, including the Mg- and Fe-rich olivines discussed in this Section. The increased Fe content (5 to 20%) does not dramatically change the absorption coefficients in the near-IR. This is illustrated in Fig. 4 where we show the mass absorption coefficients Q_{abs} for Mg- and Fe-rich olivine $0.1 \mu\text{m}$ -sized grains (solid black and dashed red lines, respectively). To compute these Q_{abs} values for different grain sizes, we first use the DHS theory (with $f_{\text{max}} = 1$) for different grain sizes (between 0.1 and $1.0 \mu\text{m}$) on near-IR optical constant from Zeidler et al. (2011). We then combine the resulting mass absorption coefficients with the aerosol measurements from Tamanai & Mutschke (2010) at $\lambda \sim 7 \mu\text{m}$ (where data from Zeidler et al. 2011 stop) that are representative of grains smaller than $1 \mu\text{m}$ (the size distribution in the aerosol apparatus being unknown). The motivation behind this work being that Q_{abs} values are required in the near-IR regime, in order to compute grain temperatures consistently (Eqn. 3). Additionally, as the aerosol measurements are “non-quantitative”, data from Zeidler et al. (2011) enable us to have a more accurate overall level for the aerosol Q_{abs} values.

3.2.4. Silica

Silica grains have recognizable emission features around 9 , ~ 12.5 , 16 , and $\sim 20 \mu\text{m}$. Several polymorphs of silica have been studied in detail by Sargent et al. (2009a). According to the authors, spectroscopic observations of protoplanetary disks are best reproduced by annealed silica, which contains β -cristobalite. They used the KBr pellet transmission measurements from Fabian et al. (2000), in which β -cristobalite was formed by thermal annealing of SiO_2 during 5 h at a temperature of 1220 K . This motivates our choice for using this par-

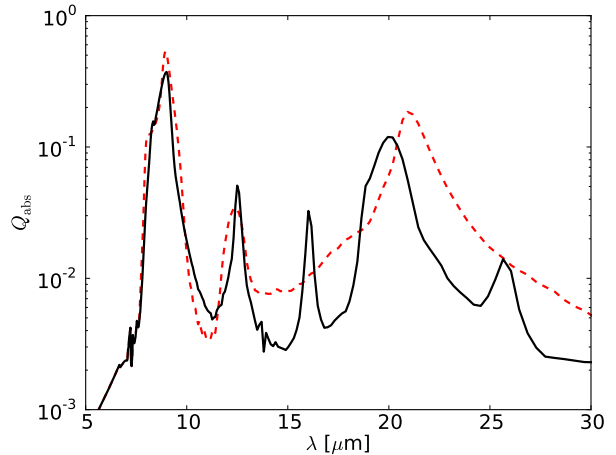


Fig. 5. Q_{abs} values for β -cristobalite silica grains from aerosol measurements (solid black line, Tamanai 2010), and $0.1 \mu\text{m}$ -sized grains computed with optical constant from Henning & Mutschke (1997, dashed red line).

ticular polymorph to represent for silica grains. However, even if Sargent et al. (2009a) discussed the possible influence of the KBr medium before concluding it should not have altered too much the band positions, we prefer to use aerosol measurements from Tamanai (2010) in order to avoid any possible contamination by the pellet medium. The maximum grain size inside the aerosol chamber is again $1 \mu\text{m}$. Figure 5 shows the mass absorption coefficients for β -cristobalite silica derived from the aerosol measurements (black line). The major emission features peak at 9.0 , 12.5 , 16.0 and $20 \mu\text{m}$, which compares well with measurements from Fabian et al. (2000, 9.1 , 12.6 , 16.1 and $20 \mu\text{m}$).

To obtain Q_{abs} values, we could not follow the same procedure as for Mg- and Fe-rich olivine grains (Sect. 3.2.3) since no data are available for wavelengths shorter than $\sim 7 \mu\text{m}$. Therefore, we use optical constant from Henning & Mutschke (1997, SiO_2 at 300 K), and combine them with optical constant from iron-free MgSiO_3 in the spectral range 0.3 – $7 \mu\text{m}$. With the DHS theory ($f_{\text{max}} = 1$) we then compute Q_{abs} values for 10 grain sizes between 0.1 and $1 \mu\text{m}$ (dashed red line for $0.1 \mu\text{m}$ in Fig. 5). We then replace these mass absorption coefficients with the aerosol measurements from Tamanai (2010) in the range 7 – $40 \mu\text{m}$. The choice of using MgSiO_3 data has in the end very little influence on the results as silica most likely needs to be mixed with impurities as discussed in Sect. 4.1, which will dominate the near-IR mass absorption coefficients.

3.3. Determining uncertainties

To assess the significance of the mineralogical results, we compute uncertainties on the relative abundances of the dust species. For each source, we use a Monte-Carlo method (as described in van Boekel et al. 2005 for instance) on the original spectrum and its uncertainties σ , which represent the statistical noise in the data. For a given object, we generated n_λ spectra, n_λ being the total number of data points in the spectrum. For each synthetic spectrum, we randomly added Gaussian noise for one wavelength point, within a distribution of width σ . We then run the DEBRA code on all synthetic spectra, with the disk parameters found for the best fit (r_{in} , r_{out} , α and p) to find the best dust composition for the n_λ synthetic spectra. The final uncertainties

on the relative abundances correspond to the standard deviations of all the realizations.

We choose not to determine uncertainties for the disk parameters. The reason being that spectroscopic observations are spatially unresolved and do not contain informations on the location of the dust. Spectral decomposition modeling of the emission features provide information about typical temperatures of the grains, not about their distance r directly. As we will detail in the following Section 4.1, the relation between r and T_d is influenced by additional effects that are poorly known at present days. Without spatially resolved observations constraining the location of the dust, assumptions have to be made in order to model available data, even though they contain a limited amount of informations.

4. Early results on SED modeling

In the course of this study, we encounter a couple of issues when modeling Irs spectra (e.g., low temperatures for grains with small Q_{abs} values in the near-IR, or maximum grain size s_{max} used for the amorphous dust grains). We detail in this Section why these issues are considered as problematic, how we correct for them and what are the implications on the dust content in debris disks. In parallel, we emphasize the limits and possible biases of SED modeling and spectral decomposition for debris disks.

4.1. “Dirty” silicates

The fitting process used in this study highly depends on the grain temperatures. As explained in Sect. 3.2.2, these temperatures are determined using the mass absorption coefficients (from optical to mm wavelengths) of the considered dust species and grain sizes. It is therefore necessary to have a precise knowledge of the mass absorption coefficients. If the Q_{abs} values are too low (e.g., $\sim 10^{-4}$) in the optical and near-IR where the stellar emission peaks, the grains will become almost “transparent” to stellar flux. Consequently, at the same distance from the star, such grains will have a much smaller temperature than grains with higher Q_{abs} values in the near-IR. If all the dust species had similar mass absorption coefficients in the near-IR this would simply translate into smaller distances r to get appropriate temperatures. However, grains such as amorphous MgFeSiO_4 or $\text{MgFeSi}_2\text{O}_6$ have higher Q_{abs} values in the near-IR and are expected to be detected (as in Class II disks). If one consider the best fit will match crystalline features such as the ones from Mg- and Fe-rich olivines, in that case, the temperature of amorphous MgFeSiO_4 or $\text{MgFeSi}_2\text{O}_6$ grains, for which the Q_{abs} values are higher in the near-IR because of a increased Fe content (50% versus 5 and 20%), will most likely be above the sublimation temperature.

To illustrate the problem, we model the spectrum of HD 69830 using the dust species listed in Table 2. For Mg- and Fe-rich crystalline olivine grains, we use the near-IR data from Zeidler et al. (2011) as described previously, which have small mass absorption coefficients in the near-IR ($\sim 10^{-4}$, see Fig. 4). In that case, the final temperatures of $0.1 \mu\text{m}$ -sized Mg- and Fe-rich crystalline olivine grains are of about 380 and 475 K, respectively, at $r = r_{\text{in}} = 4.4 \times 10^{-2}$ AU. The corresponding temperatures for $0.1 \mu\text{m}$ -sized amorphous MgFeSiO_4 grains is ~ 1970 K, a temperature at which they should vaporize. If one limits the r_{in} value for which these grains will have a $T_d \sim 1500$ K, then the crystalline olivine grains will be too cold and will consequently not contribute to the emission, resulting in a poor fit to the spectrum.

To circumvent this temperature issue, one can consider “dirty” silicates. In other words, one can assume that grains contain a fraction of impurities, such as transition metal ions, or are in contact with other dust grains that have high mass absorption coefficients in the near-IR (see, Ossenkopf et al. 1992). To account for this dirtiness, we use the Maxwell-Garnett effective medium theory to mix different fractions of carbon (“cel600”, Jäger et al. 1998b) with Mg- and Fe-rich crystalline olivine grains. The choice of using carbon over other species (e.g., metallic iron) has very little effect on the results, since the primary goal is simply to heat up the grains. This is achieved by choosing featureless Q_{abs} values for the impurities. Similar low-temperature problems will affect enstatite, silica (see Sect. 3.2.4), amorphous Mg_2SiO_4 or MgSiO_3 grains that have low Q_{abs} values in the near-IR. For these dust species, from now on, we will consider that grains can be mixed with different amounts of carbon: 1, 5 and 20% to represent for low-, intermediate- and high-temperature grains. The need for dirty grains can be justified by dynamical process at stake in debris disks. Dust grains have most likely encountered several collisions during their lifetime, which renders the assumption of pure dust grains unlikely. Consequently we exclude from the modeling the mass absorption coefficients with no impurities for Mg-, Fe-rich olivines, enstatite, silica and Fe-free amorphous grains. However, as we have no available diagnostics for this effect, we allow different levels of dirtiness for the aforementioned dust grains, by mixing different fractions of mixed carbon.

Introducing several impurities fractions reveals some of the unfortunate degeneracies of SED modeling, especially on the disk geometry. Spectroscopic observations contain a finite amount of information; the “star + disk” system is not spatially resolved by the Irs instrument. By modeling the emission features, one can determine typical temperatures of the dust grains. However, the distance r from the star, and hence the geometry of the system, will depend on how efficiently these grains can absorb the stellar radiation. In other words, if one does not know the dirtiness of grains, then one should not claim to constrain the exact morphology of the disk. At that point, spatially resolved observations become mandatory to break down the number of degeneracies on the morphology of the system in order to refine the model.

This underlines the overall complexity of modeling mid-IR spectra of debris disks in the optically thin regime. The need for dirty silicates, which is also valid for the other sources of our sample, is the first result of this study. Even if the dust grains are not in thermal contact (one single temperature $T_d(r)$ for all species and sizes), pure grains as opposed to grains mixed with impurities, cannot provide satisfying results, as it would result either in too weak contributions to the emission or, in temperatures above the sublimation radius for other dust species.

4.2. Continuum emission

When modeling the dust mineralogy of optically thick protoplanetary disks, a continuum component is included in the modeling (e.g., Juhász et al. 2010, Olofsson et al. 2010). This continuum is originally meant to account for the emission of optically thick regions of the disk (the emission features arising from the optically thin layers). As grains larger than $\sim 10 \mu\text{m}$ do not show significant emission features, they are not included in the spectral decomposition procedure. Their actual presence in the disk is not questioned, but their contribution to the emission is assumed to be part of the continuum. This assumption no longer holds for

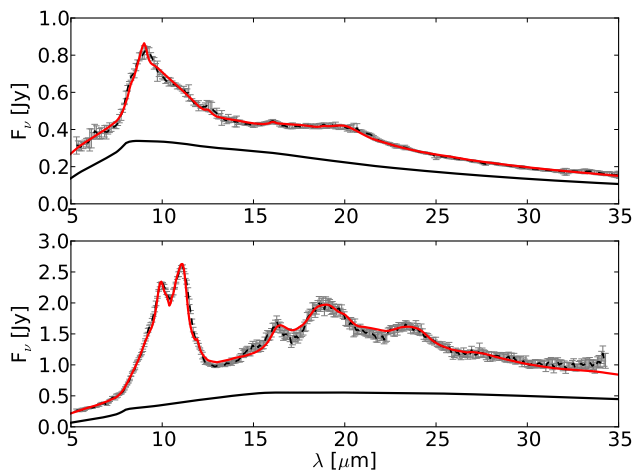


Fig. 6. Fits to stellar subtracted spectra of HD 15407 A (top) and HD 113766 A (bottom). Fits are in red, data in dashed black. Contribution of grains larger than $10\ \mu\text{m}$ is shown in solid black.

debris disks that are optically thin at all wavelengths, therefore the influence of the maximum grain size has to be tested.

As big grains are featureless, if we aim at modeling the emission features detected in spectra, the simplest assumption is to opt for $s_{\text{max}} = 10\ \mu\text{m}$. But this is not without any consequences on the results. We do not expect grain size distributions to be truncated for $s \geq 10\ \mu\text{m}$, therefore the contribution of large grains to the total emission has to be compensated by other means. One should also note that, at a given distance r from the star, 100 or $1000\ \mu\text{m}$ -sized grains are colder than μm -sized grains. The additional emission from large grains, which depends on $T_d(s)$, will still be significant at longer wavelengths. Preliminary tests have shown there are a couple of possible ways to compensate for the absence of large grains if $s_{\text{max}} = 10\ \mu\text{m}$. The first possibility is an increased abundance of carbon grains, which are featureless as well and can serve as continuum. However, given that the Q_{abs} values are rather high in the near-IR, such grains will have high temperatures and will mostly contribute at short wavelengths. Increasing the width of the disk may circumvent this issue since the outermost carbon grains will be colder and can contribute as a “cold” continuum. However it may provide a worst fit to the emission features. We assume the dust composition to be homogeneous over the entire disk, therefore if r_{out} increases by too much, some of the crystalline grains (e.g., Mg-, Fe-rich olivines or silica) will be colder as well. This will boost their emission features at longer wavelengths. For all the objects (except [GBR2007] ID 9), such features are not detected, which practically rule out the presence of such cold crystalline grains.

The second possibility to compensate for the absence of large grains is to mimic their presence by a shallower grain size distribution exponent. One can understand p as a competition between continuum emission and emission features. A steep grain size distribution will favor emission from small grains and their associated emission features, while a shallow p will favor contribution from large grains and will produce more continuum emission. Since we consider that the crystalline species have sizes between 0.1 and $1\ \mu\text{m}$, the grain size distribution exponent has little to no effect on their respective emission profiles. However, this might affect the amorphous dust grains contribution. As an example, when fitting the spectrum of HD 15407 A

with $s_{\text{max}} = 10\ \mu\text{m}$, we obtain a value of -1.70 for p , which is far from $p = -3.5$ expected for a collisional disk (Dohnanyi 1969). Such result would most likely require unusual constraints on the dynamical processes at stake in the disk. Again, as we do not expect grain size distributions to be suddenly truncated for sizes larger than $10\ \mu\text{m}$, we estimate the most straightforward (and less biased) solution is to simply increase s_{max} up to $1000\ \mu\text{m}$ so that the large grains can account for the continuum emission. In Figure 6, we show fits to stellar subtracted spectra of HD 15407 A and HD 113766 A (top and bottom panels, respectively). The solid black lines represent the flux contribution of grains with sizes between 10 and $1000\ \mu\text{m}$. The flux fractions over the Irs range are of about $\sim 60\%$ and 36% of the total flux, for both objects and this underlines the need for large grains as their contribution is non-negligible.

5. Individual source descriptions

As several sources in our sample are well-studied debris disks, we present and discuss in the following the best fits we obtain for each source and compare our results to previous studies. One should note that we focus our study solely on Spitzer/Irs data and we do not include other measurements such as MIPS photometric observations. But when available, we compare the flux predictions of our best model with these measurements. Tables 3 and 4 show the disk parameters and relative abundances for the best models, respectively. Figure 11 shows the fits to the stellar subtracted spectra as well as the residuals from the fits, which are computed as follows:

$$\text{Diff [\%]} = 100 \times \frac{F_{v,\text{obs}} - F_{v,\text{model}}}{F_{v,\text{obs}}}, \quad (4)$$

with $F_{v,\text{obs}}$ and $F_{v,\text{model}}$ being the observed and modeled fluxes, respectively. Figures 12 shows the entire SEDs as modeled using the DEBRA code. We check the optical depth is always smaller than unity meaning all sources can be modeled within the optical thin regime.

5.1. HD 113766 A

Previous studies: The Irs spectrum of HD 113766 A (the primary of the binary system) shows several strong emission features attributed to crystalline grains (9.95 , 11.0 , 16.25 and $23.5\ \mu\text{m}$). The circumprimary dust orbiting the F3-type star has been studied in detail by Lisse et al. (2008). To model the Spitzer data, the authors used several species among a vast library of 80 different chemical components. Among them, they found two types of Mg_2SiO_4 forsterite, labeled as “ForsteriteKoike” and “Forsterite038”. Unfortunately, the authors do not discuss how different these measurements are (in Lisse et al. 2008 or in the supporting online material of Lisse et al. 2006), or the implications of why both species are required. In addition to olivine and pyroxene dust species, the authors state they need additional dust species to model the spectrum, such as ferrosilite ($\text{Fe}_2\text{Si}_2\text{O}_6$), diopside ($\text{CaMgSi}_2\text{O}_6$), phyllosilicates (smectite notronite), water ice and metal sulfides (ningerite $\text{Mg}_{10}\text{Fe}_{90}\text{S}$). Concerning the latter dust species, to our knowledge, there is a major gap in the wavelength range of laboratory measurements. Optical constants are available between 0.3 and $1.1\ \mu\text{m}$ (Egan & Hilgeman 1977) and between 10 and $500\ \mu\text{m}$ (Begemann et al. 1994). The Q_{abs} values remain unknown both in the near-IR (rendering the temperature determination uncertain), and in the 5 – $10\ \mu\text{m}$ range, i.e., in the Irs spectral range. Additionally, in the range 10 – $40\ \mu\text{m}$

Table 3. Disk parameters for the best fits, as well as the maximum temperature reached

Star	r_{in} [AU]	r_{out} [AU]	α	p	M_{dust} [M_{\oplus}]	T_{max} [K]
HD 113766 A	0.4	7.1	-1.25	-3.23	4.63×10^{-3}	1452
HD 69830	0.7	1.5	-2.14	-3.54	3.75×10^{-7}	686
BD+20 307	0.2	5.2	-2.15	-3.36	2.31×10^{-4}	1510
HD 15407 A	0.4	19.2	-3.00	-3.12	7.73×10^{-5}	1531
HD 169666	59.0	60.8	-2.47	-3.60	1.12×10^{-3}	207
HD 98800 B	1.5	5.4	-0.86	-2.98	4.41×10^{-2}	510
ID 9	0.2	48.8	-1.27	-3.71	5.63×10^{-2}	934
ID 8	0.4	0.5	-2.31	-4.01	2.45×10^{-6}	885

Table 4. Relative abundances (%) and corresponding Monte-Carlo uncertainties for the best fits

Star	MgFeSiO ₄ amorphous	MgFeSi ₂ O ₆ amorphous	MgSiO ₃ amorphous	Mg ₂ SiO ₄ amorphous	Fe-rich olivine	Mg-rich olivine	Enstatite	β -cristobalite	C
HD 113766 A	0 ± 0	39.3 ± 0.2	0 ± 0	29.8 ± 0.2	17.4 ± 0.1	2.0 ± 0.1	0 ± 0	0 ± 0	11.6 ± 0.1
HD 69830	16.8 ± 0.5	27.2 ± 1.1	0 ± 0	0 + 0.5	29.3 ± 1.4	0 + 0.1	1.7 ± 0.1	1.2 ± 0.3	23.9 ± 0.2
BD+20 307	55.6 ± 0.2	15.1 ± 0.2	0 ± 0	9.1 ± 0.2	15.4 ± 0.1	0 ± 0	4.8 ± 0	0 ± 0	0 ± 0
HD 15407 A	0 ± 0	25.0 ± 0.3	28.0 ± 1.1	0 ± 0	0 ± 0	0 ± 0	0 ± 0	17.3 ± 0.1	29.7 ± 0.9
HD 169666	0 + 0.1	0 ± 0	0 ± 0	61.7 ± 1.2	0 ± 0	10.5 ± 0.6	8.0 ± 0.4	3.3 ± 0.2	16.6 ± 0.4
HD 98800 B	0 ± 0	34.3 ± 0.5	0 ± 0	37.4 ± 0.8	0 ± 0	1.8 ± 0.3	5.6 ± 0.1	0 ± 0	20.8 ± 0.2
ID 9	19.9 ± 0.8	19.9 ± 1.0	0 ± 0	23.8 ± 1.0	0 + 0.1	26.8 ± 1.4	2.0 ± 0.9	0.1 ± 0	7.5 ± 0.7
ID 8	23.0 ± 2.5	0 ± 0	0 ± 0	64.3 ± 4.0	2.9 ± 0.8	0 + 1.0	0 + 0.2	0 ± 0	9.7 ± 1.0

only one emission feature is visible and is very broad ($\sim 26\text{--}35\ \mu\text{m}$). We estimate that the lack of a clear, narrow emission feature makes the identification of this particular dust species difficult, the risk being to use metal sulfides as a proxy for continuum emission (see Sect. 4.2). Finally, concerning the geometry of the dust belt, Lisse et al. (2008) found the dust to be located at 1.8 AU from the star, with a width of 0.4 AU with typical temperatures between 420 and 490 K. The grain size distribution for their best fit has an exponent of -3.5 ± 0.2 for grain sizes between 0.1 and $20\ \mu\text{m}$, and finally the authors found a mass of $\sim 5 \times 10^{-5} M_{\oplus}$ for grain sizes up to $20\ \mu\text{m}$.

The amorphous phase: According to our model, the abundances of amorphous grains with olivine and pyroxene stoichiometries are almost comparable (30 versus 39%, respectively), with a slight edge for grains of the pyroxene group, a result different from the analysis of Lisse et al. (2008) who found a strong depletion of amorphous grains with pyroxene stoichiometry. We find that amorphous grains of the pyroxene group are all Fe-bearing while grains from the olivine group are Fe-free (no MgFeSiO₄). The abundance of amorphous carbon grains is of about 11.6%. Grains between 10 and $1000\ \mu\text{m}$ in size contribute to about 36% of the total emission over the entire Irs range.

The crystalline phase: According to our model, the crystalline content simply consists of Fe-rich olivine grains (17.4%) and a small amount of Mg-rich olivine grains ($\sim 2\%$), both species being mixed with 5% of carbon. No enstatite nor silica grains are used (and needed) in the best fit model. Inspecting the residuals of the best fit, we do not see any conclusive evidences for additional species such as the ones used by Lisse et al. (2008).

Disk geometry: The dust belt is found to be rather extended, with $r_{\text{in}} = 0.4\ \text{AU}$, close to the sublimation radius ($T_{\text{max}} \sim 1450\ \text{K}$), and $r_{\text{out}} = 7.1\ \text{AU}$. The grain size distribution follows a power-law with an exponent of $p = -3.23$ and we obtain a total mass of dust of $\sim 4.6 \times 10^{-3} M_{\oplus}$ for grain sizes up to $1000\ \mu\text{m}$.

Additional observations and comments: Our best fit returns a $70\ \mu\text{m}$ flux of 250 mJy, to be compared to the MIPS flux of $390 \pm 28\ \text{mJy}$ (Chen et al. 2011, revised from $350 \pm 35\ \text{mJy}$ originally reported in Chen et al. 2005). The $\sim 150\ \text{mJy}$ difference can be explained in two different ways. First, one may consider time variability between Irs and MIPS observations. In the case of HD 113766 A, this explanation is quite unlikely as only a few days separate the two observations (23rd of February 2004 and 1st of March 2004 for MIPS and Irs, respectively). Second, an outer dust belt might surround the inner warm dust belt. Lisse et al. (2008) considered an equivalent scenario as their model suggests the presence of two additional dust belts, one as far out as 9 AU where water ice may be located and another one at 30–80 AU. We note that according to our best model, the dust belt is quite extended (up to 7 AU), suggesting the need for cold dust grains, which may support the possibility of an outer belt. To further investigate this scenario, we have an on-going Herschel/Pacs photometry program (PI: J. Olofsson, observations shared with another team, PI: B. Zuckerman) which will provide new observations at longer wavelengths (70, 100 and $160\ \mu\text{m}$) and thus, additional constraints on the presence of an outer belt.

5.2. HD 69830

Previous studies: In addition to show emission features in its Irs spectrum, HD 69830 is known to host 3 Neptune-sized planets (Lovis et al. 2006), with semi-major axis of ~ 0.08 , 0.186 and 0.63 AU, with respective minimum masses of 10.2, 11.8 and $18.1 M_{\oplus}$. The mineralogy of the dust content has been studied in detail by Lisse et al. (2007), who modeled the Irs spectrum in a similar way as the one of HD 113766 A. According to their model, the authors detected several species. Similarly to Lisse et al. (2008), they used two types of Mg₂SiO₄ forsterite (“ForsteriteKoike” and “Forsterite038”). Fayalite, the Fe-rich end member of the olivine group was also included in their model. No amorphous grains with a pyroxene stoichiom-

etry were detected, but bronzite ($\text{Mg}_{1-x}\text{Fe}_x\text{Si}_2\text{O}_6$) ferrosilite ($\text{Fe}_2\text{Si}_2\text{O}_6$) and diopside ($\text{CaMgSi}_2\text{O}_6$) were necessary for their model to reproduce the data. Finally, the authors state they detected several carbonate species, including magnesite (MgCO_3), dolomite (CaMgC_2O_6), water ice and amorphous carbon. Lisse et al. (2007) concluded the dust belt is located at 1.04 ± 0.12 AU from the central star. According to the stability analysis from Lovis et al. (2006), this value is consistent with the 2:1 and 5:2 orbital resonances with the outer Neptune-sized planet. The typical temperature for the dust is ~ 400 K and the grain size distribution exponent for grain sizes between 0.1 and $10 \mu\text{m}$ is -3.9 ± 0.2 , with a dust mass for these grain sizes of $\sim 5 \times 10^{-8} M_\oplus$.

In a recent study, Beichman et al. (2011) investigated the time variability of the dust around HD 69830. With new spectroscopic observations and using the model from Lisse et al. (2007), they led a similar mineralogical analysis. Their model confirmed the previous detections of amorphous grains with olivine stoichiometry, both species of “ForsteriteKoike” and “Forsterite038” and fayalite. For the pyroxene group, only bronzite was detected in the new analysis (while ferrosilite and diopside went from “detected” to “marginally detected” in the new analysis). The carbonates species of magnesite and dolomite were also considered as marginally detected, as opposed to the best fit results of Lisse et al. (2007). Sulfides ($\text{Fe}_{90}\text{Mg}_{10}\text{S}$) for which an upper limit were found in the 2007 analysis was detected in the latest best fit model. According to Beichman et al. (2011), the most significant difference between the two models, is that water ice was no longer needed to reproduce the spectrum, therefore concluding the debris disk to be really dry. Based on the main mineralogical markers from both analysis, the authors state that the dust has not evolved significantly in time, but instead that the aforementioned differences arise from higher signal-to-noise ratio (SNR hereafter) of the new observations and a better data calibration for the Irs spectra. This comforts us in our conclusion that spectral decomposition is a non-trivial problem, strongly dependent on several aspects, from SNR in the data, to imperfect knowledge of mass absorption coefficients. This is why we prefer to use less dust species but focus on those that have been intensively studied in laboratory experiments, and have a good coverage of their optical data. Finally, Beichman et al. (2011) mentioned that no gas lines could be convincingly detected in their data.

The amorphous phase: According to our model, the amorphous dust content is best described by Fe-bearing amorphous grains with olivine stoichiometry (16.8%) and Fe-bearing amorphous grains with pyroxene stoichiometry (27.2%). No Fe-free amorphous grains (MgSiO_3 or Mg_2SiO_4) are required in the best fit model. The abundance of amorphous carbon is 23.9%. Grains with sizes between 10 and $1000 \mu\text{m}$ amounts for 14.7% of the total flux in the Irs range.

The crystalline phase: The total crystallinity fraction of $\sim 32\%$ is the sum of contributions from Fe-rich olivines (29.3%), with carbon additions of 1% and 5% (abundances of 17.3 and 12%, respectively), enstatite (1.7%) mixed with 5% carbon, and β -cristobalite (1.2%) most of it being mixed with 1% carbon. No Mg-rich olivine grains are required to fit the data. Inspecting the residuals between the modeled and observed fluxes, we do not see evidences for other emission features.

Disk geometry: The dust belt is located between 0.7 and 1.5 AU, values that are not in contradiction with the position of the outermost Neptune-sized planet. Since we assume the disk to be circular (eccentricity $e = 0$), the solution we find is in agreement with the stability analysis from Lovis et al. (2006). But again, we cannot claim to have precise measurements for both

r_{in} and r_{out} as the “dirtiness” of dust grains is hardly constrained. The grain size distribution exponent is $p = -3.54$, with a total dust mass $M_{\text{dust}} = 3.75 \times 10^{-7} M_\oplus$ for sizes up to $1000 \mu\text{m}$.

Additional observations and comments: Beichman et al. (2005) reported a Spitzer/MIPS flux of 19 ± 3 mJy at $70 \mu\text{m}$. Our model returns a slightly higher flux of 22 mJy, within the uncertainties of the $70 \mu\text{m}$ observations. This suggests there is no colder material that contribute up to $70 \mu\text{m}$ and that the system can be successfully described with one single dust belt close to the star. Our best fit model is consistent with the VLTI/MIDI observation and VLT/VISIR imaging of Smith et al. (2009), who constrain the warm dust in HD 69830 system to lie within 0.05–2.4 AU.

5.3. BD+20 307

Previous studies: The dust in the debris disk of BD+20 307 surrounds a spectroscopic binary of similar F-type stars, as explained by Weinberger (2008). Zuckerman et al. (2008) estimated an age of a few Gyr for this binary system. The Irs spectrum shows prominent feature at $10 \mu\text{m}$, with a shoulder at $11 \mu\text{m}$. At longer wavelengths, smaller emission features are visible at 16.5, 18.7 and $23.5\text{--}24 \mu\text{m}$. No detectable feature is seen around $33 \mu\text{m}$. Weinberger et al. (2011) modeled the spectrum of this object, using amorphous grains with olivine (MgFeSiO_4) and pyroxene ($\text{MgFeSi}_2\text{O}_6$) stoichiometries plus forsterite grains. According to their modeling, the dust is located at 0.85 AU from the binary, with typical temperatures between 450 and 500 K for the dust grains. The inferred crystallinity fraction is $\sim 20\%$ in forsterite grains. Finally, their model slightly over-predicts the MIPS flux at $70 \mu\text{m}$ (measured: 27.8 mJy, against 40.2 mJy for their model).

The amorphous phase: The amorphous dust content is dominated by Fe-bearing grains from the olivine group (55.6%), but also consists of Fe-free amorphous grains with olivine stoichiometry (9.1%) and Fe-bearing grains with pyroxene stoichiometry (15.1%). No amorphous carbon grains are required in the best fit model. Grains with sizes between 10 and $1000 \mu\text{m}$ contribute to 20% of the total flux in the 5– $35 \mu\text{m}$ range.

The crystalline phase: The remaining relative abundances are divided between Fe-rich olivine grains (15.4%) with a 5% mixed carbon fraction and warm enstatite (4.8%) in contact with 20% of carbon. Overall the dust composition we find compares relatively well with the one inferred by Weinberger et al. (2011).

Disk geometry: We find the dust belt to be located between $r_{\text{in}} = 0.2$ AU and $r_{\text{out}} = 5.2$ AU. According to Weinberger (2008), the spectroscopic binary has a separation of 0.06 AU at present days, meaning our estimation of r_{in} is compatible with the geometry of the system. The grain size distribution exponent p is -3.36 and the total dust mass is of about $2.3 \times 10^{-4} M_\oplus$ for sizes up to $1000 \mu\text{m}$.

Additional observations and comments: Our model slightly over-predicts the MIPS observation, as we obtain a $70 \mu\text{m}$ flux of 34 mJy, to be compared with 27.8 mJy. Interestingly, our fit to the $10 \mu\text{m}$ feature does not match perfectly the data, as in Weinberger et al. (2011). The feature at $10 \mu\text{m}$ looks “primitive” (or pristine, very similar to ISM dust), with a shoulder around $11 \mu\text{m}$ but no clear emission features on top of it, while our best fit model shows small features caused by the presence of crystalline olivine grains and enstatite. The inferred dust composition contains Fe-rich olivine grains because of the features detected at longer wavelengths (16.5 and $24 \mu\text{m}$). In other words, the spectrum seems to show two different dust compositions, a first one with very low crystallinity and a second one with an

increased crystallinity fraction. To reproduce the strong, pristine $10\mu\text{m}$ feature, amorphous grains need to be warm enough to reproduce the data, therefore they need to be located close to the star (0.2 AU in our case, 0.85 AU according to Weinberger et al. 2011). As we assume the disk to be homogenous, as soon as crystalline olivine grains are present they will contribute around $10\mu\text{m}$. One may argue that the temperature of the olivine grains could in principle be lower and thus contribute only at longer wavelengths. In our best fit, the mixed carbon fraction for the Fe-rich olivine grains is 5%, while we allow it to be 1% (which would return lower temperatures). Such solution is apparently not providing a better fit to the data. Allowing smaller fractions of carbon would basically mean that olivine grains are in their pure form, an unlikely prediction for dust grains that have most likely encountered collisions. The small uncertainties on the relative abundances of Fe-rich grains confirm that the fit cannot be improved by much, with the assumption of one single belt. This overall means we should consider alternative solutions.

As discussed previously, spectroscopic observations are unresolved and do not contain spatial informations. The simplest solution is always to assume that the emission probed by Irs is arising from one single dust belt. With this assumption, it appears the fit to the data could be improved. Additionally, the shape of the emission features, and thus the mineralogy suggests a two-components dust composition. Even though the $70\mu\text{m}$ MIPS observations do not suggest the presence of colder material, one can imagine the emission in excess is arising from two separated dust belts: a warm inner dust belt, containing mainly amorphous grains and a second, slightly colder dust belt with a higher crystallinity fraction. This possible scenario is also supported by the extension of the dust belt, up to ~ 5 AU, which suggests the presence of rather cold dust, further away from the star. For this object, we have an on-going Herschel/PACS photometry program that will provide observations at 70, $100\mu\text{m}$ (and hopefully $160\mu\text{m}$ if the source is bright enough). These new observations will provide additional constraints on the system and we will investigate the possibility of another dust belt.

5.4. HD 15407 A

HD 15407 A is the primary of a binary system ($21.25''$ separation), with an age estimate of ~ 80 Myr (Melis et al. 2010). The Irs spectrum mostly displays one strong emission feature peaking at $\sim 9.1\mu\text{m}$ that can be attributed to β -cristobalite SiO_2 grains.

The amorphous phase: No amorphous grains with olivine stoichiometry are required in order to reproduce the spectrum of HD 15407 A. Fe-bearing amorphous grains from the pyroxene group amount for 25%, and Fe-free grains represent 28% of the total abundance. The abundance of amorphous carbon is 29.7%. Grains with sizes larger than $10\mu\text{m}$ contribute significantly to the total emission (60%, see Fig. 6).

The crystalline phase: Beside the aforementioned species, the only remaining dust composition required to reproduce the data is β -cristobalite, for which we find an abundance of 17.3% with a mixed carbon fraction of 1%. Even though the emission feature at $12.5\mu\text{m}$ is slightly under-estimated (but still visible in the model), the $16\mu\text{m}$ feature is well-reproduced, as well as the broader feature at $20\mu\text{m}$. As shown in Fig. 5 (as well as Fig. 2 to 4 of Sargent et al. 2009a), β -cristobalite (annealed silica) is the only polymorph of SiO_2 that displays a strong feature at $16\mu\text{m}$. Therefore we are confident in the detection of this particular dust species in the disk around HD 15407 A. No additional crystalline

grains are found in the best fit, and the residuals do not show evidences for other emission features.

Disk geometry: For the best model, the dust is located in an extended belt, between 0.4 and 19.2 AU. The grain size distribution follows a power-law distribution with a slope in $p = -3.1$, and the total dust mass is $\sim 7.72 \times 10^{-5} M_{\oplus}$, for grain with sizes smaller than $1000\mu\text{m}$.

Additional observations and comments: The large extension of the dust belt (up to ~ 19 AU) underlines the need for cold dust. HD 15407 A is part of our Herschel/PACS photometry program (shared with another team, PI: B. Zuckerman) and the resulting far-IR observations will reveal if colder dust is detected around this star.

5.5. HD 169666

The emission in excess seen in the SED of HD 169666 starts at longer wavelengths ($\sim 10\mu\text{m}$) compared to the other objects of our sample, but still present emission features characteristic of small grains, especially at 24, 28 and $33\mu\text{m}$, that can be attributed to crystalline olivine grains. Moór et al. (2009) reported a MIPS flux of 21.8 ± 4.0 mJy at $70\mu\text{m}$.

The amorphous phase: The composition of the amorphous content is solely represented by Fe-free grains with olivine stoichiometry (61.7%) mixed with 1 and 20% carbon (abundances of 47.5 and 14.2%, respectively). The abundance of amorphous carbon grains is of about 16.6%. Grains with sizes larger than $10\mu\text{m}$ contribute to 4.8% of the total emission in the $5\text{--}35\mu\text{m}$ range.

The crystalline phase: Emission features seen at 24, 28 and $33\mu\text{m}$ can be reproduced using Mg-rich olivine grains, with an abundance of 10.5%. The rather broad features around $28\mu\text{m}$ are best matched with a additional contribution of enstatite (8%) and finally 3.3% of β -cristobalite are used in order to reproduce the shape around $20\mu\text{m}$. All these three crystalline species are mixed with 1% carbon.

Disk geometry: According to our best model, the dust is located at rather large distances from the star, with $r_{\text{in}} = 59$ AU and $r_{\text{out}} = 60.8$ AU. The grain size distribution exponent is $p = -3.60$ and we find a total dust mass of $1.3 \times 10^{-3} M_{\oplus}$ for grains with sizes smaller than $1000\mu\text{m}$.

Additional observations and comments: Our best model predicts a flux at $70\mu\text{m}$ of 35 mJy, which does not compare so well with the MIPS measurement of 21.8 ± 4 mJy reported by Moór et al. (2009). The current best fit reproduces well the emission features detected at longer wavelengths, however it does not match properly other, smaller emission features, especially those around 11 and $16\mu\text{m}$. Among all the models tested by the genetic algorithm, we find another solution that reproduces these two specific features slightly better (but has a higher χ_r^2). For this model, the disk is located much closer to the star ($r_{\text{in}} = 10.0$, $r_{\text{out}} = 14.8$ AU versus 59 and ~ 61 AU), therefore the olivine grains are warmer and can better reproduce the emission features around 11- $16\mu\text{m}$. To distinguish between both solutions, we refer to this one as the “warm” solution and to the original one as the “cold” solution, and both fits to the data are shown in Figure 7. As can be seen on the right panel of Fig. 7, the warm solution does not provide a good fit to the $33\mu\text{m}$ feature, which is the most prominent feature in the spectrum, and this model still over-predicts the MIPS flux at $70\mu\text{m}$ (31 versus 21.8 mJy). Compared to the cold model, the χ_r^2 increases by $\sim 6\%$. In the rest of the study, we will still consider the cold solution as the best fit model, since it provides the smallest χ_r^2 and reproduces well the strongest emission features. But, this simple exercise underlines

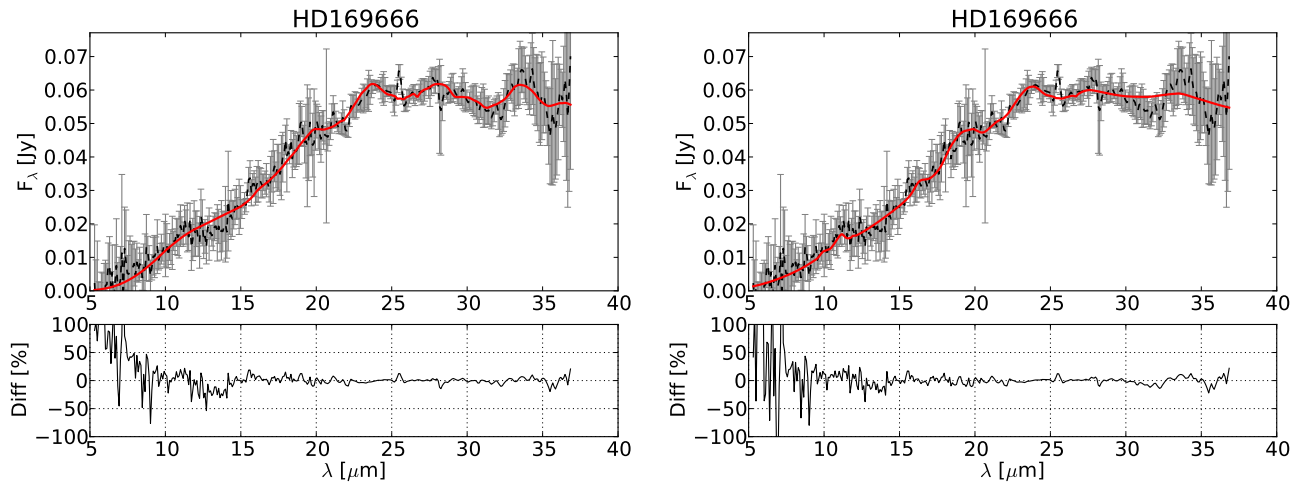


Fig. 7. Fits to stellar subtracted spectra of HD 169666 (fits in red, data in dashed black), for the “cold” (left) and “warm” solutions (right).

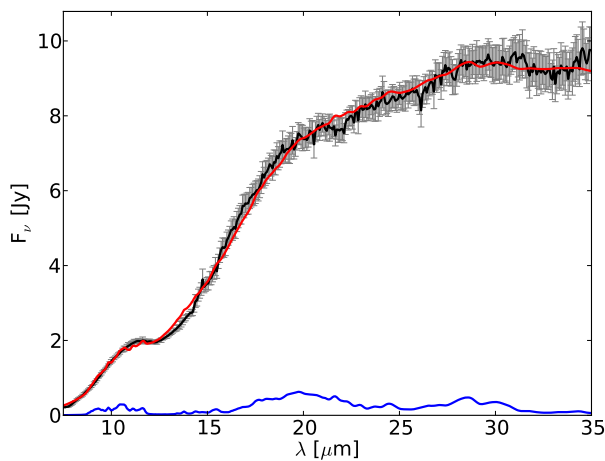


Fig. 8. Fits to stellar subtracted spectra of HD 98800 B (fits in red, data in dashed black). Blue line shows the contribution of enstatite grains.

the importance of SNR when modeling spectroscopic observations. HD 169666 is the source with the lowest SNR once the stellar spectrum is subtracted, and this renders both the analysis and interpretations difficult.

5.6. HD 98800 B

Previous studies: As the title of Verrier & Evans (2008) suggests, HD 98800 B is an unusual debris disk. The system is composed of two bound spectroscopic binaries A and B. Both systems were estimated to be 10 Myr old (Soderblom et al. 1998). With their Keck diffraction-limited observations, Prato et al. (2001) found that the dust belt surrounds the second spectroscopic binary HD 98800 B, and that the dust belt is located between ~ 2 and 5 AU. The structure of the dusty component has also been studied by several authors, for instance, Koerner et al. (2000) estimated the disk to be located at $r_{\text{in}} = 5 \pm 2.5$ AU with a width $r_{\text{out}} - r_{\text{in}}$ of 13 ± 8 AU. Furlan et al. (2007) also studied the dust content around HD 98800 B with a two components scenario: an optically thin disk located between 1.5 and 2 AU from the spectroscopic binaries and an optically thick disk at a

distance of 5.9 AU. Low et al. (2005) reported MIPS fluxes of 6.26 ± 1.25 Jy and 2.1 ± 0.42 Jy at 70 and 160 μm , respectively.

The amorphous phase: A broad feature is seen at around 10 μm , but no other emission features can be firmly detected in the Irs spectrum. Fe-bearing amorphous grains with pyroxene stoichiometry accounts for 34.3% of the total abundance. For the olivine group, no Fe-bearing grains are used in the best fit model, but 37.4% of Fe-free amorphous grains with olivine stoichiometry are used to reproduce the data. The abundance of amorphous carbon grains is 20.8%. Grains with sizes between 10 and 1000 μm contribute to $\sim 63\%$ of the total emission in the Irs range.

The crystalline phase: Even though no emission features attributed to crystalline grains are detected in the spectrum, the crystallinity fraction is not null, with overall small uncertainties. We obtain an abundance of Mg-rich olivine grains of 1.8% and an abundance of enstatite of 5.6%. The contributions of both crystalline species peak around 19–20 μm in order to try to reproduce the rather broad band between 18 and ~ 21.5 μm . Enstatite grains additionally contribute around 28–30 μm in order to match the other broad feature at the end of the spectrum. Figure 8 shows the fit to the stellar subtracted spectrum as well as the contribution of enstatite grains to the total emission. We try to fit the spectrum without crystalline dust species. The overall shape of the spectrum is reproduced, however the best fit cannot account for the broad features around 20 and 30 μm . Since the presence of crystalline grains is not supported by the detection of recognizable emission features, we prefer to consider the presence of Mg-rich olivine and enstatite grains as tentative.

Disk geometry: We find the dust to be located between $r_{\text{in}} = 1.5$ AU and $r_{\text{out}} = 5.4$ AU, values consistent with previous analysis led on this object. The grain size distribution has an exponent $p = -2.98$, and we find a total dust mass of $4.4 \times 10^{-2} M_{\oplus}$, for grain with sizes smaller than 1000 μm .

Additional observations and comments: Our model predicts fluxes of 5.33 and 1.3 Jy at 70 and 160 μm , respectively, to be compared with the values from Low et al. (2005) of 6.26 ± 1.25 Jy and 2.1 ± 0.42 Jy at the same wavelengths. The small differences may be the consequence of even larger grains contributing to the emission, but obtaining a better match of the photometric measurement at 160 μm is out of the scope of this study.

5.7. [GBR2007] ID 9

Emission in excess at mid-IR wavelengths for both [GBR2007] ID 8 and [GBR2007] ID 9 (ID 8 and ID 9 hereafter) was first identified by Gorlova et al. (2007), who estimate an age of about 30 Myr for both stars. Both sources display departure from photospheric emission at ~ 4 and $\sim 2 \mu\text{m}$, respectively. Based on SEDs, the authors propose three possibilities to explain the measured excess. First, disks are transiting between primordial toward debris disks, second a collision between two planetesimals occurred, producing large excesses and finally both objects could be background post-MS stars experiencing mass loss. However, based on near-IR spectra, the authors conclude both objects are most likely members of the NGC 2547 cluster. Therefore we assume in the following they are debris disks, and we can successfully model both SEDs including their Irs spectra in the optically thin regime. The modeling of ID 8 will be discussed in Sect. 5.8.

The spectrum of ID 9 is incredibly rich in emission features that can be attributed to crystalline olivine grains (11, 16.5, 19, 23, 28 and $33 \mu\text{m}$). The excess emission peaks at rather long wavelengths compared to other objects, which suggest either an extended or distant belt from the star.

The amorphous phase: We find a similar amount of Fe-bearing amorphous grains with olivine and pyroxene stoichiometries (19.9% for both species). For the Fe-free phase, grains from the olivine group are preferentially used (23.8%). The abundance of amorphous carbon grains is of about 7.5% and grains with sizes between 10 and $1000 \mu\text{m}$ contribute to 5% of the total emission.

The crystalline phase: The total crystallinity fraction is 29%, dominated by Mg-rich olivine grains (26.8%). Most of the Mg-rich olivine grains (19.9%) are mixed with 1% carbon, while the rest is mixed with 5% of amorphous carbon. Abundances for enstatite and β -cristobalite are of $\sim 2\%$ and 0.1%, respectively. The presence of enstatite is mostly inferred from the presence of a shoulder on the blue-side of the $10 \mu\text{m}$ feature.

Disk geometry: We find the dust to be located in an extremely extended disk around the star, with $r_{\text{in}} = 0.2$ and $r_{\text{out}} = 48.8 \text{ AU}$. The grain size distribution exponent p is -3.71 and we obtain a total dust mass of $5.6 \times 10^{-2} M_{\oplus}$ for grains with sizes smaller than $1000 \mu\text{m}$.

Additional observations and comments: The extreme extension of the disk can be explained by the presence of strong emission features at all wavelengths in the Irs spectral range. This is the only source in our sample with such strong features. Consequently, warm crystalline olivine grains must be present at the same time as much colder olivine grains that will contribute significantly at longer wavelengths. Even if all the emission features are overall well reproduced by our best model, the fit may be improved. For instance, our model over-estimates the emission feature at $23 \mu\text{m}$ while it slightly under-estimates the one at $28 \mu\text{m}$ and more significantly at $33 \mu\text{m}$. This can be explained by our assumption of a homogenous, continuous composition through the entire disk. If the model would reproduce exactly the strength of the $33 \mu\text{m}$ feature, then it would greatly over-predict the flux at $23 \mu\text{m}$. We estimate that this competition between the strength of the 23 and $33 \mu\text{m}$ features can be avoided assuming the presence of an outer belt, containing cold forsterite grains. If there is a gap in density, the flux around $23 \mu\text{m}$ will most likely decrease and will probably provide a better fit to the data. ID 9 is also one of our target in the Herschel/PACS photometry program, and the new observations will help us constraining the presence of cold dust around the star, in order to refine our model.

5.8. [GBR2007] ID 8

The Irs spectrum of ID 8 shows a strong emission feature at $10 \mu\text{m}$ as well as at 16 and $19 \mu\text{m}$. Unfortunately, as the source is faint, the Long Low module suffers from larger uncertainties, but a tentative feature at around $\sim 24 \mu\text{m}$ is detected.

The amorphous phase: Most of the dust in the disk around ID 8 is in its amorphous form (97.1%). Fe-free amorphous grains with olivine stoichiometry clearly dominate the composition, with a total abundance of 64.3%. Fe-bearing amorphous grains from the pyroxene group are also required to reproduce the data, with an abundance of 23%. Amorphous carbon grains have an abundance of 9.7%. Grains with sizes between 10 and $1000 \mu\text{m}$ contribute to 3.5% of the total emission.

The crystalline phase: The 2.9% crystallinity fraction obtained for the dust composition is represented only with Fe-rich olivine grains, mixed with 20% of carbon. The use of Fe-rich olivine grains instead of Mg-rich olivine mostly comes from the peak position of the feature around $24 \mu\text{m}$. Given the low SNR in this spectral range, the detection of Fe-rich instead of Mg-rich olivine grains is considered as tentative. The uncertainties show that Mg-rich olivine grains can also provide a satisfying fit to the data.

Disk geometry: According to our best model, the dust is located close to the star with $r_{\text{in}} = 0.4 \text{ AU}$ and $r_{\text{out}} = 0.5 \text{ AU}$. The grain size distribution is really steep, with a slope in -4.01 and the total dust mass is of about $2.45 \times 10^{-6} M_{\oplus}$ for grain with sizes smaller than $1000 \mu\text{m}$. The steep grain size distribution exponent can be explained by the rather unusual strength of the $10 \mu\text{m}$ emission feature, which requires that small grains dominate the emission.

Additional comments: Interestingly, the spectrum of ID 8 compares well with the spectrum of BD+20 307. Both spectra show a strong $10 \mu\text{m}$ emission feature, and a smaller flux level at longer wavelengths. Emission features from crystalline olivine grains are detected at 16, 19 and $24 \mu\text{m}$. The main difference between the two spectra is the emission feature at $11 \mu\text{m}$ on top of the emission feature from the amorphous grains for ID 8, and this is the reason why we obtain a good fit to the spectrum of this source and not for BD+20 307.

6. Dust survival and origin

In this Section, we discuss the possible origin of the dust belts using diagnostics for their luminosities and we investigate the main process that can be responsible for the evacuation of the smallest grains, in order to better constrain their survival and thus their origin. We also discuss the eventual presence of outer dust belts located further away from the star in order to have a better description of the systems.

6.1. Steady-state evolution or transient dust ?

As explained in Wyatt et al. (2007), the fractional luminosity of the debris disk $f_{\text{obs}} = L_{\text{IR}}/L_{\star}$ is designed to be a proxy for the disk mass. If the observed debris disk is the result of a steady-state evolution, and assuming the disk has the same age as the host star, one can determine a maximum disk mass M_{max} (and its corresponding maximal fractional luminosity f_{max}) that can survive up to the age t_{age} of the system. To compute f_{max} , we start from Eqn. 20 of Wyatt et al. (2007), inject the modeling results ($r = r_{\text{in}}$, $dr = r_{\text{out}} - r_{\text{in}}$, both in AU) and use their prescribed values for the maximal size of planetesimals in cascade ($D_c = 2000 \text{ m}$), planetesimal strength ($Q_D^* = 200 \text{ J.kg}^{-1}$),

Table 5. Observed (f_{obs}) and maximal (f_{max}) fractional luminosities for our target list.

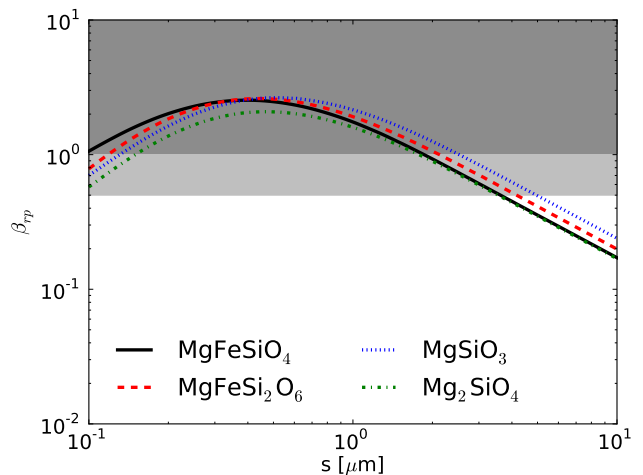
Star	f_{obs}	f_{max}	$f_{\text{obs}}/f_{\text{max}}$	Transient
HD 113766 A	7.2×10^{-2}	1.5×10^{-5}	$> 10^3$	✓
HD 69830	2.8×10^{-4}	6.2×10^{-8}	$> 10^3$	✓
BD+20 307	5.2×10^{-2}	7.1×10^{-8}	$> 10^3$	✓
HD 15407 A	7.0×10^{-3}	7.0×10^{-6}	$\sim 10^3$	✓
HD 169666	1.8×10^{-4}	2.0×10^{-5}	~ 10	×
HD 98800 B	1.5×10^{-1}	2.2×10^{-4}	~ 700	?
ID 9	4.3×10^{-1}	1.0×10^{-4}	$> 10^3$	✓
ID 8	3.2×10^{-2}	4.5×10^{-7}	$> 10^3$	✓

and planetesimal eccentricity ($e = 0.05$). The authors defined a threshold value of 100 for the ratio $f_{\text{obs}}/f_{\text{max}}$: below this value, a steady-state evolution can explain the infrared luminosity of the debris disk. Table 5 summarizes the values we obtain for f_{obs} , f_{max} and their ratios. First, the case of HD 98800 B is difficult to assess, the ratio being of about ~ 700 . Given the relatively young age of the object, Wyatt et al. (2007) concluded the collisional cascade has probably started recently, meaning the observed debris disk could possibly be the result of a steady-state evolution. Concerning HD 169666, we find a ratio of about ~ 10 for the best model (“cold” solution as discussed in Sect. 5.5) and a ratio of about 50 for the “warm” solution. As Moór et al. (2009) found a smaller value for r_{in} (~ 4 AU), their ratio $f_{\text{obs}}/f_{\text{max}}$ is above the threshold value of 100 and the authors concluded the amount of dust is too important compared to what is expected for a steady-state model. Since the SNR is overall small for this source, we cannot easily conclude on the evolutionary status of this source since f_{max} depends on the modeling results. Finally, the remaining six objects are too old for their high infrared luminosities to be explained by a steady-state evolution and are therefore considered as being in a transient phase. A dynamical event, responsible for the production of small dust grains, must have occurred so that we can observe such high fractional luminosities in the mid-IR.

It is of interest to mention that most of the sources are either in binary systems or have known planets orbiting the central star. The disks around HD 113766 A and HD 15407 A surround the primary of a binary system, and the dust belt around HD 69830 is found to encompass the orbits of three Neptune-sized planets. In the case of BD+20 307, the dust surrounds the spectroscopic binary, while the system of HD 98800 B is slightly more complicated. The disk surrounds the spectroscopic binary HD 98800 Ba–Bb, but is located around the secondary of the HD 98800 A–B system. The multiplicities of ID 8 and ID 9 remain unknown since the stars were only recently discovered. Stellar companions or massive planets could be responsible for the stirring of smaller planetesimals (see e.g., Matranga et al. 2010). Therefore, we may be witnessing the consequences of dynamical instabilities triggered by massive objects in the vicinity of the central stars.

6.2. Radiation pressure and time variability

Most of the debris disks of our sample are in transient phase, one may therefore wonder about the survival of the fresh, recently produced dust we observe. In Section 5, we have not discussed the effect of radiation pressure of the stellar radiation field on the small grains. This effect acts as an apparent reduction of the stellar mass and is the fastest way to evacuate μm -sized grains from

**Fig. 9.** Radiation pressure efficiency (β_{rp}) as a function of grain size s for amorphous MgFeSiO₄ (solid black), MgFeSi₂O₆ (red dashed), MgSiO₃ (blue dotted), Mg₂SiO₄ grains (green dotted-dashed). Shaded areas show regions where grains are gravitationally unbound.

dust belts. Using the mass absorption and scattering coefficients, and assuming isotropic scattering we compute the efficiency ratio β_{rp} , as detailed in Burns et al. (1979) for different grain sizes and different dust compositions. This ratio quantifies the competition between gravitational and radiation pressure forces. As both forces decrease as r^{-2} , the ratio between them no longer depends on r . A β_{rp} value greater than 1 means that radiation pressure overcome gravitational forces and such grains will be evacuated from the system. In the range $0.5 < \beta_{\text{rp}} < 1.0$, grains are in the unbound regime with hyperbolic orbits (see Krivov 2010), and will be evacuated from the system as well. Figure 9 shows an example of the β_{rp} -ratio for HD 15407 A, with s ranging between 0.1 and 10 μm for different amorphous dust species. An immediate result is that grains smaller than typically 3.5–4 μm are rapidly ejected from the system ($\beta_{\text{rp}} > 0.5$). If a grain is subjected to a β_{rp} ratio greater than 0.5, then the eccentricity of its orbit will increase and the grain will quickly be evacuated from the system (comparable to the dynamical timescale). In the case of HD 15407 A, the evacuation of grains smaller than $\sim 3.5 \mu\text{m}$ is expected to be very efficient, of the order of a few months.

Among the sources studied in our sample, grains around HD 98800 B, ID 9 and ID 8 are not subjected to evacuation by radiation pressure. As discussed in Beichman et al. (2011) for HD 69830, only grains with sizes in a narrow range ($0.4 < s < 0.7 \mu\text{m}$) are affected by radiation pressure. Concerning the spectroscopic binary BD+20 307 ($M_{\star} = 2.2 M_{\odot}$), grains with sizes in the range 0.3–0.7 μm (0.4–0.9 μm , depending on the composition) are expected to be evacuated from the system. This means that grains smaller than $\sim 0.3 \mu\text{m}$ may survive to radiation pressure around the binary system, and these small grains can contribute to the emission of the strong and sharp 10 μm feature. Overall for both HD 69830 and BD+20 307, since the ranges of grain sizes with $\beta_{\text{rp}} > 0.5$ are narrow, we consider this effect to be of second order for the purpose of this study. Additional constraints are necessary to lower the degeneracies in the problem in order to refine the model and assess the influence of radiation pressure. For the five sources discussed here, the survival of the dust belt in time (and hence their detection in the mid-IR) is less constraining on the mechanisms responsible for the pro-

duction of small μm -sized grains, compared to the other three sources. Since the timescale for evacuation by P-R drag is much longer (thousands of years) than the timescale for evacuation by radiation pressure (≤ 1 yr), these small grains can survive much longer being gravitationally bound to the central star.

6.2.1. Time variability of HD 15407 A

As shown in Fig. 9, amorphous grains smaller than $\sim 3.5\text{--}4\mu\text{m}$ are short-lived around HD 15407 A. Interestingly, we find that some small β -cristobalite grains, mixed with 1% carbon and with sizes in the range $0.1\text{--}0.25\mu\text{m}$, can survive to radiation pressure. Larger grains or grains mixed with higher carbon fractions have β_{rp} greater than 0.5. According to our results (Sect. 5.4) silica grains with an addition of 1% carbon provide the best match to the data. This finding slightly relaxes the constraints on the production of silica grains. However, we could not obtain a satisfying fit to the data (especially to the emission features from β -cristobalite) without amorphous grains smaller than $3\mu\text{m}$.

To assess the time variability of the debris disk around the primary, we gather photometric observations that were performed in the last decades for this object. Table 6 summarizes our findings and reports the instrument used and wavelengths of observations as well as measured fluxes. The column labeled “P_{IRS}” reports flux predictions based on the Irs spectrum and their uncertainties. To obtain these predictions and their uncertainties, we convolve the Irs fluxes to the filters widths of other instruments. The photometric observations cover a range of several decades. Unfortunately, in some cases determining an accurate epoch for the measurement is not trivial. All of the relevant facilities and instruments listed in Tab. 6 observed the vicinity of HD 15407 A several times during their mission and the quoted flux densities in the specific catalogues have been derived as a combination of the photometry extracted from individual scans or observations. However, we know that the MSX Galactic Plane survey observations were obtained between May 1996 and January 1997, while AKARI was active between May 2006 and August 2007. The IRAS mission was accomplished in 1983. Utilizing the SCANPI software we find that the relevant IRAS scans were performed between the 17th of August and the 15th of September 1983. In the case of WISE, we can determine an averaged epoch (January 2010) using the epochs of individual observations.

The measured fluxes over almost 3 decades compare very well with the fluxes derived from the Irs spectrum (observed on the 9th of October 2008), especially at $9\mu\text{m}$ where the emission feature from small β -cristobalite grains peaks (960 ± 57 versus 927 ± 47 mJy). At least one year passed between the AKARI and Spitzer observations, a timescale to be compared with a radiation pressure timescale of a couple months. We therefore conclude that the system has been collisionally active over several decades and that the population of grains smaller than $3\mu\text{m}$ is being efficiently replenished over time.

6.2.2. Time variability of HD 113766 A

In the disk around HD 113766 A, all grains (amorphous and crystalline) smaller than $\sim 3.5\mu\text{m}$ should in principle be evacuated from the system on a timescale of months. This result contrasts with the strong and sharp emission features detected in the spectra, arising from amorphous grains as probed by the $10\mu\text{m}$ feature and from crystalline olivine grains (e.g., at $11\mu\text{m}$). As for HD 15407 A we gather photometric observa-

tions of HD 113766 A, and compare them with the Irs spectrum convolved with the corresponding photometric filters. Table 6 summarizes the measurements and predictions. Again, obtaining dates of observations is difficult, however the IRAS mission was accomplished in 1983, AKARI was active between May 2006 and August 2007, and MIPS observations were performed the 23rd of February 2004. The photometric observations compare very well with the convolved fluxes, meaning the emission in excess around HD 113766 A should have been stable for the last 25 years. It indicates that small dust grains must be replenished quickly, on a timescale of a few orbits, suggesting on-going collisions in the disk.

6.2.3. Time variability of HD 169666

Computing the efficiency ratio β_{rp} for grains of different compositions (MgFeSiO_4 , $\text{MgFeSi}_2\text{O}_6$, MgSiO_3 and Mg_2SiO_4) we find that grains smaller than $4\mu\text{m}$ are expected to be short-lived in the disk around HD 169666. And yet emission features arising from small crystalline olivine grains are detected at longer wavelengths. Moór et al. (2009) already discussed the time variability of this debris disk. First, they compared two Irs spectra, obtained the 2nd of July 2007 and on the 6th of June 2008 (we model the first one in this study as it includes the 2nd order of module SL, between $5.2\text{--}8.7\mu\text{m}$). Within 3 years, no significant variations are visible, and emission features can still be detected at longer wavelengths. In a similar way as for HD 15407 A and HD 113766 A, we gather photometric observations of HD 169666 and compare them with the Irs spectrum convolved with the photometric filters. Last two columns of Table 6 show both observed and predicted fluxes. The strongest constraint on the time variability comes from the IRAS observations (1983), for which the observed and predicted fluxes compare extremely well. Spitzer/MIPS observations were performed on the 23rd of June 2005. As for the previous two sources, there must be an on-going process in the disk that continuously replenishes the small dust grain population, as the transient phase has been stable for at last 25 years.

6.3. Presence of outer, cold belts

The dust detected around the debris disks we study is transient, except maybe in the cases of HD 98800 B (given its relatively young age) and HD 169666 (given the uncertainties on the location of the dust belt because of the low SNR in its spectrum). Two main scenarios have been invoked in the literature to explain the origin of the transient dust. First, the dust may be produced directly close to the star. This scenario can be the consequence of several events. Transient dust may originate from a massive collision between two large bodies being disrupted close to the star. Even though the possibility of a steady-state asteroid belt is most likely ruled out by the high fractional luminosities of the disks, transient dust may be produced in an asteroid belt that has been recently excited and where a collisional cascade started. The second main scenario involves the presence of an outer belt dynamically unstable feeding the inner regions of the disk. Bodies from the outer belt are being scattered inward and release small grains as they collide or sublimate. As mentioned in Beichman et al. (2005) for HD 69830, a “super-comet” from the outermost regions may have been captured on a low eccentricity orbit in the inner regions and is being evaporated and broken into smaller bodies. However, detecting an outer belt further away from the star does not enable us to disentangle between the two scenarios,

Table 6. Photometric observations of HD 15407 A, HD 113766 A and HD 169666 and flux predictions from their Irs spectra convolved with the corresponding photometric filters. Second line are the epochs of the Irs observations (day/month/year).

Facilities	Epoch	HD 15407 A (09/10/2008)			HD 113766 A (01/03/2004)		HD 169666 (02/07/2007)	
		λ [μm]	Flux [mJy]	P_{IRS} [mJy]	Flux [mJy]	P_{IRS} [mJy]	Flux [mJy]	P_{IRS} [mJy]
MSX		8.28	808 ± 49	894 ± 45	-	-	-	-
AKARI/IRC	2006–2007	9.00	960 ± 57	927 ± 47	1308 ± 92	1439 ± 72	344 ± 18	333 ± 17
WISE		11.56	656 ± 32	673 ± 34	1152 ± 56	1372 ± 71	177 ± 8	182 ± 9
IRAS	1983	12.00	1050 ± 63	891 ± 45	1590 ± 79	1746 ± 89	232 ± 16	242 ± 12
AKARI/IRC	2006–2007	18.00	497 ± 33	445 ± 22	1428 ± 87	1504 ± 79	130 ± 23	107 ± 6
WISE		22.09	386 ± 23	384 ± 19	1600 ± 93	1554 ± 81	100 ± 6	99 ± 5
Spitzer/MIPS		23.67	-	-	1459 ± 30	1431 ± 74	90 ± 4	94 ± 5
IRAS	1983	25.00	432 ± 39	402 ± 20	1760 ± 123	1664 ± 87	101 ± 13	112 ± 6

since the outer belt may be stable in time and not responsible for the production of the transient dust.

The most straightforward way to infer the presence of cold, outer dust belts is via far-IR, or mm observations. For one source (HD 113766 A) in our sample, we can already mention that another dust belt is likely to be located further away from the star. Spitzer/MIPS observations at $70 \mu\text{m}$ reveal an excess of ~ 100 mJy compared to the predicted flux of our best model (250 mJy). As Irs observations are unresolved, we cannot quantify by how much the outer belt contribute to the flux measured by the spectrograph. In this study, we assume the dust is in a single, continuous disk. Consequently, the flux contribution of the outer belt in the Irs range translates into rather large values for r_{out} (7 AU). This can be explained as a need for a contribution of colder grains. The extension of disks, in the formalism of our model, may be a possible diagnostic for the presence of outer dust belts. According to our findings, HD 113766 A is not the only source for which we find an extended disk. The width of the disk around HD 15407 A is of about ~ 19 AU, which also suggests the need for cold dust grains in order to reproduce the unresolved spectroscopic data. Further observations are required in order to distinguish between the two possibilities of either an extended disk or two spatially separated dust belts. In an even more exaggerated way, the disk around ID 9 has a width of about 48 AU according to our best model. Such large width can be explained looking at the shape of the SED, where the emission still contributes significantly at longer wavelengths ($\sim 30 \mu\text{m}$). In Sect. 5.7, based on mineralogical considerations (relative strengths of emission features at 23 and $33 \mu\text{m}$) we explain why a second, outer dust belt might be present. Again, complementary observations are required to conclude on the morphology of the system. Finally, even though the MIPS observations at $70 \mu\text{m}$ do not reveal emission in excess compared to our model, we discussed the possibility of a two belts scenario around BD+20 307 (Sect. 5.3). The justification being that we observe a gradient in the dust composition: the warm inner regions probed by the $10 \mu\text{m}$ emission feature seem to contain pristine dust, while emission features from crystalline olivine grains are detected at longer wavelengths and may arise from another dust belt. With the assumption of one continuous, homogenous disk and even allowing olivine grains to have low temperatures (mixed with 1% of carbon), we cannot achieve a satisfying fit to the $10 \mu\text{m}$ feature. For these four objects, as already mentioned in the text, we have an on-going Herschel/PACS photometry program that will provide additional constraints on the SEDs of these objects. We will be able to refine our models, and investigate the possibility of additional dust belts. Only then we will be able to discuss more accurately the

geometry of these disks and the possible origin of the transient dust.

The four remaining objects in our sample (HD 69830, HD 169666, HD 98800 B, and ID 8) could be modeled successfully with the assumption of a single dust belt. For the first three objects complementary observations at longer wavelengths do not show any emission in excess, except for the $160 \mu\text{m}$ MIPS photometric point for HD 98800 B. In this case, including larger grains in the model may compensate for the small difference. The absence of far-IR emission in excess for these four sources points towards the scenario of a dust production mechanism directly close to the star. However, determining which declinations of this scenario (massive collision, grinding of an asteroid belt, super-comet) is the correct one is difficult. In the following Section, we will investigate if the results from spectral decomposition can provide additional information on the origin of the dust.

7. The origin of crystalline grains

To have a better understanding of these transient debris disks, one can investigate the origin of crystalline grains. There are several possibilities to crystallize amorphous grains: (i) thermal annealing, where grains are continuously subjected to temperatures high enough (or timescales long enough at intermediate temperatures), for their internal structure to be re-arranged, (ii) shock-induced melting, where grains are subject to high temperatures for short amounts of time, due to collisions. Finally, (iii) crystalline grains detected in the disks may originate from the parent bodies. The parents bodies are either asteroids progressively grind down in a belt, cometary bodies that were scattered inwards in a LHB-like event, or large planetesimals that collided and were completely disrupted. In the latter case, either collisions produce enough heat for crystallization, or crystals were already formed through differentiation in parent bodies and are released through collisions. Scenarios (ii) and (iii) are intimately connected since they both involve collisions between planetesimals or smaller bodies, but in the first case, collisions release enough energy to melt the fragments, while in the second case, the fragments are not altered by the collisions. As debris disks are gas-poor, we do not consider “present-day” gas-phase condensation as a possible significant source of crystalline grains. In the following paragraphs, we present and discuss which diagnostics on crystallization can help us better understand the origin of these grains.

7.1. Where's the enstatite ?

The spectra in our target list do not display any strong, recognizable emission features that can be attributed to enstatite crystalline grains (e.g., at $9.3\ \mu\text{m}$). This translates into small abundances for this particular dust species, while enstatite has been commonly observed around a large variety of primordial disks around Class II objects, with different ages, stellar masses and star forming regions (Bouwman et al. 2008, Juhász et al. 2010, Olofsson et al. 2010, Sargent et al. 2009b). One result of these studies being that enstatite seem to dominate the crystallinity fraction of the inner regions (~ 1 AU for T Tauri stars), while forsterite dominates the outer regions of the disks (≤ 10 – 15 AU for T Tauri stars). Concerning debris disks, so far only one object seems to harbor significantly more enstatite grains than forsterite grains (HD 165014, Fujiwara et al. 2010). This contrast between gas-rich, massive disks and gas-poor debris disk is intriguing. For the objects we study, if the observed dust was the product of a collisional cascade of planetesimals formed within the first AU of the star, we may expect to detect more enstatite dust grains. This is one possible scenario mentioned by Fujiwara et al. (2010), who speculate that the debris disk around HD 165014 may be the aftermath of a collision between small bodies and the surface of a Mercury-like planet (which shows evidence for enstatite, Sprague & Roush 1998). The overall low abundances of enstatite, and the lack of strong emission features, practically rule out such events for the sources in our sample.

The observed crystalline grains may therefore originate from bodies that were formed at larger distances from the star (where enstatite is scarcer), or the observed crystalline grains were recently formed inside the disks. Both possibilities fit rather well the picture of low enstatite abundances. If the debris disks modeled in this study are the results of collisions between bodies originally formed further away from the star, the lack of enstatite grains compares well with observations of younger disks. On the other hand, if the observed crystalline grains are formed at present days in the debris disk, laboratory experiments can help us understanding the low abundances of enstatite. For instance, with their annealing experiments of amorphous grains at 1000 K, Thompson et al. (2002) have shown that enstatite is difficult to form at these temperatures, and forsterite grains are preferentially formed (as probed by X-ray powder diffraction). This result echoes the respective activation energies for annealing of enstatite and forsterite grains (see Gail 2010 and references therein). Several laboratory experiments have shown that activation energies for annealing are larger for enstatite grains compared to forsterite grains (e.g., Fabian et al. 2000). Consequently, if present day thermal annealing is the main production source of crystalline grains in debris disks, we can expect to have higher abundances of crystalline olivine grains compared to enstatite. Interestingly, our finding of low abundances of enstatite echoes the study of crystallization processes during the violent outburst around EX Lupi. Ábrahám et al. (2009) found evidences for production of forsterite, but not of enstatite, within the first AU of the disk, a result later on confirmed by VLTI/MIDI observation (Juhász et al. 2011).

7.2. The iron content

In order to distinguish between the different scenarios on the origin of crystalline grains, the petrology of the olivine dust grains, and more specifically the iron content of these grains, can provide new insights on the crystallization processes at stake in these disks. This means, we have to be confident in our re-

sults about the Fe content of the olivine grains. This study relies on the laboratory measurements for different olivine samples. We insure to use data with as little contamination as possible (aerosol data). Additional effects may come into play, such as the shape of the grains, that can possibly shift peak positions (Koike et al. 2010), but also porosity or crystalline grains being embedded in other materials (Min et al. 2008). Concerning the first point, many olivine grains are injected in the aerosol apparatus, with various sizes and shapes. Plus, we do not make use of any scattering theory, such as the DHS theory that may oversimplify some aspects of the problem (Mutschke et al. 2009). Concerning the last two additional effects, we unfortunately cannot take them into account as this will result in too many degeneracies for an already complex modeling. We would like to stress again the need for high angular resolution observations to constrain some crucial parameters in the modeling. To conclude, we are confident about the detection of Fe-rich olivine grains in several debris disks, given the available data and observations and the assumptions made on the disk geometry (e.g., continuous dust belt). We run additional tests for the two sources with the highest fractions of Fe-rich olivine grains (HD 69830 and HD 113766 A) and the outcomes of these tests are briefly described in Appendix A.

This finding of Fe-rich olivine grains contrasts with observations of Class II disks, for which most studies concluded on the absence of Fe-rich crystalline silicates. In their sample of Herbig Ae/Be stars, based on the $9.4\ \mu\text{m}$ emission feature, Juhász et al. (2010) reported a possible fraction of 10% of Fe in crystalline pyroxene grains. However, the authors do not report any shift to longer wavelengths for emission features associated with crystalline olivine grains (e.g., 19 or $24\ \mu\text{m}$), which indicates an overall small fraction of Fe in olivine grains. To shed light on these petrological differences of olivine grains, we discuss the results of annealing experiments in laboratory. If amorphous dust grains are exposed to high temperatures for a significant amount of time, internal re-arrangement processes are activated in the material and a periodic order appears in the originally disordered network (see Henning 2010). Crystallization is a process that depends on two main environmental parameters: temperature and time. Amorphous grains can crystallize either at intermediate temperatures if the exposure time is long enough, or on short timescales if the temperature is high enough. In the following, we summarize some of the relevant results from previous studies about thermal annealing, with respect to the influence of these two parameters on dust petrology.

Nuth & Johnson (2006) investigated the production of Fe-bearing crystalline grains via thermal annealing. According to their results, both Mg- and Fe-rich crystalline grains can be formed via thermal annealing, if the temperatures are high enough (~ 1100 K and 1400 K, for Mg-rich and Fe-rich, respectively). However, as the temperatures required to form crystalline iron silicates is higher than temperatures required to form crystalline magnesium silicates, the authors argue the Fe-rich grains will evaporate on timescales of several months, while Mg-rich grains can survive longer. According to Nuth & Johnson (2006), this timescale issue is the reason why Fe-rich crystalline grains are not detected in young protoplanetary disks: the grains are destroyed before they can be transported outwards and survive.

Murata et al. (2009) have conducted annealing experiments of Fe-bearing amorphous silicates during 600 min at temperatures of about ~ 1000 K, at a pressure of one atmosphere. The oxygen partial pressure in the furnace was controlled, to avoid possible oxidation of Fe^{2+} cation. They conclude that the Mg

fraction ($=\text{Mg}/[\text{Mg}+\text{Fe}]$) of the crystalline phase decreases with time, as the amorphous phase is annealed. At the end of the experiment, the crystallinity fraction in the original material reached $\sim 80\%$ and the Mg fraction was close to 0.8. Such fraction is quite similar to the Fe-rich olivine grains we use in this study. We can speculate that annealing over long period of time will lower the Mg fraction to match the one of the initial amorphous material.

7.3. Collisions and shock-induced melting

Transient dust observed around debris disks may be the consequence of collisions, where either km-sized bodies are continuously being grind down, or larger planetesimals are completely disrupted after a collision at high relative velocity ($v_{\text{rel}} \sim 10 \text{ km.s}^{-1}$). One can wonder how much energy these disruptive, or grinding collisions release, and by how much the temperature of the resulting fragments can increase. Using laboratory shock experiments as well as numerical simulations Keil et al. (1997) studied the efficiency of shock heating by collisions (with v_{rel} ranging from 3 to 7 km.s^{-1}) for planetesimals with sizes between 10 and 1000 km. Their conclusions were that impacts do not significantly heat up the fragments, mostly because of the low surface gravity of the planetesimals. For instance, they find a globally averaged increase of about 50 K for a collision involving a 1000 km-sized body. But the planetesimals considered by the simulations of Keil et al. (1997) were non-porous bodies. Since many asteroids in the Solar System have non negligible porosity fraction (Britt et al. 2002), more refined models should consider the case of porous bodies. As demonstrated by Davison et al. (2010), including porosity in simulations can make a significant difference. The main conclusions from their study were that before disruption, shock-induced heating is relatively minor and localized, but if the collision is disruptive then heating can be significant, with more than $\sim 10\%$ in mass of the parent body is melted. Collisions between porous planetesimals of comparable sizes can melt almost the entire mass of the two bodies, even at collision velocities smaller than 7 km.s^{-1} . Since many uncertainties remain, applying the results of these simulations to our results is difficult. Nonetheless, we cannot rule out disruptive collisions as a source of newly formed crystalline grains. As the maximum temperature reached during such events is unknown, we can hardly address the question of the petrology of olivine grains when produced by collisions.

7.4. Crystallization in differentiated bodies

Interiors of non-porous planetesimals are known to be regions of high temperatures. The formation of planetary cores requires the melting of accreted bodies, and consequent differentiation. As discussed by Moskovitz & Gaidos (2011), the heating source can be the decay of short-lived radionuclides ^{26}Al and ^{60}Fe , which can result in temperatures high enough for the melting to take place. Differentiated bodies will then be the result of gravitational segregation of molten silicates and metals within the planetesimal. According to studies of iron meteorites (e.g., Scherstén et al. 2006, Qin et al. 2008), such meteorites are the consequence of accretion, differentiation as well as crystallization of their original parent bodies, suggesting that such event most likely happened in the early Solar System, when the decay of ^{26}Al was significant enough to contribute to the heating. Even though the question of Fe-enrichment of olivine grains in differentiated bodies is at present day a poorly constrained question,

we hypothesize that Fe-rich olivine grains may be produced in the cores of young differentiated bodies. Disruptive collisions between these large planetesimals can then release the crystalline grains (and not altering them in non-porous collisions as studied by Keil et al. 1997), an hypothesis that could account for the origin of the chondritic meteorites, that contain Fe-rich olivine grains.

7.5. Individual source discussions

Here, we discuss for each source which scenario can best explain the inferred dust mineralogy, based on the several diagnostics presented previously.

7.5.1. HD 113766 A

According to our best model, r_{in} is located close to the sublimation radius ($T_{\text{d,max}} = 1452 \text{ K}$). Fe-rich olivine grains are found to be more abundant than Mg-rich olivine grains. Since olivine grains produced in Class II disk are preferentially Mg-rich, we are not observing the original mineralogy of parent bodies inherited from the primordial disk (e.g., comets, see Sect. 8) as the origin of the dust around HD 113766 A. Additionally, from IRAS observations, we know that the transient phase has last for at least 25 years, suggesting a continuous replenishment of μm -sized grains. Given the high temperatures, one can hypothesize that crystalline grains are produced via thermal annealing, in a similar way as described in Nuth & Johnson (2006), where crystalline Fe-rich silicates are preferentially formed at high temperatures from Fe-bearing amorphous grains. The timescale issue for the survival of these grains is then comparable to the timescale of evacuation by radiation pressure for the (sub-) μm -sized olivine grains. With this picture in mind, we do not have to assume that all the observed μm -sized *crystalline* grains are continuously replenished from the parent bodies. Instead, we can suggest that a fraction of the small amorphous grains is constantly being annealed into Fe-rich olivine grains. The non-detection of enstatite supports this scenario of present days alteration, since the timescale for production of enstatite may be longer than the timescale for evacuation of the grains.

7.5.2. HD 69830

The dust belt is located slightly further away from the star compared to the sublimation radius of silicate grains ($T_{\text{d,max}} \sim 700 \text{ K}$). Therefore, thermal annealing as described by Nuth & Johnson (2006) cannot account for the abundances of Fe-rich olivines we find. However, radiation pressure has very limited effect on sub- μm -sized grains, meaning the grains survive much longer in the disk. As explained by Murata et al. (2009), and since survival timescales are no longer an issue, long exposure times to intermediate temperatures can produce Fe-rich crystalline grains, annealed from Fe-bearing amorphous grains that dominate the amorphous phase according to our results. We could possibly be witnessing the effect of Mg/Fe fractionation in olivine grains. However, one open question would be to compare the timescales for fractionation and collisions. According to our result the system can be described as a collisional cascade ($p = -3.54$, Dohnanyi 1969), suggesting numerous collisions at all grain sizes. Coagulation of the smallest grains may prevent exposure times long enough for the Mg/Fe fractionation to be efficient (larger grains being colder, fractionation may stop). If the timescale for fractionation is larger than the timescale for co-

agulation, we should observe intermediate Mg fractions values for olivine grains, which is not the case. Additionally, on the basis of the activation energy for Mg_2SiO_4 smoke in Fabian et al. (2000), the characteristic crystallization timescale at 700K is of about 3000 yr. However, at 600K the timescale increases up to 30 Myr, to produce pure Mg-rich forsterite. Since the activation energy of iron-rich silicates should be higher than Mg-rich silicates this timescale can increase significantly for iron-bearing silicates. This renders the possibility of thermal annealing as the main production source for crystalline grains quite unlikely.

As suggested by the absence of an outer belt, the origin of the transient dust around HD 69830 is most likely the result of collisions close to the star. Since Fe-rich olivine grains are unlikely to be formed in Class II disks, we can potentially rule out that the parent bodies are inherited from the primordial disk. Therefore the possibility of an asteroid belt being slowly ground down may not fit very well in this picture. Comets like Hale-Bopp or Wild 2 are thought to have formed in the outer regions of the Solar System, and the crystalline olivine grains observed in both comets are Mg-rich (see Sect. 8). In the scenario of a super-comet being broken in smaller bodies, as described in Beichman et al. (2005), we would therefore expect to preferentially detect Mg-rich olivine grains, formed in the primordial disk and transported outside afterwards. The petrology of the olivine grains we infer in our best model does not support this scenario. Shock-induced melting could contribute as a source of crystalline grains, but it would imply that crystallization via melting strongly favors the production of Fe-rich crystalline grains, a question that is poorly constrained at the moment. However, we can hypothesize that the transient dust originate from differentiated bodies, inside which olivines grains were enriched in Fe. As mentioned by Keil et al. (1997), if such planetesimals are non-porous (a reasonable assumption for differentiation to take place), their breakup into smaller fragments should not alter the mineralogy too much since shock-induced melting then becomes negligible. Interestingly, the system is the host of three Neptune-sized planets. The observed dust belt may be the result of a dynamical instability, possibly triggered by planet migration for instance. Such an instability may have resulted in the collision and disruption of outer large planetesimals.

The small abundance of enstatite may correlate with the abundance of silica ($\text{Mg}_2\text{SiO}_4 + \text{SiO}_2 \rightleftharpoons 2\text{MgSiO}_3$). Production of enstatite via thermal annealing may not fit well in this picture, since the timescale for annealing at these intermediate temperatures (700 K at maximum) is most likely longer than the collisional timescale. The finding of low abundance of enstatite compares well with the overall scarcity of enstatite chondrites in the Solar System ($\sim 2\%$, Norton & Chitwood 2008).

7.5.3. ID 8

The origin of crystalline grains around ID 8 is more difficult to discuss because of the low SNR at longer wavelengths. This renders the spectral decomposition more uncertain on the exact crystallinity fraction as well as on the iron fraction in olivine grains. The uncertainties from the Monte-Carlo runs show that the detection of Fe-rich over Mg-rich olivine grains is less convincing compared to the previous two sources. We therefore prefer not to conclude on the exact petrology of the crystalline olivine grains in this system. Still, one of the interesting characteristic of the modeling results is the very steep grain size distribution ($p = -4.01$, for s up to $1000\mu\text{m}$). This suggests the system is at a peak of collisional activity, therefore producing large

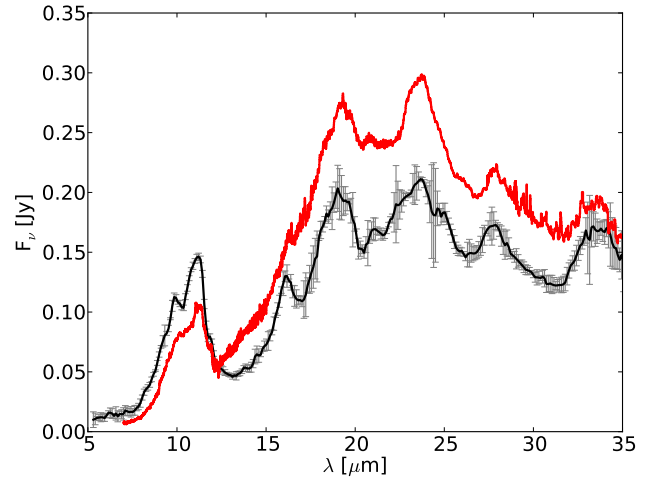


Fig. 10. Stellar subtracted Irs spectrum of ID 9 (black) and ISO spectrum of Hale-Bopp (in red, Crovisier et al. 1997).

amounts of sub- μm -sized grains. At the same time, it implies that collision timescales grains are short. This could explain the rather small crystallinity fraction we derive, as crystalline grains may collide and coagulate with larger grains in short amount of time, minimizing their relative contribution to the emission features.

7.5.4. ID 9

As discussed in Sect. 5.7, cold Mg-rich olivine grains are required to reproduce the data, and the disk extends up to ~ 50 AU according to our best model. Even though the maximum temperature is of about 950 K the strongest emission features are at long wavelengths (23, 28 and $33\mu\text{m}$) suggesting cold olivine grains must contribute significantly to the emission. In Sect. 6.3 we discuss the possibility of a second belt, further away from the star that can be responsible for the emission at longer wavelengths. With this assumption, that needs to be confirmed by far-IR observations, we can suggest a LHB-like scenario to explain the observed spectrum of ID 9. If the outer belt has been dynamically excited (e.g., planet migration) a collisional cascade may have been initiated, releasing a large amount of Mg-rich olivine grains. As radiation pressure is not efficient for μm -sized grains, they are able to survive on long timescales. Simultaneously to the collisional cascade, bodies from the outer system may be scattered inside the system, releasing Mg-rich olivine grains as well, via collisions and sublimation close to the star. The detection of Mg-rich, instead of Fe-rich, olivine grains could trace the mineralogy inherited from the primordial protoplanetary disk where crystalline iron-rich silicates do not survive long enough. For example, cold crystalline olivine grains have been detected in the disk around HD 100546 (see Mulders et al. 2011, Sturm et al. 2010). Studying the peak position of the $69\mu\text{m}$ feature, the authors concluded that the iron content must not be higher than 1%. This solution of a LHB-like scenario will be further investigated, when Herschel/PACS observations will be available. Interestingly, the resemblance of the emission features and their corresponding peak positions, between the spectrum of ID 9 and the one of Hale-Bopp (Figure 10) echoes the youth of our Solar System and further supports the possibility of an outer belt, similar to the Kuiper-belt.

7.5.5. HD 15407 A

The spectrum of HD 15407 A does not show emission features associated either with olivine or enstatite grains. But the spectrum shows several emission features associated with β -cristobalite. The nature of this particular SiO_2 polymorph is derived from the presence of the $16\ \mu\text{m}$ feature, which is usually not visible for other polymorphs of silica. It is possible that other polymorphs are present in the debris disk around HD 15407 A, but inferring their presence is a degenerate problem given the strong similarities between the Q_{abs} values, especially for the strongest emission feature, around $9\ \mu\text{m}$. The β -cristobalite form of silica we use is close to the annealed silica studied by Fabian et al. (2000), which is the result of an annealing experiment, for 5 h at a temperature of 1220 K. This indicates that the grains have experienced rather high temperatures, but as discussed in Sargent et al. (2009a), once formed cristobalite grains must be cooled quickly enough to retain their crystalline structure. If not, the grains are expected to revert to the lower temperature polymorphs of silica (β - and α -quartz).

The mineralogy of the dust around this source is puzzling. No crystalline silicate grains are detected (olivine or enstatite), but a significant fraction of SiO_2 is required to model the spectrum. According to our best model, the maximum temperature is of about 1530 K, a temperature high enough to anneal silicates grains and produce for example crystalline olivine grains. Even though small grains are short-lived around HD 15407 A because of radiation pressure, if the main source of silica was thermal annealing, we should expect to detect features associated with olivine grains (as in the case of HD 113766 A for instance). Their non-detection points towards a different source than thermal annealing, which in turn means that the maximum temperature we find is not appropriate and that grains should be colder than ~ 1500 K. Since the derived geometry of the disk is dominated by the fitting of emission features and given that there is only one strong emission feature around $9\ \mu\text{m}$ in the entire spectrum of HD 15407 A, this can explain why the temperature determination may not be accurate enough in our best model. Additionally, we discuss previously the possibility of a cold, outer dust belt further away from the central star. If confirmed, the geometry of the disk would have to be revised, and hence at the same time the maximum temperature of the grains. This temperature issue can be easily constrained with high-angular resolution observations. From the absence of crystalline olivine grains, ruling out thermal annealing as a source of crystallization process leaves us with little options to explain the peculiar dust composition of this dust belt. Since such crystalline grains are expected to be common in the outer regions of primordial disks where comets are formed, and since silica is not usually detected in cometary bodies, the scenario of an outer belt feeding the inner regions is rather unlikely. We may therefore be witnessing the aftermath of collisions between silica-rich parent bodies. But the collisions should be either non-disruptive, or the parent bodies should be non-porous, otherwise shock-induced heating may produce other crystalline grains, as discussed previously. In that scenario, the collision may be the consequence of dynamical stirring, triggered by the secondary of the binary system.

7.5.6. Other sources

For the three remaining sources, we can hardly discuss the origin of the crystalline grains or the petrology of olivine grains. For HD 98800 B, we conclude in Sect. 5.6 that the crystallinity fraction could only be considered as tentative because of the lack of

strong emission features associated with crystalline grains. For HD 169666, the overall low SNR of the data renders the analysis difficult, and we prefer not to over-interpret data. Concerning BD+20 307, since the fit is not satisfying around $10\ \mu\text{m}$, the confidence level of the crystallinity fraction and iron content is not sufficient for any interpretation of their origin.

Overall, the conclusion of this Section 7 is that many processes are active in these transient debris disks, from thermal annealing at high or intermediate temperatures, to LHB-like event or (non-)disruptive collisions of (non-)porous parent bodies with various compositions. Spectroscopic observations contain a finite amount of informations, and we have a limited knowledge about dust emissivity properties (from near- to mid-IR). We are able to suggest possible scenarios to interpret our results within these unfortunate limitations but further investigations (observations, instrumentation and laboratory experiments) are mandatory to confirm, or infirm our findings. To summarize, it is not trivial to assume that the inferred dust composition originates from parent bodies inherited from the primordial disk because of several other crystallization processes at stake.

8. Comparison with Solar System objects

In this Section, we discuss how our results compare with studies of bodies that were formed during the youth of our Sun. For instance, very few chondrites ($\sim 2\%$) that fall on Earth seem to contain enstatite, a result consistent with our modeling results (Norton & Chitwood 2008). The Fe content is also an interesting diagnostic to focus on, as it is quite a novel finding for dust in exo-zodiacal belts.

8.1. Cometary bodies

As detailed in the reviews by Hanner & Zolensky (2010) and Wooden (2008) comets are formed in the outermost regions (Oort Cloud or the Kuiper Belt region) and trace the youth of the Solar System. Informations about their dust composition is a prime objective in order to better understand the formation of our Solar System. Using spectroscopic observations at different epochs, Wooden et al. (1999) studied the spectrum of the comet C/1995 01 (Hale-Bopp). According to their modeling they found both Mg-rich olivine and crystalline pyroxene grains to be present in the coma, but no clear evidences for Fe-rich olivine grains. The Mg-rich pyroxene grains were required to reproduce the $9.3\ \mu\text{m}$ feature, which is not detected in any of our sources. The modeling results for the comet Hale-Bopp are very similar to the ones for the Oort Cloud comet C/2001 Q4 (Wooden et al. 2004). For both comets, the mineralogical findings echo the results from the Stardust mission, on comet P81/Wild 2. Based on samples collected in situ, Zolensky et al. (2008) found that olivines in the comet mostly consist of Mg-rich olivine grains, with a Mg fraction ($=\text{Mg}/[\text{Mg} + \text{Fe}]$) of about 99. However, they also found the composition range of olivine to be extremely broad, ranging from Mg fractions between 4 and 100. From the peak position of the Mg distribution, the author conclude that composition changes due to capture heating must have been insignificant, otherwise the Mg distribution should instead peak around 40–60. Interestingly, from the peak position of the Mg distribution, the authors suggest that the petrology of the olivine grains can be explained by aqueous alteration; Fe-rich olivine grains may not survive as easily as Mg-rich grains in presence of water (Wogelius & Walther 1992). The olivine grains may either trace the eventual presence of liquid water on P81/Wild 2, or trace different regions of the initial disk, where Mg-rich olivine

grains have preferentially survived aqueous alteration compared to Fe-rich olivine grains. The fact that crystalline grains detected in cometary bodies are Mg-rich rather than Fe-rich is consistent with observations of gas-rich Class II disks. Comets are expected to trace the mineralogy of the outermost regions of the primordial disk, therefore if Fe-rich olivine grains, as we detect them in debris disks, were already formed during the Class II phase, we should in principle observe significant fractions of them in comets as well. For the couple of objects for which we find more Fe-rich than Mg-rich olivine grains, this points towards a new generation of dust in the transient phase of the disks, which is subjected to alternative crystallization processes, more inclined to form Fe-rich crystalline grains.

8.2. Interplanetary Dust Particles

The exact origin of Interplanetary dust Particles (IDPs) is difficult to track, they could have been trapped in asteroids or short-period comets (originating from the Kuiper Belt, see Bradley 2010 for a review). It is in principle possible to distinguish between an asteroid or comet origin, based on He release during their entry in the atmosphere (Brownlee et al. 1995), but the place of formation of these IDPs is thought to be presolar, based on isotopic considerations. An interesting point with respect to our study is that olivine grains contained in several IDPs are found to be Fe-rich (e.g., Bradley 1994, Brunetto et al. 2011). According to Zolensky & Barrett (1994) the Mg distribution for hydrous IDPs ranges between 76 and 100, while it ranges between 52 and 100 for anhydrous IDPs, these differences being the possible consequence of aqueous alteration. Nguyen et al. (2007) discuss the origin of Fe in IDPs and suggest that the Fe enrichment may be the consequence of a secondary process taking place in the solar nebula or on the parent body, which echo our findings. Nevertheless, Nguyen et al. (2007) also discuss the possibility of condensation under non-equilibrium conditions in stellar outflows to explain the detection of Fe-rich crystalline grains. Overall, the existence of Fe-rich olivine grains in various astronomical environments is not to be questioned, the novelty of our study is to detect such grains in several exo-zodiacal belts. Such grains are not expected to be formed in gas-rich Class II disks, their detection in debris disks points toward alternative crystallization process of the transient phase of the disks.

9. Conclusion

In this study, we present a new and powerful radiative transfer code, dedicated to debris disks (DEBRA). Combining this code with recent laboratory experiments on dust mass absorption coefficients we model the Spitzer/IRs spectra of 8 debris disks around solar-analogs, in order to derive the mineralogy of the dust. The entire SEDs, from optical to millimeter wavelengths, are successfully modeled within the optically thin regime. Based on the modeling of these rare objects, we obtain the following results:

1. To contribute significantly in the mid-IR, grains must be “dirty”, especially grains with low mass absorption coefficients in the near-IR, such as Fe-free amorphous grains, enstatite and crystalline olivine grains. If Q_{abs} values are too low at the wavelengths where the stellar radiation peaks, grains will be almost transparent and consequently too cold for their emission to be significant. Even though the crystalline olivine grains we use in this study contain some Fe ($\text{Mg}_{2x}\text{Fe}_{2(1-x)}\text{SiO}_4$, with $x = 0.95$ or 0.8), Zeidler et al. (2011) have demonstrated their mass absorption coefficients to be very low in the near-IR. Adding impurities, such as carbon, helps circumventing this issue. We find that usually a fraction in mass of 1 or 5% is sufficient for the crystalline grains. One should note that we do not assume the dust species to be in thermal contact. Each grain, for all sizes and all compositions has its temperature determined as a function of its distance r and mass absorption coefficients.
2. Signal-to-noise ratio has a significant influence on the modeling results (see Juhász et al. 2009 as well). The case of HD 169666 is a good example, as we are able to find two comparable solutions (in terms of reduced χ_r^2), but with completely different results on the location of the dust belt. In the case of ID 8, where SNR decreases at longer wavelengths, the relative fractions of Fe- versus Mg-rich olivine grains become uncertain. For the remaining sources, the high SNR ensures a reliable determination of the dust mineralogy, as probed by the overall low uncertainties we find.
3. In the formalism of Wyatt et al. (2007), using the ratio between observed and maximal fractional luminosities, we find that most, if not all, of the debris disks we study are in a transient phase, that cannot be explained by a steady-state evolution from the Class II phase, meaning that μm -sized dust grains were produced relatively recently compared to the stellar ages.
4. At least 5 out of 8 sources are either in binary systems or have planets orbiting the central star, suggesting the transient phase might have been triggered by dynamical stirring from these massive bodies.
5. Computing the radiation pressure efficiency β_{rp} , we find that μm -sized grains are short-lived around at least three objects (HD 15407 A, HD 113766 A and HD 169666). Investigating their variability in time via archive mining, we conclude there must be an on-going process that replenishes efficiently the population of smallest grains on short dynamical timescales (on the order of months). The dust grains around the remaining debris disks are not subjected to evacuation by radiation pressure, meaning the disk can survive much longer, on the timescale of Poynting-Robertson drag.
6. According to far-IR observations, only one source in our sample (HD 113766 A) harbors a cold, outer belt. Based either on disk properties from our modeling results, or on mineralogical considerations, we discuss the possibility of outer belts around three other sources (HD 15407 A, ID 9 and BD+20 307). An on-going Herschel/Pacs photometry program will enable us to investigate deeper this scenario. The four remaining sources could successfully be modeled by one single continuous dust belt.
7. We find that relative abundances of enstatite grains are overall small for all the sources we study. This is the consequence of non-detections of recognizable emission features associated with pyroxene crystalline grains (especially at $9.3\mu\text{m}$). If the observed crystallinity is the consequence of present day crystallization processes, this may be explained by higher activation energies for annealing compared to crystalline olivine grains. If the observed crystallinity fraction is inherited from parent bodies that collided (with no shock-induced heating), this echoes the scarcity of enstatite chondrites in our own Solar System.
8. The quality of the fits we present in this study is intimately connected to the Q_{abs} values we use in the modeling, especially for olivine grains (which are responsible for the strongest emission features in the spectra we model). This underlines the extreme importance of laboratory experiments

- to improve our understanding of cosmic dust properties. The aerosol technique is a promising technique that provide direct measurements of Q_{abs} values, with as less contamination from the environment as possible (e.g., Tamanai et al. 2009).
9. For several sources, especially HD 113766 A and HD 69830, we find that crystalline olivine grains contain Fe inclusions, up to 20% compared to Mg. Fe-rich crystalline grains are usually not observed in Class II protoplanetary disks, which can be explained by several studies of crystallization processes (e.g., Gail 2004, Nuth & Johnson 2006, Murata et al. 2009). Therefore our findings of more Fe-rich compared to Mg-rich olivine grains points towards a new generation of crystalline dust grains, for which the crystallization processes are slightly different in debris disks than in Class II disks, leading to a higher fraction of Fe.
 10. We discuss the possibility of shock-induced heating by collisions as an eventual source of crystalline grains. While many uncertainties remain on the outcomes of collisions, it appears that disruptive collisions between porous parental bodies can generate enough energy to melt fragments. Overall, with the previous point, it seems unlikely that the inferred mineralogy of the dust traces the mineralogy of the parent bodies, but instead that we are witnessing the production of a second generation of dust, different from the primordial one inherited from the protoplanetary disk.
 11. Finally, based on several considerations (extension of dust belt, strong emission features at long wavelengths and mineralogy), we suggest that a LHB-like event is happening around one of our source (ID 9). Further modeling including PACS photometry will provide additional informations on the system.

As a conclusion, we would like to stress the fact that SED modeling is a challenging problem, for protoplanetary and debris disks. Since debris disks are optically thin at all wavelengths, the problem becomes even more difficult as several assumptions no longer hold, such as thermal contact between dust grains (which renders the temperature determination critical) or the continuum emission as commonly used when modeling spectra of Class II disks. Spectroscopic observations do not contain any spatial informations, which introduces degeneracies in the modeling (e.g., dust temperature versus dirtiness of the dust grains), therefore more refined modeling would require high angular resolution observations (e.g., near-IR interferometry or scattered light imaging, Lebreton et al. 2011) to lower the number of degenerate parameters (r_{in} , r_{out} and possibly α). Further work is also required on the dust properties to improve the modeling results, not only with laboratory experiments (e.g., Q_{abs} values with aerosol technique for other dust species), but also to investigate additional effects such as grains porosity, or inhomogeneous aggregates of different species (as described in Min et al. 2008 and Voshchinnikov & Henning 2008). This latter point may eventually provide a suitable solution for disks where μm -sized grains are subjected to efficient radiation pressure. If small grains are embedded in larger grains, they may survive longer in the dust belt. Unfortunately, this increases the complexity of the modeling, and again underlines the need for additional constraints on the disk geometry. With observations and Q_{abs} values available at present days, we are able to successfully model the IRS spectra of several warm debris disks, with a limited amount of dust species, and within the limitations of our model, we are able to provide new insights on the origin of this transient dust.

Acknowledgements. We are grateful to J. Bouwman for providing the ISO spectrum of Hale-Bopp. J. O. acknowledges the Alexander von Humboldt foundation

for critical financial support. A.M. and P.Á. appreciate financial support from the Hungarian Research Fund OTKA 101393.

References

- Ábrahám, P., Juhász, A., Dullemond, C. P., et al. 2009, *Nature*, 459, 224
 Aumann, H. H., Beichman, C. A., Gillett, F. C., et al. 1984, *ApJ*, 278, L23
 Begemann, B., Dorschner, J., Henning, T., Mutschke, H., & Thamm, E. 1994, *ApJ*, 423, L71
 Beichman, C. A., Bryden, G., Gautier, T. N., et al. 2005, *ApJ*, 626, 1061
 Beichman, C. A., Lisse, C. M., Tanner, A. M., et al. 2011, *ArXiv e-prints*
 Bouwman, J., Henning, T., Hillenbrand, L. A., et al. 2008, *ApJ*, 683, 479
 Bouwman, J., Meeus, G., de Koter, A., et al. 2001, *A&A*, 375, 950
 Bradley, J. 2010, in *Lecture Notes in Physics*, Berlin Springer Verlag, Vol. 815, Lecture Notes in Physics, Berlin Springer Verlag, ed. T. Henning, 259–276
 Bradley, J. P. 1994, *Science*, 265, 925
 Britt, D. T., Yeomans, D., Housen, K., & Consolmagno, G. 2002, *Asteroids III*, 485
 Brownlee, D. E., Joswiak, D. J., Schlutter, D. J., et al. 1995, in *Lunar and Planetary Institute Science Conference Abstracts*, Vol. 26, Lunar and Planetary Institute Science Conference Abstracts, 183
 Brunetto, R., Borg, J., Dartois, E., et al. 2011, *Icarus*, 212, 896
 Bryden, G., Beichman, C. A., Trilling, D. E., et al. 2006, *ApJ*, 636, 1098
 Burns, J. A., Lamy, P. L., & Soter, S. 1979, *Icarus*, 40, 1
 Carpenter, J. M., Bouwman, J., Mamajek, E. E., et al. 2009, *ApJS*, 181, 197
 Charbonneau, P. 1995, *ApJS*, 101, 309
 Chen, C. H., Mamajek, E. E., Bitner, M. A., et al. 2011, *ApJ*, 738, 122
 Chen, C. H., Patten, B. M., Werner, M. W., et al. 2005, *ApJ*, 634, 1372
 Chen, C. H., Sargent, B. A., Bohac, C., et al. 2006, *ApJS*, 166, 351
 Crovisier, J., Leech, K., Bockelee-Morvan, D., et al. 1997, *Science*, 275, 1904
 Davison, T. M., Collins, G. S., & Ciesla, F. J. 2010, *Icarus*, 208, 468
 Dohnanyi, J. S. 1969, *J. Geophys. Res.*, 74, 2531
 Dorschner, J., Begemann, B., Henning, T., Jaeger, C., & Mutschke, H. 1995, *A&A*, 300, 503
 Egan, W. G. & Hilgeman, T. 1977, *Icarus*, 30, 413
 Fabian, D., Jäger, C., Henning, T., Dorschner, J., & Mutschke, H. 2000, *A&A*, 364, 282
 Fujiwara, H., Onaka, T., Ishihara, D., et al. 2010, *ApJ*, 714, L152
 Furlan, E., Sargent, B., Calvet, N., et al. 2007, *ApJ*, 664, 1176
 Gail, H.-P. 2004, *A&A*, 413, 571
 Gail, H.-P. 2010, in *Lecture Notes in Physics*, Berlin Springer Verlag, Vol. 815, Lecture Notes in Physics, Berlin Springer Verlag, ed. T. Henning, 61–141
 Gorlova, N., Balog, Z., Rieke, G. H., et al. 2007, *ApJ*, 670, 516
 Hanner, M. S. & Zolensky, M. E. 2010, in *Lecture Notes in Physics*, Berlin Springer Verlag, Vol. 815, Lecture Notes in Physics, Berlin Springer Verlag, ed. T. Henning, 203–232
 Henning, T. 2010, *ARA&A*, 48, 21
 Henning, T. & Mutschke, H. 1997, *A&A*, 327, 743
 Henning, T. & Mutschke, H. 2010, *Journal of Nanophotonics*, 4, 041580
 Jäger, C., Dorschner, J., Mutschke, H., Posch, T., & Henning, T. 2003, *A&A*, 408, 193
 Jäger, C., Molster, F. J., Dorschner, J., et al. 1998a, *A&A*, 339, 904
 Jäger, C., Mutschke, H., & Henning, T. 1998b, *A&A*, 332, 291
 Juhász, A., Bouwman, J., Henning, T., et al. 2010, *ApJ*, 721, 431
 Juhász, A., Dullemond, C., van Boekel, R., et al. 2011, *ArXiv e-prints*
 Juhász, A., Henning, T., Bouwman, J., et al. 2009, *ApJ*, 695, 1024
 Keil, K., Stoeffler, D., Love, S. G., & Scott, E. R. D. 1997, *Meteoritics and Planetary Science*, 32, 349
 Kemper, F., Vriend, W. J., & Tielens, A. G. G. M. 2005, *ApJ*, 633, 534
 Koerner, D. W., Jensen, E. L. N., Cruz, K. L., Guild, T. B., & Gultekin, K. 2000, *ApJ*, 533, L37
 Koike, C., Chihara, H., Tsuchiyama, A., et al. 2003, *A&A*, 399, 1101
 Koike, C., Imai, Y., Chihara, H., et al. 2010, *ApJ*, 709, 983
 Krivov, A. V. 2010, *Research in Astronomy and Astrophysics*, 10, 383
 Lebreton, J., Augereau, J.-C., Thi, W.-F., et al. 2011, *ArXiv e-prints*
 Lisse, C. M., Beichman, C. A., Bryden, G., & Wyatt, M. C. 2007, *ApJ*, 658, 584
 Lisse, C. M., Chen, C. H., Wyatt, M. C., & Morlok, A. 2008, *ApJ*, 673, 1106
 Lisse, C. M., Chen, C. H., Wyatt, M. C., et al. 2009, *ApJ*, 701, 2019
 Lisse, C. M., VanCleve, J., Adams, A. C., et al. 2006, *Science*, 313, 635
 Lovis, C., Mayor, M., Pepe, F., et al. 2006, *Nature*, 441, 305
 Low, F. J., Smith, P. S., Werner, M., et al. 2005, *ApJ*, 631, 1170
 Matraga, M., Drake, J. J., Kashyap, V. L., Marengo, M., & Kuchner, M. J. 2010, *ApJ*, 720, L164
 Melis, C., Zuckerman, B., Rhee, J. H., & Song, I. 2010, *ApJ*, 717, L57
 Min, M., Hovenier, J. W., Waters, L. B. F. M., & de Koter, A. 2008, *A&A*, 489, 135
 Min, M., Waters, L. B. F. M., de Koter, A., et al. 2007, *A&A*, 462, 667

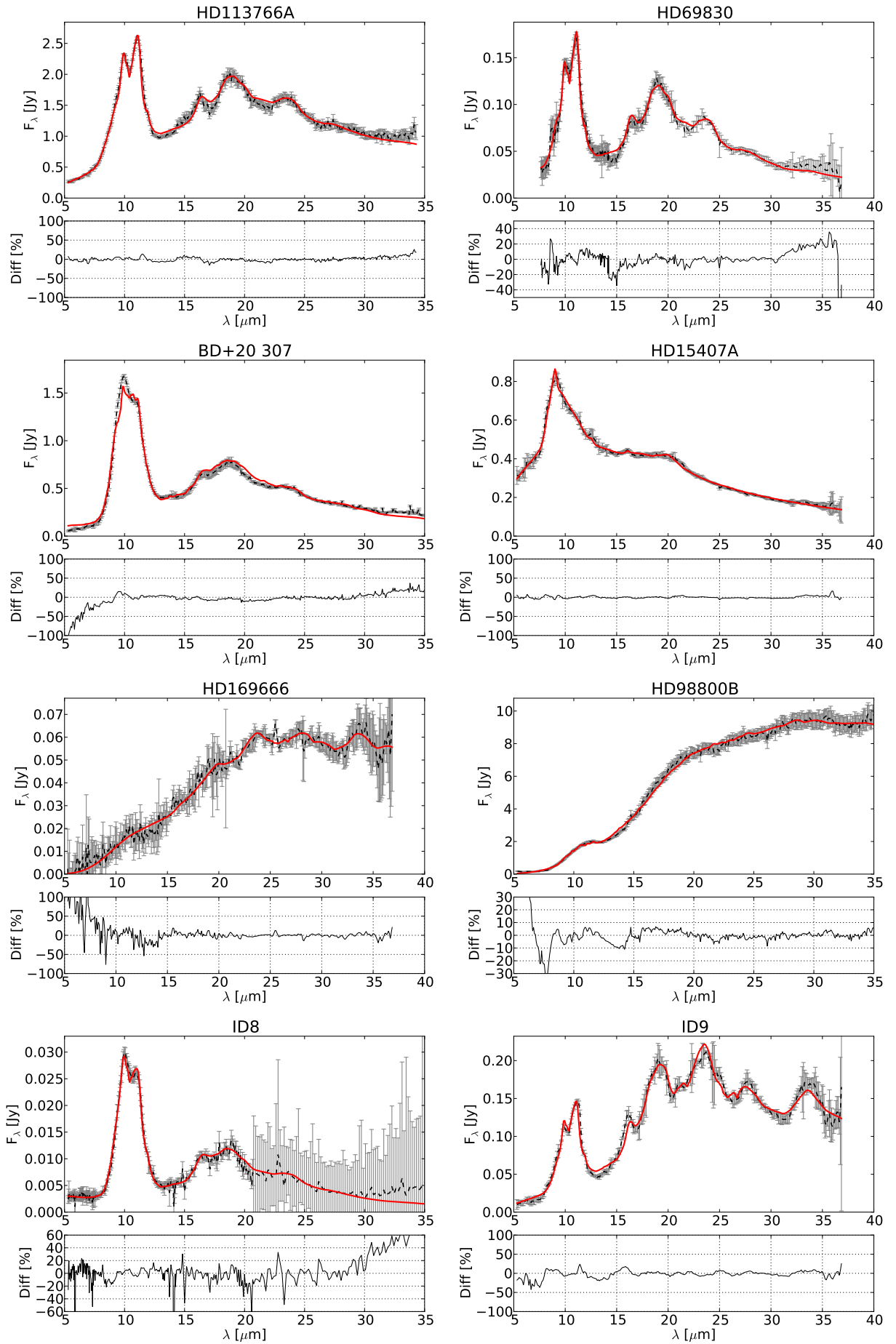


Fig. 11. Best fits (in red) to the Irs spectra (dashed lines). Uncertainties are $3\text{-}\sigma$. For each source bottom panels show the residuals.

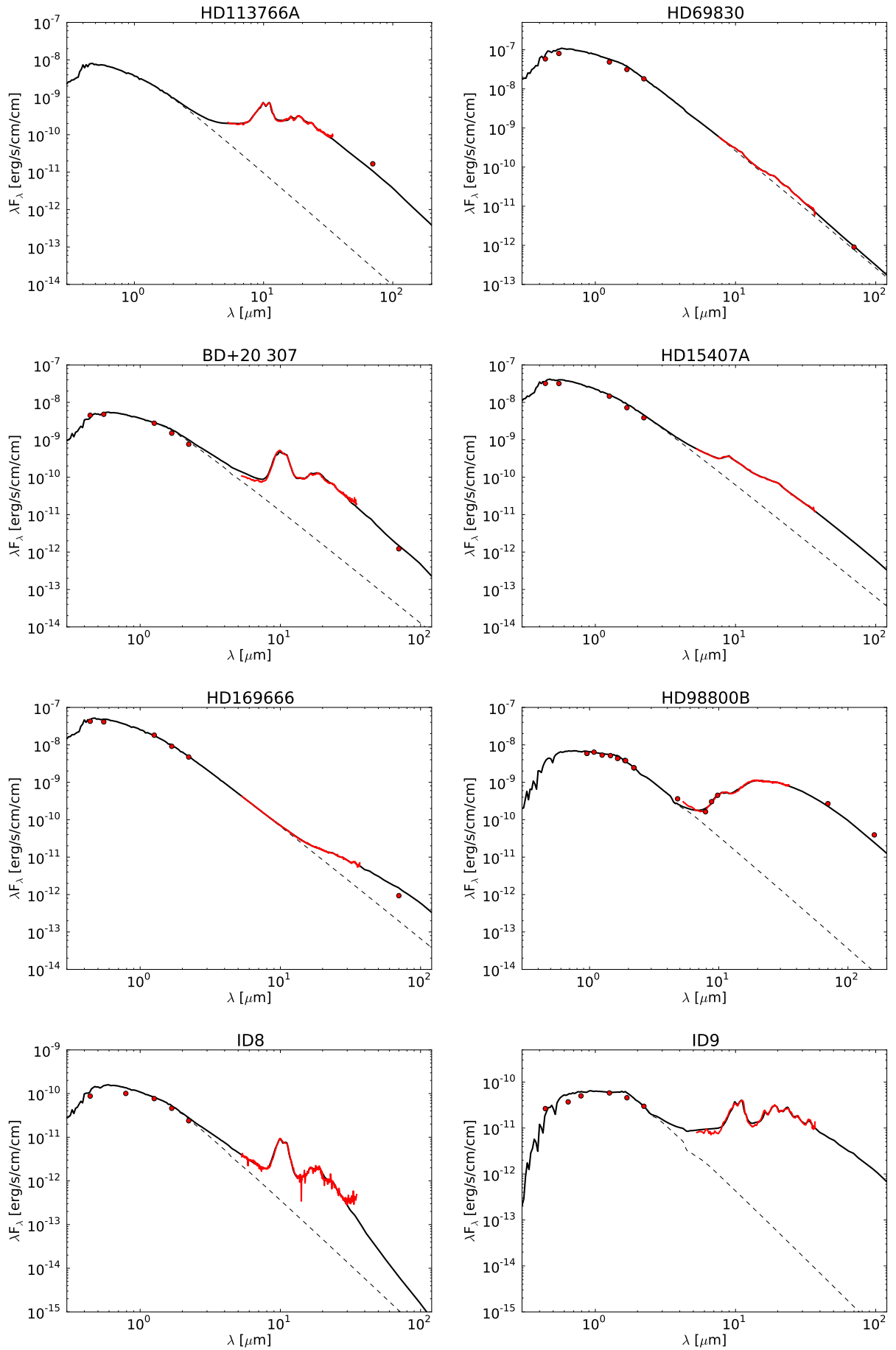


Fig. 12. Complete SEDs for all the sources, with Irs spectra and photometric observations in red. Modeled SEDs are in solid black.

- Moór, A., Apai, D., Pascucci, I., et al. 2009, *ApJ*, 700, L25
- Moskovitz, N. & Gaidos, E. 2011, *Meteoritics and Planetary Science*, 46, 903
- Mulders, G. D., Waters, L. B. F. M., Dominik, C., et al. 2011, *A&A*, 531, A93
- Murata, K., Takakura, T., Chihara, H., Koike, C., & Tsuchiyama, A. 2009, *ApJ*, 696, 1612
- Mutschke, H., Min, M., & Tamanai, A. 2009, *A&A*, 504, 875
- Nguyen, A. N., Stadermann, F. J., Zinner, E., et al. 2007, *ApJ*, 656, 1223
- Norton, O. R. & Chitwood, L. 2008, *Field Guide to Meteors and Meteorites*, ed. Norton, O. R. & Chitwood, L.
- Nuth, J. A. & Johnson, N. M. 2006, *Icarus*, 180, 243
- Oliveira, I., Olofsson, J., Pontoppidan, K. M., et al. 2011, *ApJ*, 734, 51
- Olofsson, J., Augereau, J.-C., van Dishoeck, E. F., et al. 2010, *A&A*, 520, A39+
- Olofsson, J., Augereau, J.-C., van Dishoeck, E. F., et al. 2009, *A&A*, 507, 327
- Ossenkopf, V., Henning, T., & Mathis, J. S. 1992, *A&A*, 261, 567
- Prato, L., Ghez, A. M., Piña, R. K., et al. 2001, *ApJ*, 549, 590
- Qin, L., Dauphas, N., Wadhwa, M., Masarik, J., & Janney, P. E. 2008, *Earth and Planetary Science Letters*, 273, 94
- Sargent, B. A., Forrest, W. J., Tayrien, C., et al. 2009a, *ApJ*, 690, 1193
- Sargent, B. A., Forrest, W. J., Tayrien, C., et al. 2009b, *ApJS*, 182, 477
- Scherstén, A., Elliott, T., Hawkesworth, C., Russell, S., & Masarik, J. 2006, *Earth and Planetary Science Letters*, 241, 530
- Smith, R., Wyatt, M. C., & Haniff, C. A. 2009, *A&A*, 503, 265
- Soderblom, D. R., King, J. R., Siess, L., et al. 1998, *ApJ*, 498, 385
- Sogawa, H., Koike, C., Chihara, H., et al. 2006, *A&A*, 451, 357
- Sprague, A. L. & Roush, T. L. 1998, *Icarus*, 133, 174
- Sturm, B., Bouwman, J., Henning, T., et al. 2010, *A&A*, 518, L129
- Tamanai, A. 2010, in "Dust in Planetary Systems, Workshop, Jena, Germany"
- Tamanai, A. & Mutschke, H. 2010, *Infrared spectroscopy of dust particles in aerosols for astronomical application (Fundamentals and Applications in Aerosol Spectroscopy*, eds. J. Reid and R. Signorell, pp.101-123 (Taylor & Francis Group: Florida))
- Tamanai, A., Mutschke, H., Blum, J., & Mees, G. 2006, *ApJ*, 648, L147
- Tamanai, A., Mutschke, H., Blum, J., et al. 2009, *A&A*, 501, 251
- Thompson, S. P., Fonti, S., Verrienti, C., et al. 2002, *A&A*, 395, 705
- van Boekel, R., Min, M., Waters, L. B. F. M., et al. 2005, *A&A*, 437, 189
- Verrier, P. E. & Evans, N. W. 2008, *MNRAS*, 390, 1377
- Voshchinnikov, N. V. & Henning, T. 2008, *A&A*, 483, L9
- Weinberger, A. J. 2008, *ApJ*, 679, L41
- Weinberger, A. J., Becklin, E. E., Song, I., & Zuckerman, B. 2011, *ApJ*, 726, 72
- Wogelius, R. A. & Walther, J. V. 1992, *Chemical Geology*, 97, 101
- Wolf, S. & Hillenbrand, L. A. 2005, *Computer Physics Communications*, 171, 208
- Wooden, D. H. 2008, *Space Sci. Rev.*, 138, 75
- Wooden, D. H., Harker, D. E., Woodward, C. E., et al. 1999, *ApJ*, 517, 1034
- Wooden, D. H., Woodward, C. E., & Harker, D. E. 2004, *ApJ*, 612, L77
- Wyatt, M. C. 2008, *ARA&A*, 46, 339
- Wyatt, M. C., Smith, R., Greaves, J. S., et al. 2007, *ApJ*, 658, 569
- Zeidler, S., Posch, T., Mutschke, H., Richter, H., & Wehrhan, O. 2011, *A&A*, 526, A68+
- Zolensky, M., Nakamura-Messenger, K., Rietmeijer, F., et al. 2008, *Meteoritics and Planetary Science*, 43, 261
- Zolensky, M. E. & Barrett, R. A. 1994, *Meteoritics*, 29, 616
- Zuckerman, B., Fekel, F. C., Williamson, M. H., Henry, G. W., & Muno, M. P. 2008, *ApJ*, 688, 1345

Appendix A: Detection of Fe-rich olivine grains

To assess the confidence level of the detection of Fe-rich olivine grains rather than Mg-rich olivine grains for the two sources where we detect such grains, we run the *DEBRA* code in the exact same way as discussed previously, but without including Fe-rich olivine grains. The results for HD 69830 and HD 113766 A are shown in Fig. A.1, with fits to the *Irs* spectra with and without Fe-rich olivine grains (left and right panels, respectively). Since the differences between Mg- and Fe-rich Q_{abs} values are overall small (see Fig. 2), the fits to the data are still reasonably good without the Fe-rich olivine grains. However, the peak positions and shapes of features at 19 and 24 μm are best reproduced when including them, comforting us in the detection of such grains.

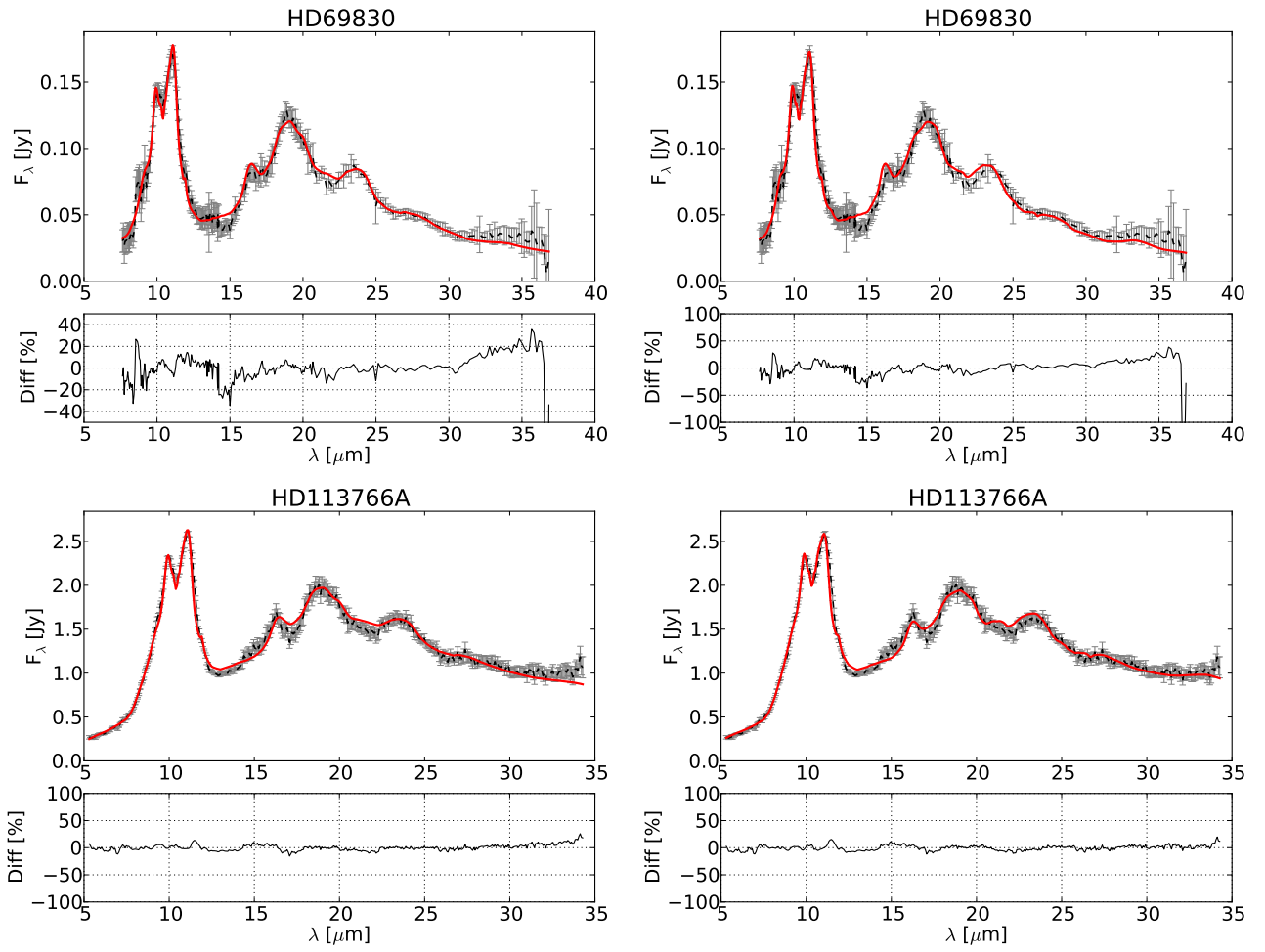


Fig. A.1. *Left panels:* fits to stellar subtracted spectra of HD 69830 and HD 113766 A (*top* and *bottom*, respectively), including Fe-rich olivine grains. *Right panels:* same, but without Fe-rich olivine grains included in the modeling.